# Uniform Asymptotic Expansion for the Incomplete Beta Function 

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#### Abstract

In [Temme N.M., Special functions. An introduction to the classical functions of mathematical physics, A Wiley-Interscience Publication, John Wiley \& Sons, Inc., New York, 1996, Section 11.3.3.1] a uniform asymptotic expansion for the incomplete beta function was derived. It was not obvious from those results that the expansion is actually an asymptotic expansion. We derive a remainder estimate that clearly shows that the result indeed has an asymptotic property, and we also give a recurrence relation for the coefficients.


Key words: incomplete beta function; uniform asymptotic expansion
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## 1 Introduction

For positive real numbers $a, b$ and $x \in[0,1]$, the (normalised) incomplete beta function $I_{x}(a, b)$ is defined by

$$
I_{x}(a, b)=\frac{1}{B(a, b)} \int_{0}^{x} t^{a-1}(1-t)^{b-1} \mathrm{~d} t
$$

where $B(a, b)$ denotes the ordinary beta function:

$$
B(a, b)=\int_{0}^{1} t^{a-1}(1-t)^{b-1} \mathrm{~d} t=\frac{\Gamma(a) \Gamma(b)}{\Gamma(a+b)}
$$

(see, e.g., [2, Section 8.17(i)]). In this paper, we will use the notation of [2, Section 8.18(ii)].
The incomplete beta function plays an important role in statistics in connection with the beta distribution (see, for instance, [1, pp. 210-275]). Large parameter asymptotic approximations are useful in these applications. For fixed $x$ and $b$, one could use the asymptotic expansion

$$
I_{x}(a, b)=\frac{x^{a}(1-x)^{b-1}}{a B(a, b)}{ }_{2} F_{1}\left(\begin{array}{c}
1,1-b  \tag{1}\\
a+1
\end{array} ; \frac{x}{x-1}\right) \sim \frac{x^{a}(1-x)^{b-1}}{a B(a, b)} \sum_{n=0}^{\infty} \frac{(1-b)_{n}}{(a+1)_{n}}\left(\frac{x}{x-1}\right)^{n}
$$

as $a \rightarrow+\infty$. The right-hand side of (1) converges only for $x \in\left[0, \frac{1}{2}\right)$, but for any fixed $x \in[0,1)$ it is still useful when used as an asymptotic expansion as $a \rightarrow+\infty$. For more details, see [3, Section 11.3.3]. However, it is readily seen that (1) breaks down as $x \rightarrow 1$. Since this limit has significant importance in applications, Temme derived in [3, Section 11.3.3.1] an asymptotic expansion as $a \rightarrow+\infty$ that holds uniformly for $x \in(0,1]$. His result can be stated as follows.

Theorem 1. Let $\xi=-\ln x$. Then for any fixed positive integer $N$ and fixed positive real $b$,

$$
\begin{equation*}
I_{x}(a, b)=\frac{\Gamma(a+b)}{\Gamma(a)}\left(\sum_{n=0}^{N-1} d_{n} F_{n}+\mathcal{O}\left(a^{-N}\right) F_{0}\right) \tag{2}
\end{equation*}
$$

as $a \rightarrow+\infty$, uniformly for $x \in(0,1]$. The functions $F_{n}=F_{n}(\xi, a, b)$ are defined by the recurrence relation

$$
\begin{equation*}
a F_{n+1}=(n+b-a \xi) F_{n}+n \xi F_{n-1} \tag{3}
\end{equation*}
$$

with

$$
F_{0}=a^{-b} Q(b, a \xi), \quad F_{1}=\frac{b-a \xi}{a} F_{0}+\frac{\xi^{b} \mathrm{e}^{-a \xi}}{a \Gamma(b)}
$$

and $Q(a, z)=\Gamma(a, z) / \Gamma(a)$ is the normalised incomplete gamma function (see [2, Section 8.2(i)]). The coefficients $d_{n}=d_{n}(\xi, b)$ are defined by the generating function

$$
\begin{equation*}
\left(\frac{1-\mathrm{e}^{-t}}{t}\right)^{b-1}=\sum_{n=0}^{\infty} d_{n}(t-\xi)^{n} \tag{4}
\end{equation*}
$$

In particular,

$$
d_{0}=\left(\frac{1-x}{\xi}\right)^{b-1}, \quad d_{1}=\frac{x \xi+x-1}{(1-x) \xi}(b-1) d_{0}
$$

