Derivation and analysis of a Fokker-Planck equation describing a population of spiking resonate-and-fire neurons

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Based on Izhikhevich's resonate-and-fire model [1] of the temporal evolution of membrane potential, we derive a kinetic Fokker-Planck equation describing the probability density $\rho(x, v, t)$ of finding neurons in a population having a potential x , with a time derivative v , at a given time t, other models having already been studied in, for example [2], [3], or [4].

The equation reads

$$
\forall x \le u_F, \frac{\partial p(x,v,t)}{\partial t} + \nabla(p(x,v,t) - \overline{D}\nabla p(x,v,t)) = \delta_{u_R}(x) \otimes \delta_0(v)N(t),
$$

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$$
N(t) = \int_{v \in \mathbb{R}^+} v \cdot p(u_F, v, t) dv,
$$

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$$
\forall v, p(-\infty, v, t) = 0; \forall v < 0, p(u_F, v, t) = 0
$$

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$$
\forall x > u_F, \forall v, p(x, v, t) = 0, p^0(x, v) \ge 0, \int_{-\infty}^{u_F} \int_{-\infty}^{\infty} p^0(x, v) dx dv = 1
$$

with $\mu(x, v, t) = \begin{pmatrix} v \\ v^2w & v^2v + b^2 \end{pmatrix}$ $-v\rightarrow -\omega_0^2 x - v/\tau + bN(t)/\tau + b\nu_{ext}/\tau$ and $\overline{\overline{D}} = (a_0 + a_1N(t))\begin{bmatrix} 0 & 0 \\ 0 & 1/2 \end{bmatrix}$ 0 $1/\tau^2$ 1 $a_0, a_1, b, \nu, \omega_0, \tau$ being neurological constants.

Let us define the following open subset of \mathbb{R}^2 : $\Omega = (-\infty, u_F) \times \mathbb{R}$. Now let us first define some notions of solutions that will allow us to work on the problem:

Definition 1. A pair of functions (p, N) is said to be a strong solution of the system (P) on $[0; T[, T \in \mathbb{R}_+^* \cap \infty$ when :

- $p \in \mathcal{C}^0([-\infty, u_F] \times] \infty, \infty[\times) \cap \mathcal{C}^{2,2,1}(] \infty, u_R[\cup] u_R, u_F] \times] \infty, \infty[\times [0,T[) \ \cap$ $L^{\infty}([0,T[, L^1(]-\infty, u_F]\times]-\infty, \infty[))$ and $N \in \mathcal{C}^0([0,T[)$
- The functions p and N are solutions of (P) in the the classical sense on $] \infty$, $u_R[\cup]u_R$, $u_F[\times] - \infty$, $\infty[$ and in the sense of distributions on $]-\infty$; $u_F[\times] - \infty$, $\infty[$.

Definition 2. A pair of non-negative functions (p, N) with $p \in L^{\infty}(\mathbb{R}^+; L^2_+((-\infty, u_F) \times]$ (∞, ∞)), $N \in L^{1}_{loc,+}(\mathbb{R}^{+})$ is a weak solution of the problem (P) if for any test function $\phi((x,v),t) \in \mathcal{C}^{\infty}([- \infty; u_F] \times]-\infty, \infty[\times [0; T])$ such that $v \frac{\partial \phi}{\partial x}$, $(x+v) \frac{\partial \phi}{\partial v} \in L^{\infty}([- \infty; u_F[\times] \infty, \infty[\times]0; T[,$ we have

$$
\int_0^T \int_{x=-\infty}^{u_F} \int_{v=-\infty}^{\infty} p(x, v, t) \left[-\frac{\partial \phi}{\partial t} - \mu(x, v, t) \nabla \phi - \nabla (\overline{\overline{D}}(t) \nabla \phi) \right] dxdvdt =
$$
\n
$$
\int_0^T N(t) (\phi(u_R, 0, t) - \phi(u_F, 0, t)) dt + \int_{-\infty}^{u_F} \int_{-\infty}^{+\infty} p^0(x, v) \phi(x, v, 0) dxdv
$$
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$$
-\int_{-\infty}^{u_F} \int_{-\infty}^{+\infty} p(x, v, T) \phi(x, v, T) dxdv
$$

We could obtain the following theorem:

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Theorem 1. Let p^0 be in $L^2(\mathbb{R}^2 \times \mathbb{R}^+; \mathbb{R})$, and (p, N) a weak solution of our problem (P) , with p^0 verifying the following conditions, denoted (C) : $(x,v) \in \mathbb{R}^2$ $p^{0}(x, v) \, \mathrm{d}x \mathrm{d}v = 1; \; \forall (x, v) \in$

 \mathbb{R}^2 , $p^0(x,v) \geq 0$; $p^0 = 0$ a.e. on $\partial \Omega_{hyp}$ and p being possibly negative for $t > 0$ meanwhile $\forall t \geq 0, N(t) \geq 0$. Then : $\forall t \geq 0$, $\int p(x, v, t) dx dv = 1$ and $\forall t \geq 0$, $p(x, v, t) \geq 0$ a.e. in Ω $(x,v) \in \Omega$

Even though we could obtain some numerical approximations of the solutions, the classical toolbox associated with linear operators cannot be used because of the non-local linearity due to the function N, which makes the theoretical study complicated. Thus, to obtain more information, we then focus on a linearized version of the equation around 0, tackling the study of a boundary value Cauchy problem for an hypoelliptic operator, using different methods and approaches. In this linearized version, the function N disappears from the drift vector $\mu_{\mathbf{L}}$ and the diffusion matrix.

On top of the degeneracy of the problem, the difficulties associated with the right-hand side of the equation, which displays a non-local dependency on the solution as well as a measure, are of course of prime concern.

We will present steps towards existence and uniqueness in this difficult case by simplifying this right-hand side and imposing some boundedness on the spatial variable x , using the framework developped in [5] :

Theorem 2. If we consider $\Omega = x_{min}; x_{max}[\times \mathbb{R}, noting Q_T$ the classical parabolic cylinder associated with Ω , if p_0 is in $L^2(\Omega)$, the right-hand side f is in $L^2(Q_T)$ we have existence and uniqueness of a solution to the linearized problem for all initial conditions verifying (C) , the solution being in the space $Y = \{p \in H : \mathcal{T}p \in H'\}$, with $H = \{p \in L^2(Q_T) : \nabla_v p \in L^2(Q_T)\}\$ and $\mathcal{T} p = \partial_t p + \mu_L \cdot \nabla p$.

As well as another result, this time in the unbounded setting, displaying similarities but still with radically different methods and point of view in the line of the work in [6] and [7].

Theorem 3. For any boundary condition in $C_0(\partial\Omega) \times \mathbb{R}$, there exists an unique weak solution p verifying $p \in L^2_{v,t}(\Omega_v \times]0,T[, H^1_x(\Omega_x))$ and $\partial_t p + \mu_L.\nabla p \in L^2_{v,t}(\Omega_v \times]0,T[, H^{-1}_x(\Omega_x)).$

Some associated results (conservation of positivity, convergence of the transport trajectories in the Wasserstein metric) which allow for a better understanding of the problem, will also be presented.

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