

On the Fourth-Order Lattice Gel'fand–Dikii Equations

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Abstract. The fourth-order lattice Gel'fand–Dikii equations in quadrilateral form are investigated. Utilizing the direct linearization approach, we present some equations of the extended lattice Gel'fand–Dikii type. These equations are related to a quartic discrete dispersion relation and can be viewed as higher-order members of the extended lattice Boussinesq type equations. The resulting lattice equations given here are in five-component form, and some of them are multi-dimensionally consistent by introducing extra equations. Lax integrability is discussed both by direct linearization scheme and also through multi-dimensional consistent property. Some reductions of the five-component lattice equations to the four-component forms are considered.

Key words: lattice Gel'fand–Dikii type equation; direct linearization approach; multi-dimensional consistency; Lax pair

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1 Introduction

The Gel'fand–Dikii (GD) hierarchy [15, 16] is a generalization of the Korteweg–de Vries (KdV) hierarchy. The GD spectral problem is known as

$$L\psi = \lambda\psi, \quad L = \partial_x^n + u_2\partial_x^{n-2} + \cdots + u_{n-1}\partial_x + u_n,$$

where λ is a spectral parameter, $u_i = u_i(t, x)$ and $\partial_x = \frac{\partial}{\partial x}$. When $n = 2$ and $n = 3$, the above spectral problem gives rise to the KdV hierarchy and Boussinesq hierarchy, respectively. The semi-discrete (differential-difference) GD hierarchy is related to the spectral problem

$$L\psi(n) = \lambda\psi(n), \quad L = E^m + \sum_{j=1}^{m+k} u_j(n)E^{m-j}, \quad m, k \geq 1, \quad m, k \in \mathbb{Z},$$

where $u_j(n) = u_j(n, t)$ are functions defined on $\mathbb{Z} \times \mathbb{R}$, E is the shift operator defined as $E^j f(n) = f(n + j)$ for $j \in \mathbb{Z}$. The semi-discrete GD hierarchy has been systematically studied by Kupershmidt in [23]. Some explicit lower-order members can be found in [4]. Both continuous and semi-discrete GD hierarchies can be studied as reductions of the Kadomtsev–Petviashvili (KP) hierarchy related to a pseudo-differential operator or differential-difference KP hierarchy related to a pseudo-difference operator in the frame of Sato's KP theory.

Compared with continuous and semi-discrete cases, there is no pseudo-difference operator available for the fully discrete KP equations. Therefore, the discrete Boussinesq (DBSQ) type equations and higher-order GD type lattice equations are expected to play more roles in understanding the discrete KdV/KP type integrable systems. For the DBSQ equations, one can refer to [22, 39] and the references therein. As discrete models, the lattice GD hierarchy was first

introduced by Nijhoff et al. [29] via direct linearization (DL) approach and transforms to the GD differential hierarchy under continuum limits. This hierarchy has received a lot of attention in understanding the discrete KdV/KP type equations. In [24], a Lagrangian for the generic member of the lattice GD hierarchy was presented. This is a Lagrangian 2-form when embedded in a higher-dimensional lattice, obeying a closure relation. Subsequently, an integrable noncommutative modified lattice GD hierarchy was introduced [8]. The Lax integrability and multidimensional consistency (MDC) were shown. Very recently, a variational perspective on continuum limits for the lattice GD hierarchy was also discussed [37].

The simplest member of lattice GD hierarchy is the lattice (potential) KdV equation

$$(p - q + u_{n,m+1} - u_{n+1,m})(p + q + u_{n,m} - u_{n+1,m+1}) = p^2 - q^2, \quad (1.1)$$

where dependent variable $u_{n,m}$ is a function defined on the two-dimensional lattice with discrete coordinates $(n, m) \in \mathbb{Z}^2$, p and q denote lattice parameters associated with the directions in the lattice. The equation (1.1) first appeared as a nonlinear superposition formula of the Bäcklund transformation of the KdV equation [38]. This equation, together with the lattice modified KdV equation as well as the lattice Schwarzian KdV equation, composes the lattice KdV type equations. Analogously, the latter two equations were, respectively, obtained as nonlinear superposition formulas of the Bäcklund transformations for the modified KdV equation and the Schwarzian KdV equation. All these three equations can arise from the so-called Nijhoff–Quispel–Capel equation as distinct parameter choices [30]. Besides, all of them possess a multi-dimensionally consistent (MDC) property, which allows a lattice equation (or a system) to be consistently embedded into a higher dimension [1, 5, 27, 32]. With this property and three additional requirements on lattice equations: affine linear, D_4 symmetry, and tetrahedron property, Adler, Bobenko, and Suris (ABS) classified lattice models defined on an elementary quadrilateral [1]. Some of the lattice equations in the ABS list are the lattice KdV type equations mentioned above.

As a higher member in the lattice GD hierarchy, the potential DBSQ equation reads [29]

$$\begin{aligned} & \frac{p^3 - q^3}{p - q - v_{n+1,m+1} + v_{n,m+2}} - \frac{p^3 - q^3}{p - q - v_{n+2,m} + v_{n+1,m+1}} \\ &= (p - q - v_{n+1,m} + v_{n,m+1})(2p + q - v_{n+2,m+1} + v_{n,m}) \\ & \quad - (p - q - v_{n+2,m+1} + v_{n+1,m+2})(2p + q - v_{n+2,m+2} + v_{n,m+1}), \end{aligned} \quad (1.2)$$

which is defined on a 3×3 stencil on the discrete two-dimensional plane. Similar to the lattice KdV case, here dependent variable v is a function of discrete independent variables n and m , and p and q are lattice parameters. Together with the potential DBSQ equation (1.2), also lattice versions of the modified Boussinesq and Schwarzian Boussinesq equations have been found (see [26, 29, 39]). These equations constitute the DBSQ type equations. By introducing two additional variables, equation (1.2) can be expressed as a three-component form, which is defined at the vertices of an elementary square and possesses the MDC property [36]. Hietarinta [17] generalised this result and gave a classification of the three-component equations satisfying the MDC property, with some equations defined on the edges of the consistency cube and others on the faces of the cube. It was revealed that the equations in the classification followed one and the same underlying structure, and some of them can be viewed as extended DBSQ type equations [44]. We refer the reader to the recent survey [22] and the references therein for the DBSQ type equations and many interesting results.

The equations of each order in the extended lattice GD hierarchy are identified by the polynomial of symmetric form

$$G_N(\omega, k) := \sum_{j=1}^N \alpha_j (\omega^j - k^j), \quad \alpha_N \equiv 1, \quad (1.3)$$

with coefficients $\{\alpha_j\}$ and parameter k . It is worthy to note that parameter α_1 yields trivial extension for the whole extended lattice GD type hierarchies, while the parameters $\{\alpha_j\}_{j=2}^N$ lead to nontrivial extension (cf. [44]). Hence, $N = 2$ leads to the usual lattice KdV case and $N = 3$ yields the extended DBSQ case [44]. By “extended” we mean the parameters $\{\alpha_j\}$ are involved in the lattice equations. In addition, more details can be found in Section 2 on how (1.3) is used to define plane wave factors of the KdV case and DBSQ case.

This paper is devoted to lattice GD type equations arising from quartic polynomial $G_4(x, k)$, by using the DL approach. This method, based on the use of linear singular integral equations, was introduced by Fokas and Ablowitz in order to solve the KdV equation and treat the initial value problem of the Painlevé II equation [10]. Soon after the discrete DL scheme was developed [30, 33]. And then, this method has been widely used to investigate continuous, semi-discrete, and discrete integrable systems [11, 12, 13, 14, 28, 30, 33, 34, 41, 44]. In this paper, we are interested in the construction of the fourth-order lattice GD type equations as well as their MDC property and Lax integrability. For the sake of brevity, we call the resulting lattice equations as lattice GD-4 type equations. We will see that the lattice GD-4 type equations exhibit many nontrivial features compared with the DBSQ type equations.

The paper is arranged as follows: In Section 2, we give a brief review of the DL scheme for the extended lattice GD type equations. In Section 3, we construct the lattice GD-4 type equations as closed-forms, which are transformed into deformed forms in Section 4. Then, Section 5 is devoted to the MDC property and Lax representation, and Section 6 is for conclusions. There are three appendices, where we list the DBSQ equations, present details of some calculations and triply shifted quantities of the variables from MDC of some lattice GD-4 equations, respectively.

2 Framework of discrete direct linearisation approach

In this section, we briefly recall the DL framework for the extended lattice GD type equations outlined in the paper [44]. Let us introduce the following conventional short hands

$$f = f(n, m), \quad \tilde{f} = f(n + 1, m), \quad \hat{f} = f(n, m + 1), \quad \widehat{\tilde{f}} = f(n + 1, m + 1).$$

Besides, we employ $\omega_j(k)$, $j = 1, 2, \dots, N - 1$, and $\omega_N(k) = k$ to represent the roots of equation $G_N(\omega, k) = 0$ given in (1.3).

The starting point in the DL approach is the following $\infty \times \infty$ matrix \mathbf{C} defined by the integral

$$\mathbf{C} = \sum_{j=1}^N \oint_{\Gamma_j} d\lambda_j(k) \rho_k \mathbf{c}_k {}^t \mathbf{c}_{-\omega_j(k)} \sigma_{-\omega_j(k)}, \quad N \geq 2, \quad (2.1)$$

where $\{d\lambda_j(k)\}$ are certain measures for the integral on integration contours Γ_j in the space of the spectral variable k , $\mathbf{c}_k = (k^j)_{j \in \mathbb{Z}}$ stands for the column vector¹ $(\dots, k^{-2}, k^{-1}, 1, k, k^2, \dots)^T$, and ${}^t \mathbf{c}_{k'}$ denotes the transpose of $\mathbf{c}_{k'}$. The factors ρ_k , $\sigma_{k'}$ are discrete exponential functions of k and k' , respectively, given by

$$\rho_k = (p + k)^n (q + k)^m \rho_k^0, \quad \sigma_{k'} = (p - k')^{-n} (q - k')^{-m} \sigma_{k'}^0,$$

where $\rho_k^0, \sigma_{k'}^0 \in \mathbb{C}$. We assume that basic operations (such as differentiations with respect to parameters or applying shifts in the variables n and m) commute with the integrations. Note

¹We use “ t ” to denote the transpose of a matrix with infinite dimensions, e.g., ${}^t \mathbf{c}$, ${}^t \mathbf{A}$, and use “ T ” for the finite-dimensional case, e.g., equation (5.2).

that ρ_k and $\sigma_{-\omega_j(k)}$ can be considered to be defined using the roots of $G_N(-p, k) = 0$. When $N = 2$, we have

$$\rho = \rho_k \sigma_{-\omega_j(k)} = \left(\frac{p+k}{p-k} \right)^n \left(\frac{q+k}{q-k} \right)^m,$$

which leads to the lattice KdV type equations, cf. [33]. When $N = 3$, we have

$$\rho = \rho_k \sigma_{-\omega_j(k)} = \left(\frac{p+k}{p+\omega_j(k)} \right)^n \left(\frac{q+k}{q+\omega_j(k)} \right)^m,$$

for $j = 1, 2$, which leads to the extended DBSQ type equations, see more details in [44] for how these equations are derived in DL approach. Note also that by ‘‘extended’’ we mean in the lattice equations the parameters $\{\alpha_j\}$ are involved.

By introducing $\infty \times \infty$ matrices $\mathbf{\Lambda}$ and ${}^t\mathbf{\Lambda}$ defined by

$$\mathbf{\Lambda} \mathbf{c}_k = k \mathbf{c}_k, \quad {}^t\mathbf{c}_{k'} {}^t\mathbf{\Lambda} = k' {}^t\mathbf{c}_{k'},$$

one easily identify that matrix \mathbf{C} obeys the linear relations [44]

$$\tilde{\mathbf{C}}(p - {}^t\mathbf{\Lambda}) = (p + \mathbf{\Lambda})\mathbf{C}, \quad \hat{\mathbf{C}}(q - {}^t\mathbf{\Lambda}) = (q + \mathbf{\Lambda})\mathbf{C}, \quad (2.2a)$$

$$\left[\prod_{j=1}^N (\omega_j(-p) - \mathbf{\Lambda}) \right] \mathbf{C} = \mathbf{C} \prod_{j=1}^N (\omega_j(-p) + {}^t\mathbf{\Lambda}), \quad (2.2b)$$

$$\left[\prod_{j=1}^{N-1} (\omega_j(-p) - \mathbf{\Lambda}) \right] \tilde{\mathbf{C}} = \mathbf{C} \prod_{j=1}^{N-1} (\omega_j(-p) + {}^t\mathbf{\Lambda}), \quad (2.2c)$$

and similar relations with p replaced by q and $\tilde{\mathbf{C}}$ by $\hat{\mathbf{C}}$. To proceed, we introduce some auxiliary elements:

- a Cauchy kernel $\mathbf{\Omega}$ defined by the relation $\mathbf{\Omega}\mathbf{\Lambda} + {}^t\mathbf{\Lambda}\mathbf{\Omega} = \mathbf{O}$ with \mathbf{O} being a rank 1 projection matrix, obeying $\mathbf{O}^2 = \mathbf{O}$,
- an $\infty \times \infty$ matrix \mathbf{U} obeying the relation $\mathbf{U} = \mathbf{C} - \mathbf{U}\mathbf{\Omega}\mathbf{C}$,
- an $\infty \times 1$ vector \mathbf{u}_k defined by $\mathbf{u}_k = \rho_k(\mathbf{c}_k - \mathbf{U}\mathbf{\Omega}\mathbf{c}_k)$.

With these elements, from (2.2) one has

$$\tilde{\mathbf{U}}(p - {}^t\mathbf{\Lambda}) = (p + \mathbf{\Lambda})\mathbf{U} - \tilde{\mathbf{U}}\mathbf{O}\mathbf{U}, \quad (2.3a)$$

$$\begin{aligned} \mathbf{U} \left[\prod_{j=1}^{N-1} (\omega_j(-p) + {}^t\mathbf{\Lambda}) \right] &= \left[\prod_{j=1}^{N-1} (\omega_j(-p) - \mathbf{\Lambda}) \right] \tilde{\mathbf{U}} \\ &+ \mathbf{U} \sum_{j=0}^{N-2} \left[\prod_{l=1}^j (\omega_l(-p) + {}^t\mathbf{\Lambda}) \right] \mathbf{O} \left[\prod_{l=j+2}^{N-1} (\omega_l(-p) - \mathbf{\Lambda}) \right] \tilde{\mathbf{U}}, \end{aligned} \quad (2.3b)$$

$$\mathbf{U} \sum_{j=1}^N \alpha_j (-{}^t\mathbf{\Lambda})^j = \sum_{j=1}^N \alpha_j \mathbf{\Lambda}^j \mathbf{U} - \mathbf{U} \sum_{j=1}^N \alpha_j \sum_{l=0}^{j-1} (-{}^t\mathbf{\Lambda})^l \mathbf{O} \mathbf{\Lambda}^{j-1-l} \mathbf{U}, \quad (2.3c)$$

as well as

$$\tilde{\mathbf{u}}_k = (p + \mathbf{\Lambda})\mathbf{u}_k - \tilde{\mathbf{U}}\mathbf{O}\mathbf{u}_k, \quad (2.4a)$$

$$\begin{aligned}
& - \left[\prod_{j=1}^N (\omega_j(-p) - k) \right] \mathbf{u}_k = \left[\prod_{j=1}^{N-1} (\omega_j(-p) - \Lambda) \right] \tilde{\mathbf{u}}_k \\
& \quad + U \sum_{j=0}^{N-2} \left[\prod_{l=1}^j (\omega_l(-p) + {}^t\Lambda) \right] \mathcal{O} \left[\prod_{l=j+2}^{N-1} (\omega_l(-p) - \Lambda) \right] \tilde{\mathbf{u}}_k, \tag{2.4b}
\end{aligned}$$

$$\sum_{j=1}^N \alpha_j k^j \mathbf{u}_k = \sum_{j=1}^N \alpha_j \Lambda^j \mathbf{u}_k - U \sum_{j=1}^N \alpha_j \sum_{l=0}^{j-1} (-{}^t\Lambda)^l \mathcal{O} \Lambda^{j-1-l} \mathbf{u}_k. \tag{2.4c}$$

Similar relations to (2.3a), (2.3b), and (2.4a) and (2.4b) can be obtained with $(\hat{\cdot}, q)$ to replace $(\tilde{\cdot}, p)$.

Relations (2.3a) and (2.3b), together with their $(\hat{\cdot}, q)$ analogues, supply important information for the construction of lattice GD type equations. Note that although relation (2.3c) does not exert any role in the construction of discrete equations (e.g., [9, 43, 44]) and semi-discrete equations (e.g., [25]), it is indispensable in the construction of continuous equations (e.g., [40]). Relations (2.4) are linear with respect to \mathbf{u}_k and usually used, together with their $(\hat{\cdot}, q)$ analogues, to construct Lax pairs for the derived lattice equations (e.g., [44]).

3 Lattice GD-4 type equations as quadrilateral closed forms

In this section, we will restrict our attention to the relations (2.3) and (2.4) in the case of $N = 4$, and construct the lattice GD-4 type equations as quadrilateral closed forms. In this case, $\omega_j(k)$, $j = 1, 2, 3, 4$, in (2.3) and (2.4) are the roots of $G_4(\omega, k) = 0$, where $G_4(\omega, k)$ is the polynomial defined as in (1.3). In the following, we show step by step how quadrilateral closed forms of scalar functions arise from (2.3) and (2.4).