They satisfy the recurrence relation

$$
\begin{align*}
\xi(n+1)(n+2) d_{0} d_{n+2}= & \xi \sum_{m=0}^{n}(m+1)\left(n-2 m+1+\frac{m-n-1}{b-1}\right) d_{m+1} d_{n-m+1} \\
& +\sum_{m=0}^{n}(m+1)\left(n-2 m-2-\xi+\frac{m-n}{b-1}\right) d_{m+1} d_{n-m} \\
& +\sum_{m=0}^{n}(1-m-b) d_{m} d_{n-m} \tag{5}
\end{align*}
$$

In the case that $b=1$, we have $d_{0}=1$ and $d_{n}=0$ for $n \geq 1$.
Our contribution is the remainder estimate in (2) and the recurrence relation (5). In fact, it is not at all obvious from (3) that the sequence $\left\{F_{n}\right\}_{n=0}^{\infty}$ has an asymptotic property as $a \rightarrow+\infty$. We will show that for any non-negative integer $n$,

$$
\begin{equation*}
0<F_{n+1} \leq \frac{n+\beta}{a} F_{n} \tag{6}
\end{equation*}
$$

where $\beta=\max (1, b)$.
In [4, Section 38.2.8] the function $F_{n}$ is identified as a Kummer $U$-function:

$$
F_{n}=\frac{\xi^{n+b} \mathrm{e}^{-a \xi} n!}{\Gamma(b)} U(n+1, n+b+1, a \xi)
$$

## 2 Proof of the main results

We proceed similarly as in [3, Section 11.3.3.1] and start with the integral representation

$$
\begin{equation*}
I_{x}(a, b)=\frac{1}{B(a, b)} \int_{\xi}^{+\infty} t^{b-1} \mathrm{e}^{-a t}\left(\frac{1-\mathrm{e}^{-t}}{t}\right)^{b-1} \mathrm{~d} t \tag{7}
\end{equation*}
$$

We substitute the truncated Taylor series expansion

$$
\left(\frac{1-\mathrm{e}^{-t}}{t}\right)^{b-1}=\sum_{n=0}^{N-1} d_{n}(t-\xi)^{n}+r_{N}(t)
$$

into (7) and obtain

$$
I_{x}(a, b)=\frac{\Gamma(a+b)}{\Gamma(a)}\left(\sum_{n=0}^{N-1} d_{n} F_{n}+R_{N}(a, b, x)\right)
$$

where $F_{n}$ is given by the integral representation

$$
\begin{equation*}
F_{n}=\frac{1}{\Gamma(b)} \int_{\xi}^{+\infty} t^{b-1} \mathrm{e}^{-a t}(t-\xi)^{n} \mathrm{~d} t=\frac{\mathrm{e}^{-a \xi}}{\Gamma(b)} \int_{0}^{+\infty}(\tau+\xi)^{b-1} \tau^{n} \mathrm{e}^{-a \tau} \mathrm{~d} \tau \tag{8}
\end{equation*}
$$

and the remainder term $R_{N}(a, b, x)$ is defined by

$$
\begin{equation*}
R_{N}(a, b, x)=\frac{1}{\Gamma(b)} \int_{\xi}^{+\infty} t^{b-1} \mathrm{e}^{-a t} r_{N}(t) \mathrm{d} t \tag{9}
\end{equation*}
$$

The recurrence relation (3) can be obtained from (8) via a simple integration by parts.
Let, for a moment,

$$
c_{n}(a, b)=\int_{0}^{+\infty}(\tau+\xi)^{b-1} \tau^{n} \mathrm{e}^{-a \tau} \mathrm{~d} \tau
$$

Then via integration by parts we find

$$
\begin{equation*}
a c_{n+1}(a, b)=(n+b) c_{n}(a, b)+\xi(1-b) c_{n}(a, b-1) . \tag{10}
\end{equation*}
$$

We make the observation that

$$
\begin{equation*}
0 \leq \xi c_{n}(a, b-1)=\xi \int_{0}^{+\infty}(\tau+\xi)^{b-2} \tau^{n} \mathrm{e}^{-a \tau} \mathrm{~d} \tau \leq c_{n}(a, b) \tag{11}
\end{equation*}
$$

It follows from (10) and (11) that

$$
a c_{n+1}(a, b) \leq \begin{cases}(n+1) c_{n}(a, b) & \text { if } 0<b \leq 1, \\ (n+b) c_{n}(a, b) & \text { if } b \geq 1\end{cases}
$$

