

TEMPERED FRACTIONAL ORNSTEIN–UHLENBECK PROCESSES

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Tempered fractional Brownian motion is a modification of fractional Brownian motion obtained by introducing exponential tempering into the Mandelbrot–Van Ness representation. It preserves the local fractional behavior and weakens dependence at large lags.

Let $H > 0$, $\lambda > 0$, $\theta > 0$, $x_+ = \max\{x, 0\}$, and let $W = \{W_x, x \in \mathbb{R}\}$ be a two-sided Wiener process. The parameter $k \in \{\text{I}, \text{II}\}$ indicates the type of tempered fractional Brownian motion: the first kind [4] or the second kind [5]. In both cases

$$B_{H,\lambda}^k(t) = \int_{\mathbb{R}} g_{H,\lambda,t}^k(x) dW_x,$$

where

$$g_{H,\lambda,t}^{\text{I}}(x) = e^{-\lambda(t-x)_+} (t-x)_+^{H-1/2} - e^{-\lambda(-x)_+} (-x)_+^{H-1/2},$$

$$g_{H,\lambda,t}^{\text{II}}(x) = g_{H,\lambda,t}^{\text{I}}(x) + \lambda \int_0^t e^{-\lambda(u-x)_+} (u-x)_+^{H-1/2} du.$$

We study stationary and non-stationary Ornstein–Uhlenbeck processes driven by $B_{H,\lambda}^k$; classical fractional Ornstein–Uhlenbeck processes were considered in [1]. These processes solve the Langevin-type equation $dX_t = -\theta X_t dt + dB_{H,\lambda}^k(t)$. The stationary solution has the form

$$Y_t^k = B_{H,\lambda}^k(t) - \theta e^{-\theta t} \int_{-\infty}^t e^{\theta s} B_{H,\lambda}^k(s) ds, \quad t \geq 0,$$

whereas the non-stationary solution with initial condition $X_0^k = 0$ is given by

$$X_t^k = B_{H,\lambda}^k(t) - \theta e^{-\theta t} \int_0^t e^{\theta s} B_{H,\lambda}^k(s) ds.$$

The process Y_t^k admits the Wiener integral representation $Y_t^k = \int_{\mathbb{R}} h^k(t-x) dW_x$ with kernel

$$h^k(z) = \left(e^{-\lambda z} z^{H-1/2} - a_k e^{-\theta z} \int_0^z e^{(\theta-\lambda)y} y^{H-1/2} dy \right) \mathbf{1}_{(0,\infty)}(z),$$

where $\mathbf{1}_{(0,\infty)}$ is the indicator function, $a_{\text{I}} = \theta$, and $a_{\text{II}} = \theta - \lambda$. Hence

$$\text{cov}(Y_t^k, Y_s^k) = r_k(t-s), \quad r_k(\tau) = \int_{\mathbb{R}} h^k(\tau+u) h^k(u) du.$$

Since Y^k is centered Gaussian and the covariance depends only on the time difference, the process Y^k is stationary.

For the autocovariance functions $r_k(\tau)$, $k \in \{\text{I}, \text{II}\}$, the following asymptotic relations are obtained as $\tau \rightarrow \infty$:

	$r_{\text{I}}(\tau)$	$r_{\text{II}}(\tau)$
$\theta < \lambda$	$\sim -\frac{\theta \Gamma(H + \frac{1}{2})^2}{2(\lambda^2 - \theta^2)^{H+\frac{1}{2}}} e^{-\theta \tau}$	$\sim \frac{\Gamma(H + \frac{1}{2})^2}{2\theta(\lambda^2 - \theta^2)^{H-\frac{1}{2}}} e^{-\theta \tau}$
$\theta = \lambda$	$\sim -\frac{\Gamma(H + \frac{1}{2})}{(H + \frac{1}{2}) 2^{H+\frac{3}{2}} \theta^{H-\frac{1}{2}}} e^{-\theta \tau} \tau^{H+\frac{1}{2}}$	$\sim \frac{\Gamma(H + \frac{1}{2})}{2^{H+\frac{1}{2}} \theta^{H+\frac{1}{2}}} e^{-\theta \tau} \tau^{H-\frac{1}{2}}$
$\theta > \lambda$	$\sim -\frac{\Gamma(H + \frac{1}{2})}{2^{H+\frac{1}{2}} \lambda^{H-\frac{3}{2}} (\theta^2 - \lambda^2)} e^{-\lambda \tau} \tau^{H-\frac{1}{2}}$	$\sim \frac{(H - \frac{1}{2}) \Gamma(H + \frac{1}{2})}{2^{H-\frac{1}{2}} \lambda^{H-\frac{1}{2}} (\theta - \lambda)(\theta + \lambda)} e^{-\lambda \tau} \tau^{H-\frac{3}{2}}, \quad H \neq \frac{1}{2}$

In the exceptional case $k = \text{II}$, $H = 1/2$, $\theta > \lambda$, one has the exact identity $r_{\text{II}}(\tau) = (2\theta)^{-1}e^{-\theta\tau}$, $\tau \geq 0$. Consequently, $r_k(\tau) \rightarrow 0$ and Y^k is ergodic.

For drift parameter estimation with known H and λ , we use an ergodic-type approach close to [2, 3]. Let $\psi_k(\theta) = \mathbb{E}_\theta(Y_0^k)^2$. The second moments have the spectral representations

$$\psi_{\text{I}}(\theta) = \frac{\Gamma(H + 1/2)^2}{\pi} \int_0^\infty \frac{\omega^2 d\omega}{(\theta^2 + \omega^2)(\lambda^2 + \omega^2)^{H+1/2}},$$

$$\psi_{\text{II}}(\theta) = \frac{\Gamma(H + 1/2)^2}{\pi} \int_0^\infty \frac{d\omega}{(\theta^2 + \omega^2)(\lambda^2 + \omega^2)^{H-1/2}}.$$

The functions ψ_k are continuous and strictly decreasing on $(0, \infty)$; hence θ is identifiable from the variance.

For X^k , the ergodic limits

$$\frac{1}{T} \int_0^T (X_t^k)^2 dt \rightarrow \psi_k(\theta), \quad \frac{1}{N} \sum_{j=1}^N (X_{jh}^k)^2 \rightarrow \psi_k(\theta)$$

hold almost surely as $T \rightarrow \infty$ and $N \rightarrow \infty$. This yields the estimators

$$\widehat{\theta}_T^k = \psi_k^{-1} \left(\frac{1}{T} \int_0^T (X_t^k)^2 dt \right), \quad \widehat{\theta}_{N,h}^k = \psi_k^{-1} \left(\frac{1}{N} \sum_{j=1}^N (X_{jh}^k)^2 \right).$$

They are strongly consistent and asymptotically normal:

$$\sqrt{T} (\widehat{\theta}_T^k - \theta) \xrightarrow{d} \mathcal{N} \left(0, \frac{\sigma_k^2}{(\psi_k'(\theta))^2} \right), \quad \sqrt{N} (\widehat{\theta}_{N,h}^k - \theta) \xrightarrow{d} \mathcal{N} \left(0, \frac{\sigma_{k,h}^2}{(\psi_k'(\theta))^2} \right),$$

where \xrightarrow{d} denotes convergence in distribution,

$$\sigma_k^2 = 2 \int_{\mathbb{R}} r_k(\tau)^2 d\tau, \quad \sigma_{k,h}^2 = 2 \sum_{m \in \mathbb{Z}} r_k(mh)^2.$$

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