3.1 Basic objects and their relations

When $N = 4$ the set of relations (2.3) and (2.4) take the form

$$\tilde{U}(p - {}^t\Lambda) = (p + \Lambda)U - \tilde{U}OU, \tag{3.1a}$$

$$\begin{aligned}
& U(A_3(-p) + A_2(-p){}^t\Lambda + A_1(-p){}^t\Lambda^2 + {}^t\Lambda^3) \\
& \quad = (A_3(-p) - A_2(-p)\Lambda + A_1(-p)\Lambda^2 - \Lambda^3)\tilde{U} + U[A_2(-p)\mathcal{O} \\
& \quad \quad - A_1(-p)(\mathcal{O}\Lambda - {}^t\Lambda\mathcal{O}) + (\mathcal{O}\Lambda^2 - {}^t\Lambda\mathcal{O}\Lambda + {}^t\Lambda^2\mathcal{O})]\tilde{U}, \tag{3.1b}
\end{aligned}$$

$$U \sum_{j=1}^4 \alpha_j (-{}^t\Lambda)^j = \sum_{j=1}^4 \alpha_j \Lambda^j U - U \sum_{j=1}^4 \alpha_j \sum_{l=0}^{j-1} (-{}^t\Lambda)^l \mathcal{O} \Lambda^{j-1-l} U, \tag{3.1c}$$

and

$$\tilde{\mathbf{u}}_k = (p + \Lambda)\mathbf{u}_k - \tilde{U}\mathcal{O}\mathbf{u}_k, \tag{3.2a}$$

$$\begin{aligned}
& - \left[\prod_{j=1}^4 (\omega_j(-p) - k) \right] \mathbf{u}_k = (A_3(-p) - A_2(-p)\Lambda + A_1(-p)\Lambda^2 - \Lambda^3)\tilde{\mathbf{u}}_k \\
& \quad + U[A_2(-p)\mathcal{O} - A_1(-p)(\mathcal{O}\Lambda - {}^t\Lambda\mathcal{O}) + (\mathcal{O}\Lambda^2 - {}^t\Lambda\mathcal{O}\Lambda + {}^t\Lambda^2\mathcal{O})]\tilde{\mathbf{u}}_k, \tag{3.2b}
\end{aligned}$$

$$\sum_{j=1}^4 \alpha_j k^j \mathbf{u}_k = \sum_{j=1}^4 \alpha_j \Lambda^j \mathbf{u}_k - U \sum_{j=1}^4 \alpha_j \sum_{l=0}^{j-1} (-{}^t\Lambda)^l \mathcal{O} \Lambda^{j-1-l} \mathbf{u}_k, \tag{3.2c}$$

and these relations hold when $(\tilde{\cdot}, p)$ is replaced by $(\hat{\cdot}, q)$. Here $\{A_j(k)\}$ are defined as

$$\begin{aligned} A_1(k) &= \omega_1(k) + \omega_2(k) + \omega_3(k) = -(k + \alpha_3), \\ A_2(k) &= \omega_1(k)\omega_2(k) + \omega_1(k)\omega_3(k) + \omega_2(k)\omega_3(k) = k^2 + \alpha_3k + \alpha_2, \\ A_3(k) &= \omega_1(k)\omega_2(k)\omega_3(k) = -(k^3 + \alpha_3k^2 + \alpha_2k + \alpha_1), \end{aligned}$$

and $\{\omega_j(k)\}$ are the roots of $G_4(\omega, k) = 0$, i.e.,

$$G_4(\omega, k) = \prod_{j=1}^4 (\omega - \omega_j(k)), \quad \omega_4(k) = k.$$

Next, in order to construct the lattice equations, let us introduce the following scalar objects:

$$v_a := 1 - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}\mathbf{e}, \quad w_b := 1 + {}^t\mathbf{e}\mathbf{U}(b - {}^t\mathbf{\Lambda})^{-1}\mathbf{e}, \quad (3.3a)$$

$$s_a := a - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}{}^t\mathbf{\Lambda}\mathbf{e}, \quad t_b := -b + {}^t\mathbf{e}\mathbf{\Lambda}\mathbf{U}(b - {}^t\mathbf{\Lambda})^{-1}\mathbf{e}, \quad (3.3b)$$

$$r_a := a^2 - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}{}^t\mathbf{\Lambda}^2\mathbf{e}, \quad z_b := b^2 + {}^t\mathbf{e}\mathbf{\Lambda}^2\mathbf{U}(b - {}^t\mathbf{\Lambda})^{-1}\mathbf{e}, \quad (3.3c)$$

$$f_a := a^3 - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}{}^t\mathbf{\Lambda}^3\mathbf{e}, \quad g_b := -b^3 + {}^t\mathbf{e}\mathbf{\Lambda}^3\mathbf{U}(b - {}^t\mathbf{\Lambda})^{-1}\mathbf{e}, \quad (3.3d)$$

$$u_{i,j} := {}^t\mathbf{e}\mathbf{\Lambda}^i\mathbf{U}{}^t\mathbf{\Lambda}^j\mathbf{e}, \quad s_{a,b} := {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e}, \quad i, j \in \mathbb{Z}. \quad (3.3e)$$

Here a and b are parameters, \mathbf{e} is a fixed column vector $(\dots, 0, 0, 0, 1, 0, 0, \dots)^T$ where the only nonzero entry 1 is in the central, the adjoint vector ${}^t\mathbf{e}$ is the transpose of \mathbf{e} . It then turns out that

$${}^t\mathbf{e}\mathbf{c}_k = 1, \quad {}^t\mathbf{c}_{k'}\mathbf{e} = 1, \quad \mathbf{O} = \mathbf{e}{}^t\mathbf{e}.$$

Except f_a and g_b , all others take the same forms as in [44] for the DBSQ case. Based on the shift relations (3.1a) and (3.1b), for the objects (3.3) we have the following relations: multiplying (3.1a) and (3.1b) from the left by ${}^t\mathbf{e}$ and from the right by \mathbf{e} , we have

$$p\tilde{u}_{i,j} - \tilde{u}_{i,j+1} = pu_{i,j} + u_{i+1,j} - \tilde{u}_{i,0}u_{0,j}, \quad (3.4a)$$

$$\begin{aligned} &A_3(-p)u_{i,j} + A_2(-p)u_{i,j+1} + A_1(-p)u_{i,j+2} + u_{i,j+3} \\ &= A_3(-p)\tilde{u}_{i,j} - A_2(-p)\tilde{u}_{i+1,j} + A_1(-p)\tilde{u}_{i+2,j} - \tilde{u}_{i+3,j} \\ &\quad + A_2(-p)u_{i,0}\tilde{u}_{0,j} + A_1(-p)(u_{i,1}\tilde{u}_{0,j} - u_{i,0}\tilde{u}_{1,j}) + u_{i,0}\tilde{u}_{2,j} - u_{i,1}\tilde{u}_{1,j} + u_{i,2}\tilde{u}_{0,j}, \end{aligned} \quad (3.4b)$$

where p_a, p_b are defined as

$$p_a = G_4(-p, -a)/(p - a) = A_3(-p) + aA_2(-p) + a^2A_1(-p) + a^3, \quad (3.5a)$$

$$p_b = G_4(-p, -b)/(p - b) = A_3(-p) + bA_2(-p) + b^2A_1(-p) + b^3; \quad (3.5b)$$

we can also derive (as examples of derivation, we provide details in Appendix B for equations (3.6b), (3.9a) and (3.10a))

$$1 + (p - a)s_{a,b} - (p - b)\tilde{s}_{a,b} = \tilde{v}_a w_b, \quad (3.6a)$$

$$\begin{aligned} &[A_2(-p) + (a + b)A_1(-p) + a^2 + ab + b^2] + p_b s_{a,b} - p_a \tilde{s}_{a,b} \\ &= v_a \tilde{z}_b - s_a \tilde{t}_b + r_a \tilde{w}_b + A_2(-p)v_a \tilde{w}_b - A_1(-p)(v_a \tilde{t}_b - s_a \tilde{w}_b), \end{aligned} \quad (3.6b)$$

and

$$\tilde{s}_a = (p + u_0)\tilde{v}_a - (p - a)v_a, \quad (3.7a)$$

$$t_b = (p - b)\tilde{w}_b - (p - \tilde{u}_0)w_b, \quad (3.7b)$$

and

$$\tilde{r}_a = p\tilde{s}_a - (p - a)s_a + \tilde{v}_a u_{0,1}, \quad (3.8a)$$

$$z_b = (p - b)\tilde{t}_b - pt_b + \tilde{u}_{1,0}w_b, \quad (3.8b)$$

and

$$\tilde{f}_a = p\tilde{r}_a - (p - a)r_a + \tilde{v}_a u_{0,2}, \quad (3.9a)$$

$$g_b = (p - b)\tilde{z}_b - pz_b + \tilde{u}_{2,0}w_b, \quad (3.9b)$$

and

$$\begin{aligned} f_a &= p_a\tilde{v}_a - [A_3(-p) - A_2(-p)\tilde{u}_0 + A_1(-p)\tilde{u}_{1,0} - \tilde{u}_{2,0}]v_a \\ &\quad - [A_2(-p) - A_1(-p)\tilde{u}_0 + \tilde{u}_{1,0}]s_a - (A_1(-p) - \tilde{u}_0)r_a, \end{aligned} \quad (3.10a)$$

$$\begin{aligned} \tilde{g}_b &= -p_b w_b + [A_3(-p) + A_2(-p)u_0 + A_1(-p)u_{0,1} + u_{0,2}]\tilde{w}_b \\ &\quad - [A_2(-p) + A_1(-p)u_0 + u_{0,1}]\tilde{t}_b + (A_1(-p) + u_0)\tilde{z}_b, \end{aligned} \quad (3.10b)$$

where we have used a simplified notation $u_0 = u_{0,0}$. All these relations, (3.6)–(3.10), also hold for their hat- q counterparts obtained by replacing (\sim, p) with $(\hat{\sim}, q)$. It should be emphasized that q_a and q_b appear in the rest part of the paper have the same form with p_a and p_b just with replacement $p \rightarrow q$.

A further set of relations is

$$\begin{aligned} &[H(p, q) + \hat{u}_{1,0} - (p + q - \alpha_3)\hat{u}_0 + ((p + q - \alpha_3 - \hat{u}_0)s_a + r_a)/v_a](p - q + \hat{u}_0 - \tilde{u}_0) \\ &= (p_a\tilde{v}_a - q_a\hat{v}_a)/v_a, \end{aligned} \quad (3.11a)$$

$$p - q + \hat{u}_0 - \tilde{u}_0 = [(p - a)\hat{v}_a - (q - a)\tilde{v}_a]/\hat{v}_a, \quad (3.11b)$$

$$(\tilde{r}_a - p\tilde{s}_a)/\tilde{v}_a - (\hat{r}_a - q\hat{s}_a)/\hat{v}_a = -(p - a)s_a/\tilde{v}_a + (q - a)s_a/\hat{v}_a, \quad (3.11c)$$

which are derived from equations (3.4a), (3.7a), (3.8a), (3.10a) and their hat- q counterparts, where

$$H(p, q) = (A_3(-p) - A_3(-q))/(p - q).$$

Similarly, from (3.4a), (3.7b), (3.8b), (3.10b) and their hat- q versions one has

$$\begin{aligned} &[H(p, q) + u_{0,1} + (p + q - \alpha_3)u_0 - ((p + q - \alpha_3 + u_0)\hat{t}_b - \hat{z}_b)/\hat{w}_b] \\ &\quad \times (p - q + \hat{u}_0 - \tilde{u}_0) = (p_b\hat{w}_b - q_b\tilde{w}_b)/\hat{w}_b, \end{aligned} \quad (3.12a)$$

$$p - q + \hat{u}_0 - \tilde{u}_0 = [(p - b)\tilde{w}_b - (q - b)\hat{w}_b]/w_b, \quad (3.12b)$$

$$(\hat{z}_b + p\hat{t}_b)/\hat{w}_b - (\tilde{z}_b + q\tilde{t}_b)/\tilde{w}_b = (p - b)\hat{t}_b/\hat{w}_b - (q - b)\tilde{t}_b/\tilde{w}_b. \quad (3.12c)$$

3.2 Lattice equations as quadrilateral closed forms

We have achieved many relations with respect to the scalar elements defined in equation (3.3). With these relations in hand, we are able to derive quadrilateral closed forms of some elements, which can be considered as lattice equations.

We start by presenting a set of lattice equations, which are made up of variables v_a , s_a , r_a , u_0 and $u_{1,0}$. To do that, we first set $i = j = 0$ in equation (3.4a) and its hat- q counterpart, which yield

$$p\tilde{u}_0 - \tilde{u}_{0,1} = pu_0 + u_{1,0} - \tilde{u}_0u_0, \quad q\hat{u}_0 - \hat{u}_{0,1} = qu_0 + u_{1,0} - \hat{u}_0u_0. \quad (3.13)$$

It then follows that

$$\widehat{u}_{1,0} - \widetilde{u}_{1,0} = (p - q + \widehat{u}_0 - \widetilde{u}_0)\widehat{u}_0 - p\widehat{u}_0 + q\widetilde{u}_0, \quad (3.14a)$$

$$\widehat{u}_{0,1} - \widetilde{u}_{0,1} = (p - q + \widehat{u}_0 - \widetilde{u}_0)u_0 - p\widetilde{u}_0 + q\widehat{u}_0. \quad (3.14b)$$

Obviously, equation (3.7a) and its hat- q counterpart, (3.11a), (3.11c) and (3.14a) give rise to a closed-form with respect to v_a , s_a , r_a , u_0 and $u_{1,0}$, i.e.,

$$\widetilde{s}_a = (p + u_0)\widetilde{v}_a - (p - a)v_a, \quad \widehat{s}_a = (q + u_0)\widehat{v}_a - (q - a)v_a, \quad (3.15a)$$

$$(\widetilde{r}_a - p\widetilde{s}_a)/\widetilde{v}_a - (\widehat{r}_a - q\widehat{s}_a)/\widehat{v}_a = -(p - a)s_a/\widetilde{v}_a + (q - a)s_a/\widehat{v}_a, \quad (3.15b)$$

$$\widehat{u}_{1,0} - \widetilde{u}_{1,0} = (p - q + \widehat{u}_0 - \widetilde{u}_0)\widehat{u}_0 - p\widehat{u}_0 + q\widetilde{u}_0, \quad (3.15c)$$

$$\begin{aligned} & [H(p, q) + \widehat{u}_{1,0} - (p + q - \alpha_3)\widehat{u}_0 + ((p + q - \alpha_3 - \widehat{u}_0)s_a + r_a)/v_a] \\ & \times (p - q + \widehat{u}_0 - \widetilde{u}_0) = (p_a\widetilde{v}_a - q_a\widehat{v}_a)/v_a. \end{aligned} \quad (3.15d)$$

This is a set of equations consisting of five equations of five scalar functions. We will revisit it in Section 4. A similar closed-form is composed by (3.7b), its hat- q counterpart, (3.12), (3.12c) and (3.14b), i.e.,

$$t_b = (p - b)\widetilde{w}_b - (p - \widetilde{u}_0)w_b, \quad \widehat{t}_b = (q - b)\widehat{w}_b - (q - \widehat{u}_0)w_b, \quad (3.16a)$$

$$(\widehat{z}_b + p\widehat{t}_b)/\widehat{w}_b - (\widetilde{z}_b + q\widetilde{t}_b)/\widetilde{w}_b = (p - b)\widehat{t}_b/\widehat{w}_b - (q - b)\widetilde{t}_b/\widetilde{w}_b, \quad (3.16b)$$

$$\widehat{u}_{0,1} - \widetilde{u}_{0,1} = (p - q + \widehat{u}_0 - \widetilde{u}_0)u_0 - p\widetilde{u}_0 + q\widehat{u}_0, \quad (3.16c)$$

$$\begin{aligned} & [H(p, q) + u_{0,1} + (p + q - \alpha_3)u_0 - ((p + q - \alpha_3 + u_0)\widehat{t}_b - \widetilde{z}_b)/\widehat{w}_b] \\ & \times (p - q + \widehat{u}_0 - \widetilde{u}_0) = (p_b\widehat{w}_b - q_b\widetilde{w}_b)/\widehat{w}_b, \end{aligned} \quad (3.16d)$$

which are equations of t_b , w_b , z_b , $u_{0,1}$ and u_0 .