Since $F_{n}=\mathrm{e}^{-a \xi} c_{n}(a, b) / \Gamma(b)$, this inequality implies (6).
To obtain the remainder estimate in (2), we use the Cauchy integral representation

$$
\begin{equation*}
r_{N}(t)=\frac{(t-\xi)^{N}}{2 \pi \mathrm{i}} \oint_{\{\xi, t\}} \frac{\left(\frac{1-\mathrm{e}^{-\tau}}{\tau}\right)^{b-1}}{(\tau-t)(\tau-\xi)^{N}} \mathrm{~d} \tau \tag{12}
\end{equation*}
$$

where the contour encircles the points $\xi$ and $t$ once in the positive sense. From the integral representation (9), we have that $0 \leq \xi \leq t$. Thus, in the case that $N \geq 1$, we can deform the contour in (12) to the path

$$
\begin{aligned}
{[1+\infty \mathrm{i}, 1+\pi \mathrm{i}] } & \cup[1+\pi \mathrm{i},-1+\pi \mathrm{i}] \cup[-1+\pi \mathrm{i},-1-\pi \mathrm{i}] \\
& \cup[-1-\pi \mathrm{i}, 1-\pi \mathrm{i}] \cup[1-\pi \mathrm{i}, 1-\infty \mathrm{i}]
\end{aligned}
$$

For the integrals along the final three portions of the path, we have the estimates

$$
\begin{align*}
& \left|\frac{1}{2 \pi \mathrm{i}} \int_{-1+\pi \mathrm{i}}^{-1-\pi \mathrm{i}} \frac{\left(\frac{1-\mathrm{e}^{-\tau}}{\tau}\right)^{b-1}}{(\tau-t)(\tau-\xi)^{N}} \mathrm{~d} \tau\right| \leq \frac{\max \left((\mathrm{e}-1)^{b-1},\left(\frac{\mathrm{e}+1}{\sqrt{\pi^{2}+1}}\right)^{b-1}\right)}{(1+\xi)^{N+1}}, \\
& \left|\frac{1}{2 \pi \mathrm{i}} \int_{-1-\pi \mathrm{i}}^{1-\pi \mathrm{i}} \frac{\left(\frac{1-\mathrm{e}^{-\tau}}{\tau}\right)^{b-1}}{(\tau-t)(\tau-\xi)^{N}} \mathrm{~d} \tau\right| \leq \frac{\max \left(\left(\frac{\mathrm{e}^{ \pm 1}+1}{\sqrt{\pi^{2}+1}}\right)^{b-1}\right)}{\pi^{N+2}} \tag{13}
\end{align*}
$$

and

$$
\begin{align*}
\left|\frac{1}{2 \pi \mathrm{i}} \int_{1-\pi \mathrm{i}}^{1-\infty \mathrm{i}} \frac{\left(\frac{1-\mathrm{e}^{-\tau}}{\tau}\right)^{b-1}}{(\tau-t)(\tau-\xi)^{N}} \mathrm{~d} \tau\right| & \leq \frac{1}{2 \pi} \int_{\pi}^{+\infty} \frac{\max \left(\left(1 \pm \mathrm{e}^{-1}\right)^{b-1}\right)\left(s^{2}+1\right)^{(1-b) / 2}}{\sqrt{s^{2}+(1-t)^{2}}\left(s^{2}+(1-\xi)^{2}\right)^{N / 2}} \mathrm{~d} s \\
& \leq \frac{\max \left(\left(1 \pm \mathrm{e}^{-1}\right)^{b-1}\right)}{2 \pi} \int_{\pi}^{+\infty} \frac{\left(s^{2}+1\right)^{(1-b) / 2}}{s^{N+1}} \mathrm{~d} s \tag{14}
\end{align*}
$$

respectively. The integrals along the first two portions can be estimated similarly to (13) and (14). Hence, for $0 \leq \xi \leq t$ and $N \geq 1$, we have

$$
\left|r_{N}(t)\right| \leq C_{N}(b)(t-\xi)^{N}
$$

where the constant $C_{N}(b)$ does not depend on $\xi$. Using this result in the integral representation (9), we can infer that

$$
\left|R_{N}(a, b, x)\right| \leq C_{N}(b) F_{N}
$$

Finally, combining this result with the inequalities (6), we obtain the required remainder estimate in (2).

The reader can check that the function $f(t)=\left(\frac{1-\mathrm{e}^{-t}}{t}\right)^{b-1}$ is a solution of the nonlinear differential equation

$$
t f(t) f^{\prime \prime}(t)-\frac{b-2}{b-1} t f^{\prime 2}(t)+(t+2) f(t) f^{\prime}(t)+(b-1) f^{2}(t)=0 .
$$

If we substitute the Taylor series (4) into this differential equation and rearrange the result, we obtain the recurrence relation (5).

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