There are more closed forms. Taking $i = j = 0$ in (3.4b) and its hat- q counterpart yields

$$\begin{aligned} & A_3(-p)u_0 + A_2(-p)u_{0,1} + A_1(-p)u_{0,2} + u_{0,3} \\ & = A_3(-p)\widetilde{u}_0 - A_2(-p)\widetilde{u}_{1,0} + A_1(-p)\widetilde{u}_{2,0} - \widetilde{u}_{3,0} \\ & \quad + A_2(-p)u_0\widetilde{u}_0 + A_1(-p)(u_{0,1}\widetilde{u}_0 - u_0\widetilde{u}_{1,0}) + u_0\widetilde{u}_{2,0} - u_{0,1}\widetilde{u}_{1,0} + u_{0,2}\widetilde{u}_0, \end{aligned} \quad (3.17a)$$

$$\begin{aligned} & A_3(-q)u_0 + A_2(-q)u_{0,1} + A_1(-q)u_{0,2} + u_{0,3} \\ & = A_3(-q)\widehat{u}_0 - A_2(-q)\widehat{u}_{1,0} + A_1(-q)\widehat{u}_{2,0} - \widehat{u}_{3,0} \\ & \quad + A_2(-q)u_0\widehat{u}_0 + A_1(-q)(u_{0,1}\widehat{u}_0 - u_0\widehat{u}_{1,0}) + u_0\widehat{u}_{2,0} - u_{0,1}\widehat{u}_{1,0} + u_{0,2}\widehat{u}_0. \end{aligned} \quad (3.17b)$$

We further set $(i, j) = (1, 0)$ and $(0, 1)$ in (3.4a) and obtain

$$p\widetilde{u}_{1,0} - \widetilde{u}_{1,1} = pu_{1,0} + u_{2,0} - \widetilde{u}_{1,0}u_0, \quad q\widehat{u}_{1,0} - \widehat{u}_{1,1} = qu_{1,0} + u_{2,0} - \widehat{u}_{1,0}u_0, \quad (3.18a)$$

$$p\widetilde{u}_{0,1} - \widetilde{u}_{0,2} = pu_{0,1} + u_{1,1} - \widetilde{u}_0u_{0,1}, \quad q\widehat{u}_{0,1} - \widehat{u}_{0,2} = qu_{0,1} + u_{1,1} - \widehat{u}_0u_{0,1}. \quad (3.18b)$$

Eliminating $u_{1,1}$ in (3.18a) and (3.18b), respectively, we reach to

$$\widehat{u}_{2,0} - \widetilde{u}_{2,0} = -p\widehat{u}_{1,0} + q\widetilde{u}_{1,0} + (p - q + \widehat{u}_0 - \widetilde{u}_0)\widehat{u}_{1,0}, \quad (3.19a)$$

$$\widehat{u}_{0,2} - \widetilde{u}_{0,2} = -p\widetilde{u}_{0,1} + q\widehat{u}_{0,1} + (p - q + \widehat{u}_0 - \widetilde{u}_0)u_{0,1}. \quad (3.19b)$$

Next, subtracting (3.17b) from (3.17a) yields an equation without $u_{0,3}$ but containing $u_{3,0}$. To delete $u_{3,0}$, we make use of equations (3.19a) and (3.4a) with $(i, j) = (2, 0)$ and $(1, 0)$. Then, after some algebra we arrive at

$$\begin{aligned} & \frac{G_4(-p, -q)}{p - q + \widehat{u}_0 - \widetilde{u}_0} = \frac{G_4(-p, -q)}{p - q} + H(p, q)(u_0 - \widehat{u}_0) \\ & \quad - (p + q - \alpha_3)(u_0\widehat{u}_0 - \widehat{u}_{1,0} + u_{0,1}) + u_0\widehat{u}_{1,0} - \widehat{u}_0u_{0,1} - \widehat{u}_{2,0} + u_{0,2}. \end{aligned} \quad (3.20)$$

Equation (3.13) and (3.19) together with (3.20) compose a new closed-form with respect to u_0 , $u_{0,1}$, $u_{1,0}$, $u_{0,2}$ and $u_{2,0}$, i.e.,

$$p\tilde{u}_0 - \tilde{u}_{0,1} = pu_0 + u_{1,0} - \tilde{u}_0u_0, \quad q\hat{u}_0 - \hat{u}_{0,1} = qu_0 + u_{1,0} - \hat{u}_0u_0, \quad (3.21a)$$

$$\hat{u}_{2,0} - \tilde{u}_{2,0} = -p\hat{u}_{1,0} + q\tilde{u}_{1,0} + (p - q + \hat{u}_0 - \tilde{u}_0)\hat{u}_{1,0}, \quad (3.21b)$$

$$\hat{u}_{0,2} - \tilde{u}_{0,2} = -p\tilde{u}_{0,1} + q\hat{u}_{0,1} + (p - q + \hat{u}_0 - \tilde{u}_0)u_{0,1}, \quad (3.21c)$$

$$\begin{aligned} \frac{G_4(-p, -q)}{p - q + \hat{u}_0 - \tilde{u}_0} &= \frac{G_4(-p, -q)}{p - q} + H(p, q)(u_0 - \hat{u}_0) \\ &\quad - (p + q - \alpha_3)(u_0\hat{u}_0 - \hat{u}_{1,0} + u_{0,1}) + u_0\hat{u}_{1,0} - \hat{u}_0u_{0,1} - \hat{u}_{2,0} + u_{0,2}. \end{aligned} \quad (3.21d)$$

Finally, we look for a closed-form containing the element $s_{a,b}$. By transformation

$$s_{a,b} = S_{a,b} - 1/(b - a), \quad a \neq b,$$

the relation (3.6) and its hat- q counterpart are rewritten as

$$(p - a)S_{a,b} - (p - b)\tilde{S}_{a,b} = \tilde{v}_aw_b, \quad (q - a)S_{a,b} - (q - b)\hat{S}_{a,b} = \hat{v}_aw_b, \quad (3.22a)$$

$$pbS_{a,b} - p_a\tilde{S}_{a,b} = A_2(-p)v_a\tilde{w}_b - A_1(-p)(v_a\tilde{t}_b - s_a\tilde{w}_b) + v_a\tilde{z}_b - s_a\tilde{t}_b + r_a\tilde{w}_b, \quad (3.22b)$$

$$qbS_{a,b} - q_a\hat{S}_{a,b} = A_2(-q)v_a\hat{w}_b - A_1(-q)(v_a\hat{t}_b - s_a\hat{w}_b) + v_a\hat{z}_b - s_a\hat{t}_b + r_a\hat{w}_b. \quad (3.22c)$$

Deleting \hat{z}_b in (3.22b) (the hat-shifted (3.22b)) by utilizing (3.12) and deleting \hat{s}_a by using the hat- q version of (3.7a) and also making use of (3.12b) we get

$$\begin{aligned} (q - a)(v_a\hat{t}_b - s_a\hat{w}_b) &= \frac{-\hat{v}_aw_b(pb\hat{w}_b - qb\tilde{w}_b)}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b} \\ &\quad + (q - a)(p + q - \alpha_3)v_a\hat{w}_b + pb\hat{S}_{a,b} - p_a\hat{S}_{a,b}. \end{aligned} \quad (3.23)$$

Making use of (3.22a) again, equation (3.23) goes to the following desired form

$$v_a\hat{t}_b - s_a\hat{w}_b = -w_b \frac{\frac{pb}{p-a}\hat{w}_b\tilde{v}_a - \frac{qb}{q-a}\tilde{w}_b\hat{v}_a}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b} + (p + q - \alpha_3)v_a\hat{w}_b + \frac{G_4(-a, -b)}{(p - a)(q - a)}\hat{S}_{a,b}. \quad (3.24)$$

This procedure is not apparent. Let add some details. First, we rewrite (3.23) as

$$\begin{aligned} v_a\hat{t}_b - s_a\hat{w}_b &= \frac{\frac{qb\hat{v}_a\tilde{w}_b}{q-a} - \frac{pb\tilde{v}_a\hat{w}_b}{p-a}}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b}w_b + \frac{\frac{pb\tilde{v}_a\hat{w}_b}{p-a} - \frac{pb\hat{v}_a\tilde{w}_b}{q-a}}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b}w_b \\ &\quad + (p + q - \alpha_3)v_a\hat{w}_b - \frac{p_a}{q - a}\hat{S}_{a,b} + \frac{pb}{q - a}\hat{S}_{a,b}. \end{aligned} \quad (3.25)$$

Then, noticing that

$$\frac{\frac{pb\tilde{v}_a\hat{w}_b}{p-a} - \frac{pb\hat{v}_a\tilde{w}_b}{q-a}}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b}w_b = \frac{pb\hat{w}_b \frac{1}{(p-a)(q-a)} \left[-(p - a)\hat{v}_a + (q - a)\tilde{v}_a \right]}{(p - b)\frac{\tilde{w}_b}{w_b} - (q - b)\frac{\hat{w}_b}{w_b}},$$

and from (3.11b) and (3.12b) we can express v_a terms using w_b and its shifts, and then we get

$$\frac{\frac{pb\tilde{v}_a\hat{w}_b}{p-a} - \frac{pb\hat{v}_a\tilde{w}_b}{q-a}}{(p - b)\tilde{w}_b - (q - b)\hat{w}_b}w_b = -\frac{pb}{(p - a)(q - a)}\hat{w}_b\hat{v}_a.$$

The right-hand side can be expressed as the hat shift of the first part of (3.22a), i.e.,

$$(p - a)\hat{S}_{a,b} - (p - b)\tilde{S}_{a,b} = \hat{v}_a\hat{w}_b.$$

Thus, we have

$$\frac{\frac{p_b \tilde{v}_a \tilde{w}_b}{p-a} - \frac{p_b \hat{v}_a \hat{w}_b}{q-a}}{(p-b)\tilde{w}_b - (q-b)\hat{w}_b} w_b = -\frac{p_b}{q-a} \hat{S}_{a,b} + \frac{p_b(p-b)}{(p-a)(q-a)} \hat{\tilde{S}}_{a,b}. \quad (3.26)$$

Then using the relations (3.25), (3.26) and making some algebras, we get

$$\begin{aligned} v_a \hat{t}_b - s_a \hat{\tilde{w}}_b &= \frac{\frac{q_b \hat{v}_a \tilde{w}_b}{q-a} - \frac{p_b \tilde{v}_a \hat{w}_b}{p-a}}{(p-b)\tilde{w}_b - (q-b)\hat{w}_b} w_b + (p+q-\alpha_3) v_a \hat{\tilde{w}}_b \\ &\quad + \hat{\tilde{S}}_{a,b} \left(\frac{-p_a}{q-a} + \frac{p_b(p-b)}{(q-a)(p-a)} \right). \end{aligned} \quad (3.27)$$

Since

$$p_a = \frac{G_4(-p, -a)}{p-a}, \quad p_b = \frac{G_4(-p, -b)}{p-b},$$

and

$$\frac{-p_a}{q-a} + \frac{p_b(p-b)}{(q-a)(p-a)} = \frac{1}{(p-a)(q-a)} (G_4(-a, -b)),$$

(3.27) gives rise to (3.24).

In view of the equations (3.7a) and (3.7b) and their hat- q versions, we delete variable u_0 and obtain

$$\tilde{s}_a/\tilde{v}_a - \hat{s}_a/\hat{v}_a = p - q - v_a((p-a)/\tilde{v}_a - (q-a)/\hat{v}_a), \quad (3.28a)$$

$$\tilde{t}_b/\tilde{w}_b - \hat{t}_b/\hat{w}_b = p - q - \hat{w}_b((p-b)/\hat{w}_b - (q-b)/\tilde{w}_b). \quad (3.28b)$$

The system composed by (3.22a), (3.24) and (3.28) can be viewed as a closed-form for elements s_a, v_a, t_b, w_b and $S_{a,b}$, i.e.,

$$(p-a)S_{a,b} - (p-b)\tilde{S}_{a,b} = \tilde{v}_a w_b, \quad (q-a)S_{a,b} - (q-b)\hat{S}_{a,b} = \hat{v}_a w_b, \quad (3.29a)$$

$$\tilde{s}_a/\tilde{v}_a - \hat{s}_a/\hat{v}_a = p - q - v_a((p-a)/\tilde{v}_a - (q-a)/\hat{v}_a), \quad (3.29b)$$

$$\tilde{t}_b/\tilde{w}_b - \hat{t}_b/\hat{w}_b = p - q - \hat{w}_b((p-b)/\hat{w}_b - (q-b)/\tilde{w}_b), \quad (3.29c)$$

$$\begin{aligned} v_a \hat{t}_b - s_a \hat{\tilde{w}}_b &= -w_b \frac{\frac{p_b}{p-a} \hat{w}_b \tilde{v}_a - \frac{q_b}{q-a} \tilde{w}_b \hat{v}_a}{(p-b)\tilde{w}_b - (q-b)\hat{w}_b} \\ &\quad + (p+q-\alpha_3) v_a \hat{\tilde{w}}_b + \frac{G_4(-a, -b)}{(p-a)(q-a)} \hat{\tilde{S}}_{a,b}. \end{aligned} \quad (3.29d)$$

We note that the equation (3.29d) can alternatively be replaced by the following equation

$$v_a \hat{t}_b - s_a \hat{\tilde{w}}_b = -w_b \frac{\frac{p_a}{p-b} \hat{w}_b \tilde{v}_a - \frac{q_a}{q-b} \tilde{w}_b \hat{v}_a}{(p-b)\tilde{w}_b - (q-b)\hat{w}_b} + (p+q-\alpha_3) v_a \hat{\tilde{w}}_b + \frac{G_4(-a, -b)}{(p-b)(q-b)} S_{a,b}, \quad (3.30)$$

which can be derived in a similar way to (3.24). Therefore, we may have another closed form composed by (3.29a), (3.29b), (3.29c) and (3.30).

4 Lattice GD-4 type equations

We have got quadrilateral closed forms (3.15), (3.16), (3.21) and (3.29). All the elements, v_a, s_a, r_a, u_0 , etc., are well defined through (3.3), which means these closed forms (if they are considered

as equations) allow multi-soliton solutions.² In an ideal multicomponent quadrilateral equation, all dependent variables should evolve under certain boundary conditions. In this section, first, we will revisit these obtained closed forms. We will try introducing extra equations so that these closed forms (3.15), (3.16) and (3.21) can well evolve for some given boundary conditions, and therefore they can be considered in a more general sense as lattice equations of the GD-4 type. After that, we will provide deformations of these equations. Evolutions of these deformed lattice equations will be examined in Section 4.4. We will see that these lattice equations (see (4.8), (4.10), (4.12)) are the counterparts of the Hietarinta's (A-2), (B-2) and (C-3) equations of the DBSQ type [17, 44].

4.1 Closed forms: revisited

The purpose of this part is to look for extra equations to compensate for the previously obtained closed forms so that the new systems can evolve under various boundary conditions.

In the closed form (3.15), $\widehat{\tilde{r}}_a$ and $u_{0,1}$ are missed. In the next we look for an equation that involves with $\widehat{\tilde{r}}_a$ and $u_{0,1}$. For this sake, take $\widehat{\sim}$ -shift of (3.8a) and using (3.13) to delete $\widehat{u}_{0,1}$, we get

$$\widehat{\tilde{r}}_a = p\widehat{\tilde{s}}_a - (p-a)\widehat{s}_a + \widehat{v}_a(q\widehat{u}_0 - qu_0 - u_{1,0} + \widehat{u}_0u_0).$$

Replacing the shift $\widehat{\sim}$ by the shift $\widehat{\sim}$ whilst replacing the parameter p by q , we have another equation on $\widehat{\tilde{r}}_a$:

$$\widehat{\tilde{r}}_a = q\widehat{\tilde{s}}_a - (q-a)\widehat{s}_a + \widehat{v}_a(p\widehat{u}_0 - pu_0 - u_{1,0} + \widehat{u}_0u_0).$$

Adding them together yields

$$\begin{aligned} 2\widehat{\tilde{r}}_a &= (p+q)\widehat{\tilde{s}}_a - (p-a)\widehat{s}_a - (q-a)\widehat{s}_a \\ &\quad + \widehat{v}_a(p\widehat{u}_0 + q\widehat{u}_0 - (p+q - \widehat{u}_0 - \widehat{u}_0)u_0 - 2u_{1,0}). \end{aligned}$$

This equation together with the closed form (3.15) can make up a new set of lattice equations, i.e.,

$$\widehat{s}_a = (p+u_0)\widehat{v}_a - (p-a)v_a, \quad \widehat{s}_a = (q+u_0)\widehat{v}_a - (q-a)v_a, \quad (4.2a)$$

$$(\widehat{\tilde{r}}_a - p\widehat{\tilde{s}}_a)/\widehat{v}_a - (\widehat{r}_a - q\widehat{s}_a)/\widehat{v}_a = -(p-a)s_a/\widehat{v}_a + (q-a)s_a/\widehat{v}_a, \quad (4.2b)$$

$$\begin{aligned} 2\widehat{\tilde{r}}_a &= (p+q)\widehat{\tilde{s}}_a - (p-a)\widehat{s}_a - (q-a)\widehat{s}_a \\ &\quad + \widehat{v}_a(p\widehat{u}_0 + q\widehat{u}_0 - (p+q - \widehat{u}_0 - \widehat{u}_0)u_0 - 2u_{1,0}), \end{aligned} \quad (4.2c)$$

$$\widehat{u}_{1,0} - \widehat{u}_{1,0} = (p-q + \widehat{u}_0 - \widehat{u}_0)\widehat{u}_0 - p\widehat{u}_0 + q\widehat{u}_0, \quad (4.2d)$$

$$\begin{aligned} [H(p,q) + \widehat{u}_{1,0} - (p+q - \alpha_3)\widehat{u}_0 + ((p+q - \alpha_3 - \widehat{u}_0)s_a + r_a)/v_a] \\ \times (p-q + \widehat{u}_0 - \widehat{u}_0) = (p_a\widehat{v}_a - q_a\widehat{v}_a)/v_a, \end{aligned} \quad (4.2e)$$

which evolves when initial values are given on the staircase (see Section 4.4 and Figure 2). In a quite similar fashion, one can elaborate (3.16), and we get

$$t_b = (p-b)\widehat{w}_b - (p-\widehat{u}_0)w_b, \quad t_b = (q-b)\widehat{w}_b - (q-\widehat{u}_0)w_b, \quad (4.3a)$$

$$(\widehat{z}_b + p\widehat{t}_b)/\widehat{w}_b - (\widehat{z}_b + q\widehat{t}_b)/\widehat{w}_b = (p-b)\widehat{t}_b/\widehat{w}_b - (q-b)\widehat{t}_b/\widehat{w}_b, \quad (4.3b)$$

$$\widehat{u}_{0,1} - \widehat{u}_{0,1} = (p-q + \widehat{u}_0 - \widehat{u}_0)u_0 - p\widehat{u}_0 + q\widehat{u}_0, \quad (4.3c)$$

²One can introduce a particular measure $d\lambda_j(k)$ in (2.1) and this will finally lead to an explicit expression of \mathbf{U} by means of a Cauchy matrix. Thus the objects $v_a, w_b, u_{i,j}$, etc are well defined. One can refer to [44, Section 5] for the DBSQ case and [31, Section 6.4] for the elliptic solution case.

$$2z_b = (p-b)\tilde{t}_b + (q-b)\hat{t}_b - (p+q)t_b + w_b((p+q+\tilde{u}_0+\hat{u}_0)\hat{\tilde{u}}_0 - 2\hat{\tilde{u}}_{0,1} - q\tilde{u}_0 - p\hat{u}_0), \quad (4.3d)$$

$$\begin{aligned} & [H(p,q) + u_{0,1} + (p+q-\alpha_3)u_0 - ((p+q-\alpha_3+u_0)\hat{\tilde{t}}_b - \hat{\tilde{z}}_b)/\hat{\tilde{w}}_b] \\ & \times (p-q+\hat{u}_0-\tilde{u}_0) = (p_b\hat{w}_b - q_b\tilde{w}_b)/\hat{\tilde{w}}_b, \end{aligned} \quad (4.3e)$$

where (4.3d) is the extra equation we added. This system can evolve with initial values that are given on the staircase.

For the closed form (3.21), we need to have an extra equation involved with $u_{0,2}$ and $\hat{\tilde{u}}_{0,2}$. To achieve that, taking $\hat{\tilde{\cdot}}$ -shift of the first equation in (3.18b) and making use of the second equation in (3.18a), we have

$$\hat{\tilde{u}}_{0,2} = p\hat{\tilde{u}}_{0,1} - (p-\hat{\tilde{u}}_0)\hat{\tilde{u}}_{0,1} - (q+u_0)\hat{u}_{1,0} + qu_{1,0} + u_{2,0}. \quad (4.4a)$$

Similarly, we also have

$$\hat{\tilde{u}}_{0,2} = q\hat{\tilde{u}}_{0,1} - (q-\hat{\tilde{u}}_0)\hat{\tilde{u}}_{0,1} - (p+u_0)\tilde{u}_{1,0} + pu_{1,0} + u_{2,0}. \quad (4.4b)$$

Adding (4.4a) to (4.4b) yields

$$\begin{aligned} 2\hat{\tilde{u}}_{0,2} &= (p+q)(u_{1,0} + \hat{\tilde{u}}_{0,1}) - (p-\hat{\tilde{u}}_0)\hat{\tilde{u}}_{0,1} - (q-\hat{\tilde{u}}_0)\tilde{u}_{0,1} \\ & - (p+u_0)\tilde{u}_{1,0} - (q+u_0)\hat{u}_{1,0} + 2u_{2,0}. \end{aligned} \quad (4.5)$$

Inserting this equation into (3.21) leads to lattice equations

$$p\tilde{u}_0 - \tilde{u}_{0,1} = pu_0 + u_{1,0} - \tilde{u}_0u_0, \quad q\hat{u}_0 - \hat{u}_{0,1} = qu_0 + u_{1,0} - \hat{u}_0u_0, \quad (4.6a)$$

$$\hat{u}_{2,0} - \tilde{u}_{2,0} = -p\hat{u}_{1,0} + q\tilde{u}_{1,0} + (p-q+\hat{u}_0-\tilde{u}_0)\hat{\tilde{u}}_{1,0}, \quad (4.6b)$$

$$\hat{u}_{0,2} - \tilde{u}_{0,2} = -p\tilde{u}_{0,1} + q\hat{u}_{0,1} + (p-q+\hat{u}_0-\tilde{u}_0)u_{0,1}, \quad (4.6c)$$

$$\begin{aligned} 2\hat{\tilde{u}}_{0,2} &= (p+q)(u_{1,0} + \hat{\tilde{u}}_{0,1}) - (p-\hat{\tilde{u}}_0)\hat{\tilde{u}}_{0,1} - (q-\hat{\tilde{u}}_0)\tilde{u}_{0,1} \\ & - (p+u_0)\tilde{u}_{1,0} - (q+u_0)\hat{u}_{1,0} + 2u_{2,0}, \end{aligned} \quad (4.6d)$$

$$\begin{aligned} G_4(-p,-q)/(p-q+\hat{u}_0-\tilde{u}_0) &= G_4(-p,-q)/(p-q) + H(p,q)(u_0-\hat{\tilde{u}}_0) \\ & - (p+q-\alpha_3)(u_0\hat{\tilde{u}}_0 - \hat{\tilde{u}}_{1,0} + u_{0,1}) + u_0\hat{\tilde{u}}_{1,0} - \hat{\tilde{u}}_0u_{0,1} - \hat{\tilde{u}}_{2,0} + u_{0,2}, \end{aligned} \quad (4.6e)$$

which can evolve with initial values that are given on the staircase.

With regard to the closed form (3.29), unfortunately, we did not find such extra equations so that (3.29) allows more choices for boundary conditions.

4.2 Deformation of the lattice GD-4 type equations

In [17], Hietarinta obtained DBSQ type equations that are quadrilateral and multidimensionally consistent. Those equations are named (A-2), (B-2) and (C-3). Later, it is shown that Hietarinta's equations can be obtained as deformations of the extended DBSQ equations derived from the DL approach [44]. In the following, we introduce Hietarinta's form of the lattice GD-4 equations (4.2), (4.3), (4.6) and (3.29). In terms of naming conventions, we will call (4.2) the GD-4 (A-2) equation, (4.6) the GD-4 (B-2) equation, and (3.29) the GD-4 (C-3) equation, as they are the lattice GD-4 counterparts of the DBSQ (A-2), (B-2) and (C-3) equations. Equation (4.3) also belongs to the (A-2) type as it is connected to (4.2) by reflection transformations. Here, by "deformation" we mean to convert the equations obtained from the DL approach to Hietarinta's form using point transformations.

GD-4 (A-2): Through the point transformation

$$v_a = x/x_a, \quad u_0 = z - z_0, \quad s_a = (y - v_a y_a)/x_a, \quad (4.7a)$$

$$u_{1,0} = \xi - z_0 u_0 - \xi_0, \quad r_a = (\eta - z_0 y + \xi_0 x)/x_a, \quad (4.7b)$$

where

$$x_a = (p - a)^{-n} (q - a)^{-m} c_0, \quad z_0 = -pn - qm - c_1, \quad y_a = x_a z_0,$$

$$\xi_0 = ((np + mq + c_1)^2 + (np^2 + mq^2 + c_2))/2 + c_3,$$

and c_i are constants, the lattice equation set (4.2) can be transformed into

$$\tilde{y} = z\tilde{x} - x, \quad \hat{y} = z\hat{x} - x, \quad (4.8a)$$

$$\tilde{\eta}\hat{x} - \hat{\eta}\tilde{x} = y(\tilde{x} - \hat{x}), \quad (4.8b)$$

$$\hat{\xi} - \tilde{\xi} = (\hat{z} - \tilde{z})\hat{z}, \quad (4.8c)$$

$$\hat{\eta} = \frac{(\tilde{y}\hat{z} - \hat{y}\tilde{z})z - \xi(\tilde{y} - \hat{y})}{z(\tilde{z} - \hat{z})}, \quad (4.8d)$$

$$\eta = -x\hat{\xi} + y\hat{z} + \alpha_3(y - x\hat{z}) - \alpha_2 x + \frac{G_4(-p, -a)\tilde{x} - G_4(-q, -a)\hat{x}}{\hat{z} - \tilde{z}}. \quad (4.8e)$$

Note that the current equation (4.8d) has been modified. In fact, in light of the transformation (4.7), equation (4.2c) gives rise to

$$2\hat{\eta} = \hat{x}(z(\tilde{z} + \hat{z}) - 2\xi) - (\tilde{y} + \hat{y}). \quad (4.9)$$

Making use of (4.8a), we have

$$\hat{x} = \frac{\tilde{y} - \hat{y}}{z(\tilde{z} - \hat{z})},$$

substituting which into (4.9) and then we get (4.8d). We call this set of equations (4.8) the GD-4 (A-2) equation. Note that compared with (4.9), the form (4.8d) has advantage in generating Lax pair. Analogously, by point transformation

$$w_b = x/x_b, \quad u_0 = z - z_0, \quad t_b = (y - w_b y_b)/x_b,$$

$$u_{0,1} = \xi - z_0 u_0 - \xi_0, \quad z_b = (\eta - z_0 y + x\xi_0)/x_b,$$

with

$$x_b = (-p + b)^n (-q + b)^m c_0, \quad z_0 = -pn - qm - c_1, \quad y_b = x_b z_0,$$

$$\xi_0 = ((np + mq + c_1)^2 - (np^2 + mq^2 + c_2))/2 - c_3,$$

and constants c_i , the system (4.3) gives rise to

$$y = \tilde{z}x - \tilde{x}, \quad y = \hat{z}x - \hat{x}, \quad (4.10a)$$

$$\tilde{\eta}\hat{x} - \hat{\eta}\tilde{x} = \hat{y}(\tilde{x} - \hat{x}), \quad (4.10b)$$

$$\hat{\xi} - \tilde{\xi} = (\hat{z} - \tilde{z})z, \quad (4.10c)$$

$$\hat{\xi} = \frac{-\eta(\hat{x} - \tilde{x}) + y(\hat{y} - \tilde{y}) + \tilde{x}\hat{y} - \hat{x}\tilde{y}}{x(\hat{x} - \tilde{x})}, \quad (4.10d)$$

$$\hat{\eta} = -\hat{x}\hat{\xi} + \hat{y}z - \alpha_3(\hat{y} - \hat{x}z) - \alpha_2 \hat{x} - \frac{G_4(-p, -b)\hat{x} - G_4(-q, -b)\tilde{x}}{\hat{z} - \tilde{z}}. \quad (4.10e)$$

Here we already modified (4.10d). In fact, (4.3d) gives rise to

$$2\eta = -x(2\widehat{\xi} - (\widetilde{z} + \widehat{z})\widehat{z}) - (\widetilde{y} + \widehat{y}). \quad (4.11)$$

From (4.10a), we have

$$\widehat{z} = \frac{\widehat{y} - \widetilde{y}}{\widehat{x} - \widetilde{x}}, \quad x(\widetilde{z} + \widehat{x}) = 2y + \widetilde{x} + \widehat{x},$$

using which we eliminate z variable from (4.11) yields (4.10d). Equations (4.10) compose an alternative form of the GD-4 (A-2) equation. It is related to (4.8) by reflection transformations

$$p \rightarrow -p, \quad q \rightarrow -q, \quad n \rightarrow -n, \quad m \rightarrow -m, \quad \alpha_3 \rightarrow -\alpha_3, \quad \alpha_1 \rightarrow -\alpha_1, \quad b \rightarrow -a.$$

Besides, both equation (4.8) and (4.10) allow a symmetry $(\widetilde{\cdot}, p) \leftrightarrow (\widehat{\cdot}, q)$.

GD-4 (B-2): We adopt the point transformation

$$\begin{aligned} u_0 &= x - x_0, & u_{1,0} &= y - x_0 u_0 - y_0, & u_{0,1} &= z - x_0 u_0 - z_0, \\ u_{2,0} &= \xi - x_0 y + z z_0 + \mu, & u_{0,2} &= \eta - x_0 z + z y_0 + \nu, \end{aligned}$$

where

$$\begin{aligned} x_0 &= -pn - qm - c_1, \\ y_0 &= ((np + mq + c_1)^2 + (np^2 + mq^2 + c_2))/2 + c_4, \\ z_0 &= ((np + mq + c_1)^2 - (np^2 + mq^2 + c_2))/2 - c_4, \\ \mu_0 &= (np^3 + mq^3 + c_3)/3, \\ \mu &= -\frac{1}{6}x_0^3 + \frac{1}{2}(y_0 - z_0)x_0 + \frac{1}{3}\mu_0 + c_5, \\ \nu &= -\frac{1}{6}x_0^3 - \frac{1}{2}(y_0 - z_0)x_0 + \frac{1}{3}\mu_0 + c_5, \end{aligned}$$

and c_i are constants. Then from (4.6) we have the GD-4 (B-2) equation

$$\widetilde{z} = x\widetilde{x} - y, \quad \widehat{z} = x\widehat{x} - y, \quad (4.12a)$$

$$\widehat{\xi} - \widetilde{\xi} = (\widehat{x} - \widetilde{x})\widehat{y}, \quad (4.12b)$$

$$\widehat{\eta} - \widetilde{\eta} = (\widehat{x} - \widetilde{x})z, \quad (4.12c)$$

$$\widehat{\eta} = \xi + \frac{(\widehat{x}\widetilde{y} - \widetilde{x}\widehat{y})x - y(\widetilde{y} - \widehat{y})}{\widetilde{x} - \widehat{x}}, \quad (4.12d)$$

$$\eta = \widehat{\xi} - \widehat{y}x + \widehat{x}z + \alpha_3(\widehat{y} - x\widehat{x} + z) + \alpha_2(\widehat{x} - x) + \frac{G_4(-p, -q)}{\widehat{x} - \widetilde{x}}. \quad (4.12e)$$

Again, we have modified (4.12d). Under the transformation (4.6d) gives rise to

$$2\widehat{\eta} = 2\xi - x(\widehat{y} + \widetilde{y}) + (\widehat{z} + \widetilde{z})\widehat{x}. \quad (4.13)$$

From (4.12a), we can have

$$\widehat{\widetilde{x}} = \frac{\widetilde{y} - \widehat{y}}{\widetilde{x} - \widehat{x}}, \quad \widetilde{z} + \widehat{z} = (\widetilde{x} + \widehat{x})x - 2y.$$

Substituting them into (4.13) yields (4.12d). Note that (4.12) allows a symmetry $(\widetilde{\cdot}, p) \leftrightarrow (\widehat{\cdot}, q)$.

GD-4 (C-3): Inserting point transformation

$$S_{a,b} = ((p-a)/(p-b))^n ((q-a)/(q-b))^m x, \quad (4.14a)$$

$$v_a = (p - a)^n (q - a)^m y, \quad w_b = (p - b)^{-n} (q - b)^{-m} z, \quad (4.14b)$$

$$s_a = (p - a)^n (q - a)^m (y_1 - z_0 y), \quad (4.14c)$$

$$t_b = (p - b)^{-n} (q - b)^{-m} (z_1 - z_0 z), \quad (4.14d)$$

into lattice equation (3.29), where $z_0 = -pn - qm - c_1$, we get the GD-4 (C-3) equation

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (4.15a)$$

$$(\tilde{y}_1 + y)\hat{y} = (\hat{y}_1 + y)\tilde{y}, \quad (\tilde{z}_1 - \hat{z})\tilde{z} = (\tilde{z}_1 - \hat{z})\hat{z}, \quad (4.15b)$$

$$y\hat{z}_1 - y_1\hat{z} = z \frac{G_4(-p, -b)\tilde{z}\tilde{y} - G_4(-q, -b)\tilde{z}\hat{y}}{\hat{z} - \tilde{z}} - \alpha_3 y\hat{z} + G_4(-a, -b)\hat{x}. \quad (4.15c)$$

The alternative GD-4 (C-3) equation is

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (4.16a)$$

$$(\tilde{y}_1 + y)\hat{y} = (\hat{y}_1 + y)\tilde{y}, \quad (\hat{z}_1 - \tilde{z})\tilde{z} = (\hat{z}_1 - \tilde{z})\hat{z}, \quad (4.16b)$$

$$y\hat{z}_1 - y_1\hat{z} = z \frac{G_4(-p, -a)\tilde{z}\tilde{y} - G_4(-q, -a)\tilde{z}\hat{y}}{\hat{z} - \tilde{z}} - \alpha_3 y\hat{z} + G_4(-a, -b)x, \quad (4.16c)$$

where the last equation is the deformation of (3.30) in light of the transformation (4.14). Equations (4.15) and (4.16) are related by the reversal symmetry transformation

$$n \rightarrow -n, \quad m \rightarrow -m, \quad y \rightarrow z, \quad z \rightarrow -y, \quad y_1 \rightarrow z_1, \quad z_1 \rightarrow -y_1, \quad a \leftrightarrow b.$$

Both equation (4.15) and (4.16) allow a symmetry $(\sim, p) \leftrightarrow (\hat{\sim}, q)$. Besides, (4.15c) and (4.16c) are also connected through

$$z(\hat{z}\tilde{y} - \tilde{z}\hat{y}) = (\hat{x} - x)(\tilde{z} - \hat{z}),$$

which holds in light of (4.15a). In addition, we note that, when $a = b$ in (4.16c) and we eliminate x from (4.16a), we can reduce the GD-4 (C-3) equation (4.16) to the following four-component form

$$\begin{aligned} \hat{y}(\tilde{z} - \hat{z}) &= -z(\tilde{y} - \hat{y}), & (\tilde{y}_1 + y)\hat{y} &= (\hat{y}_1 + y)\tilde{y}, & (\hat{z}_1 - \tilde{z})\tilde{z} &= (\tilde{z}_1 - \hat{z})\hat{z}, \\ y\hat{z}_1 - y_1\hat{z} &= z \frac{G_4(-p, -a)\tilde{z}\tilde{y} - G_4(-q, -a)\tilde{z}\hat{y}}{\hat{z} - \tilde{z}} - \alpha_3 y\hat{z}. \end{aligned}$$

GD-4 (C-4): In [17], Hietarinta also obtained a (C-4) equation of the DBSQ type (see (A.4)), which was later shown to be another deformation of the (C-3) equation. In the following, we show a similar deformation and present the GD-4 (C-4) equation.

First, we observe that equations (4.15) and (4.16) share the same solution through (4.14). Thus, from (4.15) and (4.16), we have

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (4.17a)$$

$$(\tilde{y}_1 + y)\hat{y} = (\hat{y}_1 + y)\tilde{y}, \quad (\hat{z}_1 - \tilde{z}_1)\tilde{z} = (\tilde{z}_1 - \hat{z}_1)\hat{z}, \quad (4.17b)$$

$$y\hat{z}_1 - y_1\hat{z} = z \frac{P_{a,b}\tilde{z}\tilde{y} - Q_{a,b}\tilde{z}\hat{y}}{\hat{z} - \tilde{z}} - \alpha_3 y\hat{z} + G_{a,b}(\hat{x} + x), \quad (4.17c)$$

where

$$\begin{aligned} P_{a,b} &= (G_4(-p, -b) + G_4(-p, -a))/2, \\ Q_{a,b} &= (G_4(-q, -b) + G_4(-q, -a))/2, \quad G_{a,b} = G_4(-a, -b)/2. \end{aligned}$$

Then, consider the following transformation

$$\begin{aligned} x &= \frac{x_2 - G_{a,b}}{2G_{a,b}(x_2 + G_{a,b})}, & y &= \frac{y_2}{x_2 + G_{a,b}}, & z &= \frac{\tilde{z}_2}{x_2 + G_{a,b}}, \\ y_1 &= \frac{y_3}{x_2 + G_{a,b}}, & z_1 &= \frac{z_3}{x_2 + G_{a,b}}. \end{aligned} \quad (4.18)$$

Imposing this transformation on (4.17a) and (4.17b) gives rise to

$$\begin{aligned} x_2 - \tilde{x}_2 &= \tilde{y}_2 z_2, & x_2 - \hat{x}_2 &= \hat{y}_2 z_2, \\ (\tilde{y}_3 + y_2)\hat{y}_2 &= (\hat{y}_3 + y_2)\tilde{y}_2, & (\tilde{z}_3 - \hat{z}_2)\tilde{z}_2 &= (\hat{z}_3 - \tilde{z}_2)\hat{z}_2, \end{aligned}$$

which further yields

$$\hat{\tilde{x}}_2 = \frac{\hat{x}_2 \tilde{z}_2 - \tilde{x}_2 \hat{z}_2}{\tilde{z}_2 - \hat{z}_2}. \quad (4.19)$$

Meanwhile, by the transformation (4.18), equation (4.17c) multiplied by $(G_{a,b} + x_2)(G_{a,b} + \hat{x}_2)$ is equivalent to

$$y_2 \hat{\tilde{z}}_3 - y_3 \hat{\tilde{z}}_2 = \Delta - \alpha_3 y_2 \hat{\tilde{z}}_2 + x_2 \hat{x}_2 - G_{a,b}^2, \quad (4.20)$$

where

$$\Delta = \frac{z_2(G_{a,b} + \hat{x}_2)(P_{a,b} \hat{z}_2 \tilde{y}_2 - Q_{a,b} \tilde{z}_2 \hat{y}_2)}{(\hat{z}_2 \tilde{x}_2 - \tilde{z}_2 \hat{x}_2) + G_{a,b}(\hat{z}_2 - \tilde{z}_2)}.$$

We substitute (4.19) into Δ and rewrite (4.20) as

$$y_2 \hat{\tilde{z}}_3 - y_3 \hat{\tilde{z}}_2 = z_2 (P_{a,b} \hat{z}_2 \tilde{y}_2 - Q_{a,b} \tilde{z}_2 \hat{y}_2) / (\hat{z}_2 - \tilde{z}_2) - \alpha_3 y_2 \hat{\tilde{z}}_2 + x_2 \hat{x}_2 - G_{a,b}^2.$$

Finally, we arrive at the GD-4 (C-4) equation

$$x_2 - \tilde{x}_2 = \tilde{y}_2 z_2, \quad x_2 - \hat{x}_2 = \hat{y}_2 z_2, \quad (4.21a)$$

$$(\tilde{y}_3 + y_2)\hat{y}_2 = (\hat{y}_3 + y_2)\tilde{y}_2, \quad (\hat{z}_3 - \tilde{z}_2)\tilde{z}_2 = (\tilde{z}_3 - \hat{z}_2)\hat{z}_2, \quad (4.21b)$$

$$y_2 \hat{\tilde{z}}_3 - y_3 \hat{\tilde{z}}_2 = z_2 (P_{a,b} \hat{z}_2 \tilde{y}_2 - Q_{a,b} \tilde{z}_2 \hat{y}_2) / (\hat{z}_2 - \tilde{z}_2) - \alpha_3 y_2 \hat{\tilde{z}}_2 + x_2 \hat{x}_2 - G_{a,b}^2. \quad (4.21c)$$

Four-component GD-4 (A-2): In order to reduce the GD-4 (A-2) equation (4.8) into four-component form, we introduce variables

$$u = \frac{y}{x}, \quad v = \frac{\eta}{x}.$$

Then, equation (4.8) can be rewritten as

$$(\hat{u} - z)(\hat{u} - \hat{z}) = (\tilde{u} - z)(\tilde{u} - \tilde{z}), \quad (4.22a)$$

$$\tilde{v} - \hat{v} = u(\tilde{u} - \hat{u}), \quad (4.22b)$$

$$\hat{\xi} - \tilde{\xi} = (\hat{z} - \tilde{z})\hat{z}, \quad (4.22c)$$

$$\hat{v} = \frac{1}{z(\tilde{z} - \hat{z})} [z((\tilde{u} - \hat{u})\tilde{z}\hat{z} + (\hat{u}\tilde{z} - \tilde{u}\hat{z})\hat{u}) - \xi(\tilde{u}\tilde{z} - \hat{u}\hat{z} + (\tilde{u} - \hat{u})\hat{u})], \quad (4.22d)$$

$$(\hat{z} - \tilde{z})[v + \hat{\xi} - (u - \alpha_3)\hat{z} - \alpha_3 u + \alpha_2] = \frac{G_4(-p, -a)}{z - \tilde{u}} - \frac{G_4(-q, -a)}{z - \hat{u}}. \quad (4.22e)$$

We view this system as the four-component GD-4 (A-2) equation. Similarly, by introducing variables $w = \frac{y}{x}$, $\varpi = \frac{\eta}{x}$, system (4.10) becomes

$$(\widehat{z} - \widehat{w})(\widehat{z} - w) = (\widehat{z} - \widetilde{w})(\widetilde{z} - w), \quad (4.23a)$$

$$\widetilde{\varpi} - \widehat{\varpi} = \widehat{w}(\widetilde{w} - \widehat{w}), \quad (4.23b)$$

$$\widehat{\xi} - \widetilde{\xi} = (\widehat{z} - \widetilde{z})z, \quad (4.23c)$$

$$\widehat{\xi} = \frac{-\widetilde{\varpi}(\widehat{z} - \widetilde{z}) + w(\widehat{w}(\widehat{z} - w) - \widetilde{w}(\widetilde{z} - w)) + \widehat{w}(\widetilde{z} - w)(\widehat{z} - w) - \widetilde{w}(\widehat{z} - w)(\widetilde{z} - w)}{\widehat{z} - \widetilde{z}}, \quad (4.23d)$$

$$(\widehat{\varpi} + \xi - z\widehat{w} + \alpha_3(\widehat{w} - z) + \alpha_2)(\widehat{z} - \widetilde{z}) = \frac{G_4(-p, -b)}{\widehat{u} - \widehat{z}} - \frac{G_4(-q, -b)}{\widetilde{u} - \widetilde{z}}, \quad (4.23e)$$

which is viewed as the alternative four-component GD-4 (A-2) equation.

Note that both the equations in terms of DL notations (3.3) and in Hietarinta's form in terms of x, y, z , etc can be used to express GD-4 lattice equations. The DL approach enables us to see how these equations follow from one and the same underlying structure, while the equations in Hietarinta's form are more convenient to investigate their evolution and MDC property.

As for the names (A-2), (B-2) and (C-3), we have the following comments. On one side, as we can see, the DBSQ (A-2), (B-2) and (C-3) equations in Hietarinta's form are embedded in the corresponding GD-4 lattice equations. On the other side, in the DBSQ case, one needs to employ more than one variables so that the system is quadrilateral. Different variables from a same system can express different equations in one-component form. For example, the DBSQ (A-2) equation (A.2) can give rise to the modified DBSQ equation in terms of x and the regular DBSQ equation in terms of either z or $w = y/x$, see [22, Section 3.2]. In this context, it is not correct to call the DBSQ (A-2) equation the modified DBSQ or regular DBSQ equation. We prefer to use (A-2), (B-2) and (C-3) in the GD-4 case.

4.3 General description of MDC

In the past two decades, the property of multidimensional consistency (MDC) has been effectively used in the study of discrete integrable systems. This property allows a lattice equation (or a system) to be consistently embedded into a higher dimension [1, 5, 27, 32]. With regard to the lattice GD-4 type equations we listed in Section 4.2, they can be written as a quad equation

$$\mathcal{F}(\boldsymbol{\theta}, \widetilde{\boldsymbol{\theta}}, \widehat{\boldsymbol{\theta}}, \overline{\boldsymbol{\theta}}; p, q) = 0, \quad (4.24)$$

where \mathcal{F} is a nonlinear vector function of the vector $\boldsymbol{\theta}$ with several components. The key idea of the MDC property is to embed the equation consistently into a multi-dimensional lattice by imposing copies of the same equation, albeit with different lattice parameters in different directions. The GD-4 type equations are symmetric between the (p, n) and (q, m) coordinates of the \mathbb{Z}^2 lattice, and therefore we introduce a third dimension l and the corresponding lattice index r by $\boldsymbol{\theta} \rightarrow \overline{\boldsymbol{\theta}}$ to keep this symmetry. This means that we have the same equation on all planes around the cube (see Figure 1). The equations on the six faces can be written as

$$\mathcal{F}(\boldsymbol{\theta}, \widetilde{\boldsymbol{\theta}}, \widehat{\boldsymbol{\theta}}, \overline{\boldsymbol{\theta}}; p, q) = 0 \quad (\text{bottom}), \quad \mathcal{F}(\overline{\boldsymbol{\theta}}, \widetilde{\overline{\boldsymbol{\theta}}}, \widehat{\overline{\boldsymbol{\theta}}}, \overline{\overline{\boldsymbol{\theta}}}; p, q) = 0 \quad (\text{top}), \quad (4.25a)$$

$$\mathcal{F}(\boldsymbol{\theta}, \widetilde{\boldsymbol{\theta}}, \overline{\boldsymbol{\theta}}, \widehat{\boldsymbol{\theta}}; p, r) = 0 \quad (\text{left}), \quad \mathcal{F}(\widehat{\boldsymbol{\theta}}, \widehat{\widetilde{\boldsymbol{\theta}}}, \widehat{\overline{\boldsymbol{\theta}}}, \widehat{\overline{\boldsymbol{\theta}}}; p, r) = 0 \quad (\text{right}), \quad (4.25b)$$

$$\mathcal{F}(\boldsymbol{\theta}, \overline{\boldsymbol{\theta}}, \widehat{\boldsymbol{\theta}}, \widetilde{\boldsymbol{\theta}}; q, r) = 0 \quad (\text{back}), \quad \mathcal{F}(\widetilde{\boldsymbol{\theta}}, \widetilde{\overline{\boldsymbol{\theta}}}, \widetilde{\widehat{\boldsymbol{\theta}}}, \widetilde{\widetilde{\boldsymbol{\theta}}}; q, r) = 0 \quad (\text{front}). \quad (4.25c)$$

For a quadrilateral equation, e.g., (4.24), that can be solved for the fourth variable given the first three, its MDC property can be explained as follows: In the case of the initial values given

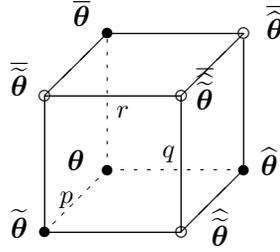


Figure 1. Consistent cube of the lattice equation (4.24).

at black dots $(\theta, \tilde{\theta}, \hat{\theta}, \bar{\theta})$ (see Figure 1), we use the left-hand side equations in (4.25a), (4.25b), (4.25c) to compute $\hat{\tilde{\theta}}, \bar{\tilde{\theta}}, \hat{\bar{\theta}}$, respectively, and this leaves the three right-hand side equations from $\bar{\tilde{\theta}}$ which we should get the $\hat{\tilde{\theta}}$. The MDC property means that all three values thus obtained for $\hat{\tilde{\theta}}$ coincide. The MDC property for quadrilateral equations can be geometrically interpreted as a consistency around a cube (CAC). In light of CAC property, in principle, the coupled system

$$\mathcal{F}(\theta, \tilde{\theta}, \bar{\theta}, \hat{\theta}; p, r) = 0, \quad (4.26a)$$

$$\mathcal{F}(\theta, \bar{\theta}, \hat{\theta}, \tilde{\theta}; q, r) = 0 \quad (4.26b)$$

act as a Bäcklund transformation for the equation (4.24) to transform solutions between θ and $\bar{\theta}$. From the Bäcklund transformation, in principle, Lax pairs can be constructed [1, 6, 27].

4.4 Evolution

To explore the MDC property of the GD-4 lattice equations, we analyze their evolutions. Let us take the GD-4 (B-2) equation (4.12) as an example. It is a system consisting of six equations for five variables x, y, z, ξ and η . In general, such a system is over-determined. We have the following remarks for such a situation:

Remark 4.1. If we ONLY require the equations to evolve with a certain staircase initial boundary, we will not need all 6 equations.

For convenience, we introduce V to denote $V = (x, y, z, \xi, \eta)^T$. It turns out that, the system (4.12), excluding (4.12d), allows an up-left evolution for given V, \tilde{V}, \hat{V} as in Figure 2 (b) or a down-right evolution for a given V, \hat{V}, \tilde{V} . In addition, without (4.12c) the system (4.12) allows an up-right evolution, and without (4.12b) it allows a down-left evolution. Now it is clear that we may exclude one equation from (4.12b, 4.12c, 4.12d), and the remaining system will consist of five equations and evolve with some suitable staircase initial values.

Remark 4.2. If we require the system (4.12) to be multidimensionally consistent, we need ALL six equations. See Section 5 and Appendix C. Similar features of evolution hold for the GD-4 (A-2) equations (4.8) and (4.10) and for the four-component GD-4 (A-2) equations (4.22) and (4.23).

Remark 4.3. The GD-4 (C-3) equation (4.15) (or (4.16)) and the GD-4 (C-4) equation (4.21) consist of 5 equations, respectively. They can have either up-left or down-right evolution, but neither up-right nor down-left evolution.

The MDC property implies certain integrability. Since for the GD-4 equations (4.8), (4.10), (4.12), (4.22) and (4.23), all equations are necessary to guarantee each system that they compose to be MDC, we prefer to keep all the equations in each system.

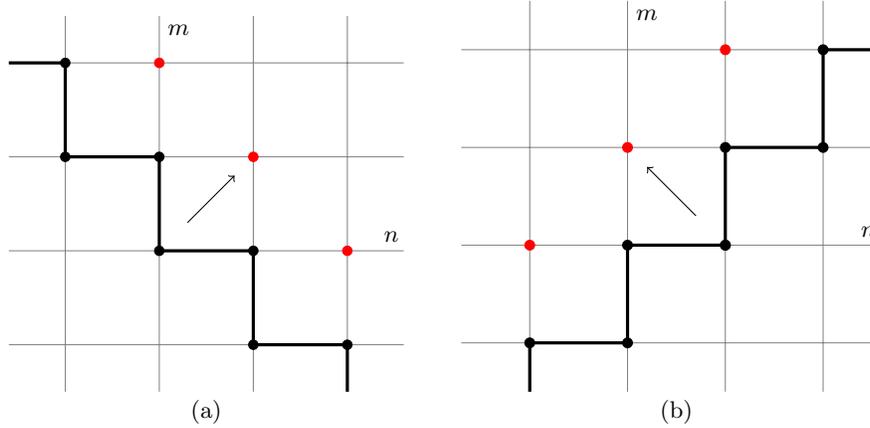


Figure 2. The initial values should be given on black dots of the staircase. One can compute the values at red dots. This way, the equations can evolve up-right as in (a) and up-left as in (b).

5 MDC property and Lax representation

5.1 MDC property of the GD-4 (A-2) and (B-2) equation

We have examined the evolution of the GD-4 type equations in Section 4.4. Among them, the GD-4 (A-2) equation (4.8) and its alternative form (4.10), the GD-4 (B-2) equation (4.12), four-component GD-4 (A-2) equation (4.22) and its alternative form (4.23), allow up-right evolution with staircase initial values. Therefore, we can check whether these equations are consistent around the cube (CAC). Note that equations (4.8b), (4.10c), (4.12b), (4.22b) and (4.23c) are not necessary to the systems where they are to get up-right evolutions. However, in the following we will see that these equations are really needed for their systems to be CAC. The checking is straightforward.

GD-4 (A-2) equation: For the GD-4 (A-2) equation (4.8), we set it to be (4.24). It is clear that from the left-hand side equations in (4.25a), (4.25b), (4.25c), we evaluate the following second-order shifts

$$\tilde{\bar{x}} = (\tilde{x} - \bar{x})/(\tilde{z} - \bar{z}), \quad \tilde{\bar{y}} = (\tilde{x}\tilde{z} - \bar{x}\bar{z})/(\tilde{z} - \bar{z}), \quad \tilde{\bar{z}} = (\tilde{\xi} - \bar{\xi})/(\tilde{z} - \bar{z}), \quad (5.1a)$$

$$\hat{\bar{x}} = (\bar{x} - \hat{x})/(\bar{z} - \hat{z}), \quad \hat{\bar{y}} = (\bar{x}\hat{z} - \hat{x}\bar{z})/(\bar{z} - \hat{z}), \quad \hat{\bar{z}} = (\bar{\xi} - \hat{\xi})/(\bar{z} - \hat{z}), \quad (5.1b)$$

$$\tilde{\bar{\xi}} = (y/x - \alpha_3) \frac{\tilde{\xi} - \bar{\xi}}{\tilde{z} - \bar{z}} + \frac{\alpha_3 y - \eta}{x} - \alpha_2 - \frac{G_4(-p, -a)\tilde{x} - G_4(-r, -a)\bar{x}}{x(\tilde{z} - \bar{z})}, \quad (5.1c)$$

$$\hat{\bar{\xi}} = (y/x - \alpha_3) \frac{\bar{\xi} - \hat{\xi}}{\bar{z} - \hat{z}} + \frac{\alpha_3 y - \eta}{x} - \alpha_2 - \frac{G_4(-r, -a)\bar{x} - G_4(-q, -a)\hat{x}}{x(\bar{z} - \hat{z})}, \quad (5.1d)$$

$$\tilde{\bar{\eta}} = \left(\tilde{y}\tilde{z} - \bar{y}\bar{z} - \frac{\tilde{\xi}}{z}(\tilde{y} - \bar{y}) \right) / (\tilde{z} - \bar{z}), \quad \hat{\bar{\eta}} = \left(\hat{y}\hat{z} - \bar{y}\bar{z} - \frac{\hat{\xi}}{z}(\hat{y} - \bar{y}) \right) / (\hat{z} - \bar{z}). \quad (5.1e)$$

Furthermore, from the three right-hand side equations in (4.25) and using the data (5.1), we have three different ways to calculate the values $\tilde{\bar{x}}$, $\tilde{\bar{y}}$, $\tilde{\bar{z}}$ and $\tilde{\bar{\eta}}$ uniquely, which are

$$\begin{aligned} \tilde{\bar{x}} &= [(\tilde{x} - \hat{x})\bar{z} + (\hat{x} - \bar{x})\tilde{z} + (\bar{x} - \tilde{x})\hat{z}] / [(\tilde{\xi} - \hat{\xi})\bar{z} + (\hat{\xi} - \bar{\xi})\tilde{z} + (\bar{\xi} - \tilde{\xi})\hat{z}], \\ \tilde{\bar{y}} &= [(\hat{x} - \tilde{x})\bar{\xi} + (\tilde{x} - \bar{x})\hat{\xi} + (\bar{x} - \hat{x})\tilde{\xi}] / [(\tilde{z} - \hat{z})\bar{\xi} + (\hat{z} - \bar{z})\tilde{\xi} + (\bar{z} - \tilde{z})\hat{\xi}], \\ \tilde{\bar{z}} &= [(\alpha_3 x - y)((\hat{z} - \bar{z})\tilde{\xi} + (\bar{z} - \tilde{z})\hat{\xi} + (\tilde{z} - \hat{z})\bar{\xi}) \\ &\quad + \tilde{x}(\hat{z} - \bar{z})G_4(-p, -a) + \hat{x}(\bar{z} - \tilde{z})G_4(-q, -a) + \bar{x}(\tilde{z} - \hat{z})G_4(-r, -a)] \end{aligned}$$

$$\begin{aligned} & / [x((\widehat{z} - \widetilde{z})\bar{\xi} + (\widetilde{z} - \bar{z})\widehat{\xi} + (\bar{z} - \widehat{z})\widetilde{\xi})], \\ \widetilde{\widetilde{\eta}} = & \frac{-\bar{\xi}\left(\frac{\widehat{xz}-\bar{xz}}{\bar{z}-\bar{z}} - \frac{\widehat{xz}-\bar{xz}}{\bar{z}-\bar{z}}\right) + \bar{z}\left(\frac{(\bar{\xi}-\widehat{\xi})(\widehat{xz}-\bar{xz})}{(\bar{z}-\bar{z})(\bar{z}-\bar{z})} - \frac{(\bar{\xi}-\widetilde{\xi})(\widehat{xz}-\bar{xz})}{(\bar{z}-\bar{z})(\bar{z}-\bar{z})}\right)}{\bar{z}\left(\frac{\bar{\xi}-\widetilde{\xi}}{\bar{z}-\bar{z}} - \frac{\bar{\xi}-\widehat{\xi}}{\bar{z}-\bar{z}}\right)}. \end{aligned}$$

For the variable ξ , the formula for $\widetilde{\widetilde{\xi}}$ is too long to be listed here. It turns out that to guarantee the relation $\widetilde{\widetilde{\xi}} = \widetilde{\xi} = \widehat{\xi}$ holds one needs

$$\widetilde{\widetilde{\eta}}\widehat{x} - \widetilde{\widetilde{\eta}}\widetilde{x} = y(\widetilde{x} - \widehat{x}), \quad \bar{\eta}\widehat{x} - \bar{\eta}\bar{x} = y(\bar{x} - \widehat{x}), \quad \widetilde{\eta}\bar{x} - \widetilde{\eta}\widetilde{x} = y(\widetilde{x} - \bar{x})$$

hold. This is nothing but the second equation (4.8b) in (4.8) and its $(-, \widehat{\cdot})$ and $(-, \widetilde{\cdot})$ versions. Thus we can conclude two points. One is that although (4.8b) is not needed for determining the up-right evolution of (4.8), it is necessary for the whole system (4.8) to be MDC. In other words, the GD-4 (A-2) equation (4.8) as a whole system satisfies the MDC property. The other point is, if, without (4.8b) the system (4.8) is integrable in the sense of being MDC, then (4.8b) is a consequence of the other five equations in (4.8).

For the alternative GD-4 (A-2) equation (4.10), the GD-4 (B-2) equation (4.12), the four-component GD-4 (A-2) equation (4.22) and its alternative form (4.23), their MDC property can be checked in a similar way and has similar results. The check is straightforward and we present formulas of those triple shifts of variables in Appendix C. What we want to emphasize is, for η in the alternative GD-4 (A-2) equation (4.10), ξ in the GD-4 (B-2) equation (4.12), ξ in the four-component GD-4 (A-2) equation (4.22), and ϖ in (4.23), their triple shifts are uniquely determined respective require (4.10c), (4.12b), (4.22c), (4.23b) and their $(-, \widehat{\cdot})$ and $(-, \widetilde{\cdot})$ versions hold. These facts provide supports for Remark 4.2. On the other hand, if we agree that these systems are MDC, then (4.10c), (4.12c), (4.22c) and (4.23b) will be the consequence of the rest equations in their own system, respectively.

5.2 Lax representations

Finding a Lax pair for a given nonlinear equation, whether continuous or discrete, is generally a difficult task. The DL scheme provides a direct procedure to construct Lax pairs of the objective equations based on the relation (2.4). One can select specific components of these vectors, or combinations thereof, to constitute the basic vector functions in terms of which we obtain the relevant linear problems. This has been illustrated in the lattice KdV and DBSQ case (e.g., [39]). Besides, the fundamental characterization of integrable partial difference equations as being MDC is intimately related to the existence of a Lax pair [1, 6, 27]. Always the information gained from the process of verifying MDC property is also crucial to the computation of the corresponding Lax pair.

For the quadrilateral ABS equations [1] and 3-point discrete Burgers equation [7, 42], both of which are CAC, they can be obtained as compatibilities of their Lax pairs constructed from the MDC property (see [6, 7, 42]). However, for a DBSQ type equation, e.g., (B-2) or ((A-2), (C-3)), its Lax pairs, no matter constructed from the DL scheme or by means of the MDC property, are incomplete in the sense that the original DBSQ equation can not be fully recovered from the compatibility of its Lax pair (see [22, 44]). Such an incompleteness of the Lax pair also happens to the GD-4 type equations. In the following, we only take the GD-4 (A-2) equation (4.8) as an example and present its Lax pairs derived from the DL scheme and the MDC property, respectively. One can see that, the equations arising from the compatibility of the Lax pairs can not yield all the equations in the GD-4 (A-2).

To determine the Lax pairs of the lattice equations using the DL approach, we should use the relation (3.2). For the sake of the construction of the Lax pairs, we introduce $(\mathbf{v})_0$ standing

for the 0-th component of the infinite-component column vector \mathbf{v} . Besides, we appoint that A^T stands for the transpose of matrix A .

We start by looking at the Lax pair of (3.15), which is the DL counterpart of the GD-4 (A-2) equation (4.8). We define an eigenvector of the form

$$\phi = (((a + \Lambda)^{-1} \mathbf{u}_k)_0, (\mathbf{u}_k)_0, (\Lambda \mathbf{u}_k)_0, (\Lambda^2 \mathbf{u}_k)_0)^T. \quad (5.2)$$

Based on (3.2), one can derive the following linear relation

$$\tilde{\phi} = \mathbf{L}\phi, \quad \hat{\phi} = \mathbf{M}\phi, \quad (5.3a)$$

in which

$$\mathbf{L} = \begin{pmatrix} p - a & \tilde{v}_a & 0 & 0 \\ 0 & p - \tilde{u}_0 & 1 & 0 \\ 0 & -\tilde{u}_{1,0} & p & 1 \\ G_4(k, -a)/v_a & *_{42} & (\alpha_3 s_a - r_a)/v_a - \alpha_2 & A_1(-p) + s_a/v_a \end{pmatrix}, \quad (5.3b)$$

where

$$*_{42} = p_a \tilde{v}_a / v_a - [A_2(-p) + (A_1(-p) s_a + r_a) / v_a] (p - \tilde{u}_0) - \tilde{u}_{1,0} (A_1(-p) + s_a / v_a), \quad (5.3c)$$

and \mathbf{M} is obtained from (5.3b) by replacing p by q and $\tilde{\cdot}$ by $\hat{\cdot}$. The compatibility of (5.3a), i.e., $\tilde{\mathbf{L}}\mathbf{M} = \tilde{\mathbf{M}}\mathbf{L}$, leads to equations (3.15b), (3.15c), (3.15d), (3.11b) and

$$p - q + \hat{s}_a / \hat{v}_a - \tilde{s}_a / \tilde{v}_a = v_a ((p - a) / \tilde{v}_a - (q - a) / \hat{v}_a).$$

The latter two equations follows readily from (3.15a). The incompleteness of the Lax pair is demonstrated by the missing of (3.15a).

It is well known that the MDC property can provide Lax pairs automatically. All the deformed lattice equations listed above admit the MDC property. As an example, in this part we discuss the Lax integrability of the GD-4 (A-2) equation (4.8) thanks to its MDC property. As the discussion of the MDC property, here we extend the aforesaid lattice equations into a third dimension by introducing a new variable l associated with a new shift $^-$ and a new complex parameter r . For the lattice equation (4.24) which is MDC, the essential procedure to construct its Lax pair is to introduce fractional expressions (e.g., f/F , g/G , etc.) for the various components in the third dimension in the Bäcklund transformation (4.26) in order to linearize the numerators and denominators of the expressions for the components of $\tilde{\theta}$ and $\hat{\theta}$ in terms of f , F , g , G , etc. This manipulation generates the Lax matrices.

For the GD-4 (A-2) equation (4.8), we consider its Bäcklund transformation (5.1). Introducing

$$\bar{x} = \frac{\varphi_1}{\varphi_0}, \quad \bar{z} = \frac{\varphi_2}{\varphi_0}, \quad \bar{y} = \frac{\varphi_3}{\varphi_0}, \quad \bar{\xi} = \frac{\varphi_4}{\varphi_0}, \quad \bar{\eta} = \frac{\varphi_5}{\varphi_0},$$

and defining

$$\phi_1 = (\varphi_0, \varphi_1, \varphi_2, \varphi_3, \varphi_4, \varphi_5)^T,$$

from (5.1) we have

$$\tilde{\phi}_1 = \mathbf{L}_1 \phi_1, \quad \hat{\phi}_1 = \mathbf{M}_1 \phi_1,$$

where

$$\mathbf{L}_1 = \begin{pmatrix} \tilde{z} & 0 & -1 & 0 & 0 & 0 \\ \tilde{x} & -1 & 0 & 0 & 0 & 0 \\ \tilde{\xi} & 0 & 0 & 0 & -1 & 0 \\ 0 & -\tilde{z} & \tilde{x} & 0 & 0 & 0 \\ *_{51} & G_4(-r, -a)/x & (\eta - \alpha_3 y)/x + \alpha_2 & 0 & \alpha_3 - y/x & 0 \\ -(\xi/z)\tilde{y} & 0 & \tilde{y} & \xi/z - \tilde{z} & 0 & 0 \end{pmatrix},$$

where $*_{51} = \tilde{z}((\alpha_3 y - \eta)/x - \alpha_2) + (y/x - \alpha_3)\tilde{\xi} - G_4(-p, -a)\tilde{x}/x$ and the matrix \mathbf{M}_1 is the hat- q version of \mathbf{L}_1 . The compatibility gives the relations:

$$\hat{\tilde{x}} = \frac{\tilde{x} - \hat{x}}{\tilde{z} - \hat{z}}, \quad \hat{\tilde{z}} = \frac{\hat{\xi} - \tilde{\xi}}{\hat{z} - \tilde{z}}, \quad x = \frac{\hat{x}\tilde{y} - \hat{y}\tilde{x}}{\tilde{x} - \hat{x}}, \quad \hat{y} = \frac{\hat{x}\tilde{z} - \hat{z}\tilde{x}}{\hat{z} - \tilde{z}}, \quad y = \frac{\hat{x}\tilde{\eta} - \hat{\eta}\tilde{x}}{\hat{x} - \tilde{x}}, \quad (5.4)$$

together with (4.8e). The second equation in (5.4) is (4.8b), while the rest four equations in (5.4) are consequences of the equations (4.8a) and (4.8b) in the GD-4 (A-2) equation. Similar to the DBSQ case [44], here we can remove the last columns and rows in \mathbf{L}_1 and \mathbf{M}_1 , and examine the remains:

$$\tilde{\phi}_2 = \mathbf{L}_2 \phi_2, \quad \hat{\phi}_2 = \mathbf{M}_2 \phi_2, \quad (5.5)$$

in which $\phi_2 = (\varphi_0, \varphi_1, \varphi_2, \varphi_3, \varphi_4)^T$ and

$$\mathbf{L}_2 = \begin{pmatrix} \tilde{z} & 0 & -1 & 0 & 0 \\ \tilde{x} & -1 & 0 & 0 & 0 \\ \tilde{\xi} & 0 & 0 & 0 & -1 \\ 0 & -\tilde{z} & \tilde{x} & 0 & 0 \\ *_{51} & G_4(-r, -a)/x & (\eta - \alpha_3 y)/x + \alpha_2 & 0 & \alpha_3 - y/x \end{pmatrix},$$

where the matrix \mathbf{M}_2 is the hat- q version of \mathbf{L}_2 . The compatibility of (5.5) gives (5.4) and (4.8e) as well.

The above are Lax pairs of the GD-4 (A-2) equation. Similarly, we can get Lax pairs for other GD-4 type lattice equations we have obtained. However, all the Lax pairs are incomplete. In other words, not all the equations in the GD-4 type equations can be obtained from the compatibility of their Lax pairs.

6 Conclusions and remarks

In this paper, we have analyzed the construction of the extended lattice GD type equations with $N = 4$ by DL approach. We first reviewed the DL scheme for the extended generic lattice GD type equations, which appeared in [44]. Then we restricted ourselves to the case of $N = 4$ and showed the derivation of the lattice GD-4 type equations, including the lattice GD-4 (A-2), (B-2) and (C-3) equations. In the DL scheme, these equations appear as closed forms, each of which includes five (independent) variables. We introduced point transformations to convert them into Hietarinta's version, which are the GD-4 counterparts of the DBSQ type equations (see Section 4.2). We also obtained reduced GD-4 (A-2) equations with four variables. In addition, we investigated the evolutions and MDC property of these equations. As examples, for the GD-4 (A-2) equation we gave its Lax pairs constructed from the DL scheme as well as the MDC property. One can also see that some lower equations are embedded into the GD-4 type equations due to the coefficients α_3 and α_2 in $G_4(\omega, k) = 0$. For example, in the GD-4 (B-2) equation (4.12), equation (4.12a) and the coefficients of α_3 and α_2 in (4.12e) compose the

DBSQ (B-2) equation (see (A.1)), while the coefficient of α_2 in (4.12e) gives rise to the lattice potential KdV equation (i.e., H1 equation in the ABS list) $(x - \widehat{\widehat{x}})(\widehat{x} - \widetilde{x}) = p^2 - q^2$.

We finish the paper with the following remarks. First of all, compared with the DBSQ case ($N = 3$, see [17, 22, 44] and also Appendix A) which are all MDC, the GD-4 type lattice equations have something special. One needs to introduce more variables to obtain quadrilateral closed forms in the DL scheme. All the obtained lattice GD-4 equations allow up-left as well as down-right evolutions with staircase initial values, but not all of them allow up-right evolution. For the GD-4 (B-2) equation, we introduced the sixth equation (4.5) (or (4.12d) or (4.13) in Hietarinta's form so that the equation set gets an up-right evolution. The enlarged system (4.12) therefore contains six equations but with five variables. However, all six equations are necessary to guarantee the system to be MDC. Among the six equations in the system (4.12), equation (4.12c) is defined on the triangle $(V, \widetilde{V}, \widehat{V})$ (not defined on the edge, cf. (4.12a)), which has zero contribution in the up-right evolution, but is necessary for the relation $\widetilde{\widetilde{\xi}} = \widehat{\widehat{\xi}} = \widetilde{\widetilde{\xi}}$ to be held (see (C.2) in Appendix C). On the other side, this fact also means that, if we admit that the GD-4 (B-2) system (4.12) excluding (4.12c) is integrable, which indicates the system is MDC, then (4.12c) will be a consequence of the other five equations in system (4.12). Similar feature holds for the GD-4 (A-2) equation, its alternative form and their four-component versions. Note that the MDC property has recently been used to construct Bäcklund transformation and find one-soliton solution to the GD-4 (B-2) equation, see [35]. Secondly, it is not clear how such GD-4 lattice equations are related to the discrete KP-type equations. In [2], the lattice modified and lattice Schwarzian GD hierarchies were obtained as reductions of the discrete KP system (see [2, equations (C.7) and (C.12)], and see also [3]), and those equations are presented in quadrilateral form in terms of multi-components. For the GD-4 case, they are 3-component quad systems. It remains open that how they are related to our GD-4 (A-2) and (C-3) equations. What's more, compared with the DBSQ equations (cf. [22]), there are still many topics to be considered in order to understand the GD-4 type lattice equations. For example, first, the one-component forms of these equations need to be determined. They are supposed to live on a 4×4 (16-point) stencil, considering the $(\widetilde{\cdot}, p) \rightarrow (\widehat{\cdot}, q)$ symmetry. However, it is not easy to eliminate auxiliary variables to get an equation for a single component, (compared with the DBSQ case in [22]). We will leave this as a topic in the next step. Besides, one can choose specific measures $d\lambda_j(k)$ and integration contour Γ_j in (2.1) to get soliton solutions, which can be operated along the lines of [44]. In addition, we have not touched the τ -function, which in the case of $N = 3$ satisfies a trilinear equation (see [44] or [22]). It is undoubtedly that figuring out the τ -function of the lattice GD-4 case and its bilinear equations (cf. [18, 19, 20, 21, 45] for the ABS and DBSQ equations) is meaningful and worth considering.

A The DBSQ-type equations

The DBSQ-type equations, including (A-2), (B-2), (C-3) and (C-4) equations [17], which are related to the third-order polynomial [44]

$$G_3(\omega, k) := \omega^3 - k^3 + \alpha_2(\omega^2 - k^2) + \alpha_1(\omega - k),$$

are listed as follows.

(A-2) equation: This equation has two forms. One is

$$\widetilde{y} = z\widetilde{x} - x, \quad \widehat{y} = z\widehat{x} - x, \quad y = x\widehat{z} - b_0x - (G_3(-p, -a)\widetilde{x} - G_3(-q, -a)\widehat{x})/(\widehat{z} - \widetilde{z}),$$

and the other is

$$y = \widetilde{z}x - \widetilde{x}, \quad y = \widehat{z}x - \widehat{x}, \quad \widehat{y} = \widehat{x}\widehat{z} - b'_0\widehat{x} - (G_3(-p, -b)\widehat{x} - G_3(-q, -b)\widetilde{x})/(\widehat{z} - \widetilde{z}).$$

Here b_0 and b'_0 are constants relating to parameter α_2 . Both of these two equations are associated with each other by reversal symmetry

$$p \rightarrow -p, \quad q \rightarrow -q, \quad n \rightarrow -n, \quad m \rightarrow -m, \quad \alpha_2 \rightarrow -\alpha_2, \quad b \rightarrow -a.$$

(B-2) equation: This equation reads

$$\tilde{z} = x\tilde{x} - y, \quad \hat{z} = x\hat{x} - y, \quad (\text{A.1a})$$

$$z = x\hat{x} - \hat{y} - \alpha_2(\hat{x} - x) - \alpha_1 - G_3(-p, -q)/(\hat{x} - \tilde{x}). \quad (\text{A.1b})$$

The parameter α_1 can be removed by transformation.

(C-3) equation: One equation of this type is

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (\text{A.2a})$$

$$y\hat{z} = z(G_3(-p, -b)\tilde{z}\tilde{y} - G_3(-q, -b)\tilde{z}\hat{y})/(\hat{z} - \tilde{z}) + G_3(-a, -b)\hat{x}, \quad (\text{A.2b})$$

and the other is

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (\text{A.3a})$$

$$y\hat{z} = z(G_3(-p, -a)\tilde{z}\tilde{y} - G_3(-q, -a)\tilde{z}\hat{y})/(\hat{z} - \tilde{z}) + G_3(-a, -b)x. \quad (\text{A.3b})$$

The equations (A.2) and (A.3) are related by reversal symmetry

$$n \rightarrow -n, \quad m \rightarrow -m \quad y \rightarrow z, \quad z \rightarrow -y, \quad a \leftrightarrow b.$$

(C-4) equation: This equation arises from the combination of equations (A.2) and (A.3)

$$x - \tilde{x} = \tilde{y}z, \quad x - \hat{x} = \hat{y}z, \quad (\text{A.4a})$$

$$y\hat{z} = z(P'_{a,b}\tilde{z}\tilde{y} - Q'_{a,b}\tilde{z}\hat{y})/(\hat{z} - \tilde{z}) + x\hat{x} - G'^2_{a,b}, \quad (\text{A.4b})$$

where

$$\begin{aligned} P'_{a,b} &= (G_3(-p, -b) + G_3(-p, -a))/2, \quad Q'_{a,b} = (G_3(-q, -b) + G_3(-q, -a))/2, \quad G'_{a,b} \\ &= G_3(-a, -b)/2. \end{aligned}$$

B Derivation of some equations in Section 3.1

B.1 Equation (3.6b)

For the both sides of equation (3.1b), we multiply from the left by ${}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}$ and from the right by $(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e}$. For the left-hand side, we have

$${}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}(A_3(-p) + A_2(-p){}^t\mathbf{\Lambda} + A_1(-p){}^t\mathbf{\Lambda}^2 + {}^t\mathbf{\Lambda}^3)(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e}. \quad (\text{B.1})$$

Calculating it term by term and express the results in terms of the variables introduced in (3.3), we have

$$\begin{aligned} A_3(-p){}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e} &= A_3(-p)s_{a,b}, \\ A_2(-p){}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}{}^t\mathbf{\Lambda}(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e} \\ &= A_2(-p){}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}({}^t\mathbf{\Lambda} - b + b)(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e} \\ &= A_2(-p)[{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}\mathbf{e} + b{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}(-b + {}^t\mathbf{\Lambda})^{-1}\mathbf{e}] \end{aligned}$$

$$\begin{aligned}
&= -A_2(-p)(v_a - 1) + bA_2(-p)s_{a,b}, \\
A_1(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} {}^t\mathbf{\Lambda}^2 (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= A_1(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} ({}^t\mathbf{\Lambda}^2 - b^2 + b^2) (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= A_1(-p) [{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} ({}^t\mathbf{\Lambda} + b) \mathbf{e} + b^2 {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e}] \\
&= A_1(-p) [(a - s_a) - b(v_a - 1)] + b^2 A_1(-p) s_{a,b},
\end{aligned}$$

and

$$\begin{aligned}
&{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} {}^t\mathbf{\Lambda}^3 (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} ({}^t\mathbf{\Lambda}^3 - b^3 + b^3) (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} ({}^t\mathbf{\Lambda}^2 + b {}^t\mathbf{\Lambda} + b^2) \mathbf{e} + b^3 {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= a^2 - r_a + b(a - s_a) - b^2(v_a - 1) + b^3 s_{a,b}.
\end{aligned}$$

Thus, equation (B.1) yields

$$\begin{aligned}
&[A_2(-p) + (a + b)A_1(-p) + a^2 + ab + b^2] + p_b s_{a,b} \\
&\quad - A_2(-p)v_a - A_1(-p)(s_a + bv_a) - r_a - bs_a - b^2 v_a,
\end{aligned}$$

where p_b is defined as in (3.5).

For the right-hand side, we can calculate

$${}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} (A_3(-p) - A_2(-p)\mathbf{\Lambda} + A_1(-p)\mathbf{\Lambda}^2 - \mathbf{\Lambda}^3) \tilde{\mathbf{U}} (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e}$$

in a similar way, which gives rise to

$$\begin{aligned}
&-[A_2(-p) + (a + b)A_1(-p) + a^2 + ab + b^2] + p_a \tilde{s}_{a,b} \\
&\quad + A_2(-p) \tilde{w}_b - A_1(-p) (\tilde{t}_b - a \tilde{w}_b) + \tilde{z}_b - a \tilde{t}_b + a^2 \tilde{w}_b,
\end{aligned}$$

where p_a is defined as in (3.5). For those terms in (3.1b) where \mathbf{O} is involved, for example, $-\mathbf{U} {}^t\mathbf{\Lambda} \mathbf{O} \mathbf{\Lambda} \tilde{\mathbf{U}}$, we make use of $\mathbf{O} = \mathbf{e} {}^t\mathbf{e}$ and have

$$\begin{aligned}
&-{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} {}^t\mathbf{\Lambda} \mathbf{O} \mathbf{\Lambda} \tilde{\mathbf{U}} (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} \\
&= -{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \mathbf{U} {}^t\mathbf{\Lambda} \mathbf{e} \cdot {}^t\mathbf{e} \tilde{\mathbf{U}} (-b + {}^t\mathbf{\Lambda})^{-1} \mathbf{e} = -(s_a - a)(1 - \tilde{w}_b).
\end{aligned}$$

We can combine all these results together. Finally, we arrive at equation (3.6b).

B.2 Equation (3.9a)

For the both sides of equation (3.1a), multiplying from the left by ${}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}$ and from the right by ${}^t\mathbf{\Lambda}^2 \mathbf{e}$, we have

$$\begin{aligned}
&{}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \tilde{\mathbf{U}} (p - {}^t\mathbf{\Lambda}) {}^t\mathbf{\Lambda}^2 \mathbf{e} \\
&= {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} (p + \mathbf{\Lambda}) \mathbf{U} {}^t\mathbf{\Lambda}^2 \mathbf{e} - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \tilde{\mathbf{U}} \mathbf{O} \mathbf{U} {}^t\mathbf{\Lambda}^2 \mathbf{e}.
\end{aligned} \tag{B.2}$$

The left-hand side of the above equation gives rise to

$$\begin{aligned}
(\text{B.2})|_{\text{l.h.s}} &= p {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \tilde{\mathbf{U}} {}^t\mathbf{\Lambda}^2 \mathbf{e} - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1} \tilde{\mathbf{U}} {}^t\mathbf{\Lambda}^3 \mathbf{e} \\
&= pa^2 - p\tilde{r}_a + \tilde{f}_a - a^3,
\end{aligned}$$

and the right-hand side yields

$$\begin{aligned}
(\text{B.2})|_{\text{r.h.s}} &= {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}(p - a + a + \mathbf{\Lambda})\mathbf{U} {}^t\mathbf{\Lambda}^2\mathbf{e} - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\tilde{\mathbf{U}}\mathbf{e} \cdot {}^t\mathbf{e}\mathbf{U} {}^t\mathbf{\Lambda}^2\mathbf{e} \\
&= (p - a) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U} {}^t\mathbf{\Lambda}^2\mathbf{e} + {}^t\mathbf{e}\mathbf{U} {}^t\mathbf{\Lambda}^2\mathbf{e} + (\tilde{v}_a - 1)u_{0,2} \\
&= (p - a)(a^2 - r_a) + u_{0,2} + (\tilde{v}_a - 1)u_{0,2} \\
&= (p - a)(a^2 - r_a) + \tilde{v}_a u_{0,2}.
\end{aligned}$$

Equation (3.9a) follows from the combination of the above results.

B.3 Equation (3.10a)

To obtain equation (3.10a), we multiply (3.1b) from the left by ${}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}$ and from the right by ${}^t\mathbf{\Lambda}^2\mathbf{e}$. The left-hand side yields

$$\begin{aligned}
& {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}(A_3(-p) + A_2(-p) {}^t\mathbf{\Lambda} + A_1(-p) {}^t\mathbf{\Lambda}^2 + {}^t\mathbf{\Lambda}^3)\mathbf{e} \\
&= A_3(-p)(1 - v_a) + A_2(-p)(a - s_a) + A_1(-p)(a^2 - r_a) + (a^3 - f_a) \\
&= p_a - A_3(-a)v_a - A_2(-p)s_a - A_1(-p)r_a - f_a,
\end{aligned}$$

where p_a is defined as in (3.5). The right-hand side yields

$$\begin{aligned}
& {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}[(A_3(-p) - A_2(-p)\mathbf{\Lambda} + A_1(-p)\mathbf{\Lambda}^2 - \mathbf{\Lambda}^3)\tilde{\mathbf{U}} \\
&\quad + \mathbf{U}[A_2(-p)\mathbf{O} - A_1(-p)(\mathbf{O}\mathbf{\Lambda} - {}^t\mathbf{\Lambda}\mathbf{O}) + (\mathbf{O}\mathbf{\Lambda}^2 - {}^t\mathbf{\Lambda}\mathbf{O}\mathbf{\Lambda} + {}^t\mathbf{\Lambda}^2\mathbf{O})]\tilde{\mathbf{U}}]\mathbf{e} \\
&= A_3(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\tilde{\mathbf{U}}\mathbf{e} - A_2(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}(\mathbf{\Lambda} + a - a)\tilde{\mathbf{U}}\mathbf{e} \\
&\quad + A_1(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}(\mathbf{\Lambda}^2 - a^2 + a^2)\tilde{\mathbf{U}}\mathbf{e} + {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}(\mathbf{\Lambda}^3 + a^3 - a^3)\tilde{\mathbf{U}}\mathbf{e} \\
&\quad + A_2(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}\mathbf{e} \cdot {}^t\mathbf{e}\tilde{\mathbf{U}}\mathbf{e} - A_1(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}\mathbf{e} \cdot {}^t\mathbf{e}\mathbf{\Lambda}\tilde{\mathbf{U}}\mathbf{e} \\
&\quad + A_1(-p) {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U} {}^t\mathbf{\Lambda}\mathbf{e} \cdot {}^t\mathbf{e}\tilde{\mathbf{U}}\mathbf{e} + {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U}\mathbf{e} \cdot {}^t\mathbf{e}\mathbf{\Lambda}^2\tilde{\mathbf{U}}\mathbf{e} \\
&\quad - {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U} {}^t\mathbf{\Lambda}\mathbf{e} \cdot {}^t\mathbf{e}\mathbf{\Lambda}\tilde{\mathbf{U}}\mathbf{e} + {}^t\mathbf{e}(a + \mathbf{\Lambda})^{-1}\mathbf{U} {}^t\mathbf{\Lambda}^2\mathbf{e} \cdot {}^t\mathbf{e}\tilde{\mathbf{U}}\mathbf{e} \\
&= A_3(-p)(1 - \tilde{v}_a) - A_2(-p)\tilde{u}_0 + aA_2(-p)(1 - \tilde{v}_a) + A_1(-p)\tilde{u}_{1,0} \\
&\quad - aA_1(-p)\tilde{u}_0 + a^2A_1(-p)(1 - \tilde{v}_a) - \tilde{u}_{2,0} + a\tilde{u}_{1,0} - a^2\tilde{u}_0 \\
&\quad + a^3(1 - \tilde{v}_a) + A_2(-p)(1 - v_a)\tilde{u}_0 - A_1(-p)(1 - v_a)\tilde{u}_{1,0} \\
&\quad + A_1(-p)(a - s_a)\tilde{u}_0 + (1 - v_a)\tilde{u}_{2,0} - (a - s_a)\tilde{u}_{1,0} + (a^2 - r_a)\tilde{u}_0 \\
&= p_a - p_a\tilde{v}_a - v_a(A_2(-p)\tilde{u}_0 - A_1(-p)\tilde{u}_{1,0} + \tilde{u}_{2,0}) - s_a(A_1(-p)\tilde{u}_0 - \tilde{u}_{1,0}) - r_a\tilde{u}_0.
\end{aligned}$$

Combining them together we get (3.10a).

C MDC property of the GD-4 (A-2) and (B-2) equations

We present procedure of checking MDC property of the alternative GD-4 (A-2) equation (4.10). Since the checking is straightforward, for the GD-4 (B-2) equation (4.12), the four-component GD-4 (A-2) equation (4.22) and its alternative form (4.23), we only list out the triple shifts of the involved variables.

Alternative GD-4 (A-2) equation (4.10): We have

$$\tilde{\tilde{x}} = (\tilde{x}\tilde{y} - \tilde{x}\tilde{y})/(\tilde{x} - \tilde{x}), \quad \tilde{\tilde{y}} = (\tilde{x}\tilde{\eta} - \tilde{x}\tilde{\eta})/(\tilde{x} - \tilde{x}), \quad \tilde{\tilde{z}} = (\tilde{y} - \tilde{y})/(\tilde{x} - \tilde{x}), \quad (\text{C.1a})$$

$$\tilde{\tilde{x}} = (\hat{x}\hat{y} - \hat{x}\hat{y})/(\hat{x} - \hat{x}), \quad \tilde{\tilde{y}} = (\hat{x}\hat{\eta} - \hat{x}\hat{\eta})/(\hat{x} - \hat{x}), \quad \tilde{\tilde{z}} = (\hat{y} - \hat{y})/(\hat{x} - \hat{x}), \quad (\text{C.1b})$$

$$\tilde{\tilde{\xi}} = \frac{\eta\tilde{x} - y\tilde{y} - (\eta + \tilde{y})\tilde{x} + (y + \tilde{x})\tilde{y}}{x(\tilde{x} - \tilde{x})}, \quad \frac{\eta\hat{x} - y\hat{y} - (\eta + \hat{y})\hat{x} + (y + \hat{x})\hat{y}}{x(\hat{x} - \hat{x})}, \quad (\text{C.1c})$$

$$\bar{\bar{\eta}} = (\alpha_3 z - \xi - \alpha_2) \frac{\tilde{x}\bar{y} - \bar{x}\tilde{y}}{\bar{x} - \tilde{x}} + (z - \alpha_3) \frac{\tilde{x}\bar{\eta} - \bar{x}\tilde{\eta}}{\bar{x} - \tilde{x}} - \frac{G_4(-p, -b)\bar{x} - G_4(-r, -b)\tilde{x}}{\bar{z} - \tilde{z}}, \quad (\text{C.1d})$$

$$\bar{\bar{\eta}} = (\alpha_3 z - \xi - \alpha_2) \frac{\hat{x}\bar{y} - \bar{x}\hat{y}}{\bar{x} - \hat{x}} + (z - \alpha_3) \frac{\hat{x}\bar{\eta} - \bar{x}\hat{\eta}}{\bar{x} - \hat{x}} - \frac{G_4(-q, -b)\bar{x} - G_4(-r, -b)\hat{x}}{\bar{z} - \hat{z}}. \quad (\text{C.1e})$$

Furthermore, from the three right-hand side equations in (4.25) and using the data (C.1) we have three different ways to calculate the values $\bar{\bar{x}}$, $\bar{\bar{y}}$, $\bar{\bar{z}}$ and $\bar{\bar{\xi}}$ uniquely, which are

$$\begin{aligned} \bar{\bar{x}} &= [(\hat{x}\bar{\eta} - \bar{x}\hat{\eta})\tilde{y} + (\bar{x}\tilde{\eta} - \tilde{x}\bar{\eta})\hat{y} + (\tilde{x}\hat{\eta} - \hat{x}\tilde{\eta})\bar{y}] / [(\hat{x} - \bar{x})\tilde{y} + (\bar{x} - \tilde{x})\hat{y} + (\tilde{x} - \hat{x})\bar{y}], \\ \bar{\bar{z}} &= [(\hat{x} - \tilde{x})\bar{\eta} + (\tilde{x} - \bar{x})\hat{\eta} + (\bar{x} - \hat{x})\tilde{\eta}] / [(\tilde{y} - \hat{y})\bar{x} + (\hat{y} - \bar{y})\tilde{x} + (\bar{y} - \tilde{y})\hat{x}], \\ \bar{\bar{\xi}} &= B_1 \times \left[\frac{(\tilde{x}\bar{y} - \bar{x}\tilde{y})\bar{\eta}}{\bar{x} - \tilde{x}} - \frac{(\bar{x}\tilde{\eta} - \tilde{x}\bar{\eta})\bar{y}}{\tilde{x} - \bar{x}} - \frac{(\bar{\eta} + \frac{\tilde{x}\bar{\eta} - \bar{x}\tilde{\eta}}{\bar{x} - \tilde{x}})(\hat{x}\bar{y} - \bar{x}\hat{y})}{\bar{x} - \hat{x}} + \frac{(\bar{y} + \frac{\tilde{x}\bar{y} - \bar{x}\tilde{y}}{\bar{x} - \tilde{x}})(\bar{x}\hat{\eta} - \hat{x}\bar{\eta})}{\hat{x} - \bar{x}} \right], \\ \bar{\bar{y}} &= \frac{B_2 + B_3 + x((\bar{y} - \tilde{y})y + (\tilde{y}\bar{z} - \bar{y}\tilde{z})x)G_4(-q, -b) + x((\tilde{y} - \hat{y})y + (\hat{y}\bar{z} - \tilde{y}\hat{z})x)G_4(-r, -b)}{x((\hat{z} - \tilde{z})\bar{y} + (\tilde{z} - \bar{z})\hat{y} + (\bar{z} - \hat{z})\tilde{y})}, \end{aligned}$$

where

$$\begin{aligned} B_1 &= \frac{1}{\left(\frac{\tilde{x}\bar{y} - \bar{x}\tilde{y}}{\bar{x} - \tilde{x}} + \frac{\hat{x}\bar{y} - \bar{x}\hat{y}}{\bar{x} - \hat{x}}\right)\bar{x}}, & B_2 &= x((\hat{y} - \bar{y})y + (\bar{y}\hat{z} - \hat{y}\bar{z})x)G_4(-p, -b), \\ B_3 &= (\alpha_3 - z)[\bar{\eta}((\tilde{y} - \hat{y})y + (\hat{y}\bar{z} - \tilde{y}\hat{z})x) + \hat{\eta}((\tilde{y} - \bar{y})y + (\bar{y}\tilde{z} - \tilde{y}\bar{z})x) \\ &\quad + \tilde{\eta}((\bar{y} - \hat{y})y + (\hat{y}\bar{z} - \bar{y}\hat{z})x)]. \end{aligned}$$

For the variable η , the formula for $\bar{\bar{\eta}}$ is somewhat long to be listed here. One can check the coincide relation $\bar{\bar{\eta}} = \hat{\bar{\eta}} = \tilde{\bar{\eta}}$ by means of mathematical softwares, e.g., MATHEMATICA, which holds only if (4.10c) and its $(\bar{\cdot}, \hat{\cdot})$ and $(\bar{\cdot}, \tilde{\cdot})$ versions hold, i.e.,

$$\hat{\bar{\xi}} - \tilde{\bar{\xi}} = (\hat{z} - \tilde{z})z, \quad \bar{\bar{\xi}} - \tilde{\bar{\xi}} = (\bar{z} - \tilde{z})z, \quad \hat{\bar{\xi}} - \bar{\bar{\xi}} = (\hat{z} - \bar{z})z.$$

In conclusion, the alternative GD-4 (A-2) equation (4.10) as a system is MDC; although (4.10c) is not needed for determining the up-right evolution of (4.10), it is necessary for the whole system (4.10) to be MDC.

GD-4 (B-2) equation (4.12): We have unique expressions for the values $\bar{\bar{x}}$, $\bar{\bar{z}}$, $\bar{\bar{\eta}}$ and $\bar{\bar{y}}$, which are

$$\begin{aligned} \bar{\bar{x}} &= [(\hat{x} - \tilde{x})\bar{\xi} + (\tilde{x} - \bar{x})\hat{\xi} + (\bar{x} - \hat{x})\tilde{\xi}] / [(\tilde{y} - \hat{y})\bar{x} + (\hat{y} - \bar{y})\tilde{x} + (\bar{y} - \tilde{y})\hat{x}], \\ \bar{\bar{z}} &= [(\hat{y} - \tilde{y})\bar{\xi} + (\tilde{y} - \bar{y})\hat{\xi} + (\bar{y} - \hat{y})\tilde{\xi}] / [(\tilde{y} - \hat{y})\bar{x} + (\hat{y} - \bar{y})\tilde{x} + (\bar{y} - \tilde{y})\hat{x}], \\ \bar{\bar{\eta}} &= \frac{(\hat{x}\bar{\xi} - \bar{x}\hat{\xi})\tilde{y} + (\bar{x}\tilde{\xi} - \tilde{x}\bar{\xi})\hat{y} + (\tilde{x}\hat{\xi} - \hat{x}\tilde{\xi})\bar{y}}{(\hat{x} - \bar{x})\tilde{y} + (\bar{x} - \tilde{x})\hat{y} + (\tilde{x} - \hat{x})\bar{y}}, \\ \bar{\bar{y}} &= \frac{1}{\frac{\tilde{y} - \bar{y}}{\bar{x} - \tilde{x}} + \frac{\hat{y} - \bar{y}}{\bar{x} - \hat{x}}} \left[\frac{G_4(-p, -r) + (x - \alpha_3)(\bar{\xi} - \tilde{\xi}) - (\tilde{y} - \bar{y})(\alpha_2 - \alpha_3 x + z)}{\tilde{x} - \bar{x}} \right. \\ &\quad \left. + \frac{-G_4(-q, -r) + (-x + \alpha_3)(\hat{\xi} - \bar{\xi}) + (\hat{y} - \bar{y})(\alpha_2 - \alpha_3 x + z)}{\hat{x} - \bar{x}} \right]. \end{aligned}$$

The formula for $\bar{\bar{\xi}}$ is too long to be listed here. The coincident relation $\bar{\bar{\xi}} = \hat{\bar{\xi}} = \tilde{\bar{\xi}}$ holds only if (4.12c) and its $(\bar{\cdot}, \hat{\cdot})$ and $(\bar{\cdot}, \tilde{\cdot})$ versions hold, i.e.,

$$\hat{\bar{\eta}} - \tilde{\bar{\eta}} = (\hat{x} - \tilde{x})z, \quad \bar{\bar{\eta}} - \tilde{\bar{\eta}} = (\bar{x} - \tilde{x})z, \quad \hat{\bar{\eta}} - \bar{\bar{\eta}} = (\hat{x} - \bar{x})z. \quad (\text{C.2})$$

Thus, the GD-4 (B-2) equation (4.12) as a system is MDC; (4.12c) is not needed for determining the up-right evolution of (4.12), but it is necessary for the whole system (4.12) to be MDC.

Four-component GD-4 (A-2) equation (4.22): \widetilde{u} , \widetilde{z} and \widetilde{v} are uniquely expressed as

$$\begin{aligned}\widetilde{u} &= [(\widetilde{\xi} - \widehat{\xi})z\bar{u} + \widetilde{u}(-\bar{u}\widetilde{\xi} + \widehat{u}(\widetilde{\xi} - \widehat{\xi}) + \bar{u}\bar{\xi} + \widehat{\xi}z - \widetilde{\xi}z) + \widehat{u}(\bar{u}\widehat{\xi} - \bar{u}\bar{\xi} - \widetilde{\xi}z + \widetilde{\xi}z)] / \\ &\quad / [\widetilde{u}\widehat{u}\widetilde{z} - \widetilde{u}\bar{u}\widetilde{z} - \widehat{u}z\widetilde{z} + \bar{u}z\widetilde{z} - \widetilde{u}\widehat{u}\widehat{z} + \widehat{u}\bar{u}\widehat{z} + \widetilde{u}z\widehat{z} - \bar{u}z\widehat{z} + (\widetilde{u} - \widehat{u})(\bar{u} - z)\bar{z}], \\ \widetilde{z} &= \frac{1}{\frac{\widetilde{\xi} - \widehat{\xi}}{\widetilde{z} - \bar{z}} + \frac{\widehat{\xi} - \widetilde{\xi}}{\widehat{z} - \bar{z}}} \left[G_4(-q, -a)/(\widehat{u} - z)(\widehat{z} - \bar{z}) - G_4(-p, -a)/(\widetilde{u} - z)(\widetilde{z} - \bar{z}) \right. \\ &\quad \left. + \frac{(\alpha_3 - u)(\bar{u} - z)(-\bar{\xi}\widetilde{z} + (\bar{\xi} - \widetilde{\xi})\widehat{z} + \widehat{\xi}(\widetilde{z} - \bar{z}) + \widetilde{\xi}\bar{z}) + (\widetilde{z} - \widehat{z})G_4(-r, -a)}{(\bar{u} - z)(\widetilde{z} - \bar{z})(\widehat{z} - \bar{z})} \right], \\ \widetilde{v} &= \frac{z((\widehat{z} - \widetilde{z})\bar{u}\bar{\xi} + (\bar{z} - \widehat{\xi})z\widetilde{z} + (\widetilde{z} - \bar{z})z\widehat{z} + (\widetilde{\xi} - \widehat{\xi})(\bar{u} - z)\bar{z}) + C_1 + C_2}{(\widehat{u} - \bar{u})\widetilde{u}\widetilde{z} + (\bar{u} - \widehat{u})z\widetilde{z} + (\bar{u} - \widehat{u})\widehat{u}\widehat{z} + (\widetilde{u} - \bar{u})z\widehat{z} + (\widetilde{u} - \widehat{u})(\bar{u} - z)\bar{z}},\end{aligned}$$

where

$$\begin{aligned}C_1 &= \widetilde{u}[(\bar{u}\bar{\xi} - \widehat{u}\widehat{\xi})\widetilde{z} + (z - \bar{u})\widetilde{\xi}\widetilde{z} - (z - \widehat{u})\widetilde{\xi}\widehat{z} + (\widetilde{\xi} - \bar{\xi})z\widetilde{z}], \\ C_2 &= \widehat{u}[-\bar{u}\bar{\xi}\widehat{z} + (-\widetilde{\xi} + \bar{\xi})z\widehat{z} + \widehat{\xi}(\widetilde{z} - \bar{z})z + \bar{u}\bar{z}].\end{aligned}$$

For the variable ξ , the formula for $\widehat{\xi}$ is somewhat long to be listed here. The coincident relation $\widetilde{\xi} = \widehat{\xi} = \widetilde{\xi}$ holds only if (4.22b) and its $(-, \widehat{\cdot})$ and $(-, \widetilde{\cdot})$ versions hold, i.e.,

$$\widetilde{v} - \widehat{v} = u(\widetilde{u} - \widehat{u}), \quad \widetilde{v} - \bar{v} = u(\widetilde{u} - \bar{u}), \quad \widehat{v} - \bar{v} = u(\widehat{u} - \bar{u}).$$

This indicates that the four-component GD-4 (A-2) equation (4.22) as a system is MDC; (4.22b) is not needed for determining the up-right evolution of (4.22), but it is necessary for the whole system (4.22) to be MDC.

Alternative four-component GD-4 (A-2) equation (4.23): For z, w and ξ , their triple shifts are uniquely expressed as

$$\begin{aligned}\widetilde{z} &= \frac{\bar{\omega}(\widetilde{z} - \widehat{z})(w - \bar{z}) - \widehat{\omega}(w - \widehat{z})(\widetilde{z} - \bar{z}) + \widetilde{\omega}(w - \widetilde{z})(\widehat{z} - \bar{z})}{(\widetilde{z} - \bar{z})\widehat{w}\widehat{z} + (\widehat{z} - \widetilde{z})\bar{w}\bar{z} + (\bar{z} - \widehat{z})\widetilde{w}\widetilde{z} + w(\widetilde{z} - \bar{z})\bar{w} + (\widehat{z} - \bar{z})\widetilde{w} + (\bar{z} - \widetilde{z})\widehat{w}}, \\ \widetilde{w} &= \frac{1}{\frac{\widetilde{\omega} - \widehat{\omega}}{\widetilde{w} - \bar{w}} + \frac{\widehat{\omega} - \widetilde{\omega}}{\widehat{w} - \bar{w}}} \left[\frac{1}{\bar{w} - \widehat{w}} \left(\frac{(\widetilde{w} - \widehat{w})\bar{\omega} + (\widehat{w} - \bar{w})\widetilde{\omega} + \widehat{\omega}(\bar{w} - \widetilde{w})(\alpha_3 - z)}{\widetilde{w} - \bar{w}} - \frac{G_4(-q, -b)}{w - \widehat{z}} \right) \right. \\ &\quad \left. + \frac{1}{\widetilde{w} - \bar{w}} \left(\frac{(\widehat{w} - \widetilde{w})G_4(-r, -b)}{(\widehat{w} - \bar{w})(w - \bar{z})} - \frac{G_4(-p, -b)}{w - \widetilde{z}} \right) \right], \\ \widetilde{\xi} &= \frac{-\bar{\omega}(w\widetilde{z} - w\widehat{z} - \widetilde{w}\widetilde{z} + \widehat{w}\widehat{z})(w - \bar{z}) + \widetilde{\omega}(w - \widetilde{z})(-w\widehat{w} + w\bar{w} + \widehat{w}\widehat{z} - \bar{w}\bar{z}) + C_3}{\widehat{w}\widetilde{z}\widehat{z} - \bar{w}\widetilde{z}\bar{z} - \widehat{w}\widehat{z}\widehat{z} + \bar{w}\widehat{z}\bar{z} + \widetilde{w}\widetilde{z}(\widetilde{z} - \widehat{z}) + w(\bar{w}\widetilde{z} + \widetilde{w}\widehat{z} - \bar{w}\widehat{z} - \widetilde{w}\bar{z} + \widehat{w}(\bar{z} - \widetilde{z}))},\end{aligned}$$

where

$$C_3 = \bar{\omega}(w - \widehat{z})(w(\widetilde{w} - \bar{w}) - \widetilde{w}\widetilde{z} + \bar{w}\bar{z}).$$

For the variable ω , the formula for $\widetilde{\omega}$ is long and one can check the coincident relation $\widetilde{\omega} = \widehat{\omega} = \widetilde{\omega}$ holds only if (4.23c) and its $(-, \widehat{\cdot})$ and $(-, \widetilde{\cdot})$ versions hold, i.e.,

$$\widehat{\xi} - \widetilde{\xi} = (\widehat{z} - \widetilde{z})z, \quad \bar{\xi} - \widetilde{\xi} = (\bar{z} - \widetilde{z})z, \quad \widehat{\xi} - \bar{\xi} = (\widehat{z} - \bar{z})z.$$

In conclusion, the four-component alternative GD-4 (A-2) equation (4.23) as a system is MDC; although (4.23c) is not needed for determining the up-right evolution of (4.23), it is necessary for the whole system (4.23) to be MDC.

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