C^* -Algebras Associated with Quadratic Dynamical System

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In this paper we consider enveloping C^* -algebras of *-algebras given by generators and defining relations of the following form $A = \mathbb{C}\langle X, X^* | XX^* = f(X^*X) \rangle$, where f is a Hermitian mapping. Some properties of these algebras associated with simple dynamical systems (f, \mathbb{R}) are studied. As an example quadratic dynamical systems are considered.

1 Introduction

It is well known that there is close connection between the representation theory of C^* -algebras and structure of dynamical systems (f(), X). In the case when f is one-to-one mapping, the C^* -algebra associated with the transformation group had been studied by many authors, for example by Glimm, Effros and Hahn. The general theory of cross-products of C^* -algebras was elaborated by Doplicher, Kastler and Robinson.

In recent papers (see [8] and references given there in) a special class of *-algebras given by generators and relations was considered and some of the results from the theory of crossproduct C^* -algebras were transferred into non-bijective settings, which may be important in studying of multi-dimensional non-linear deformation (see [11, 10, 3]), such as Witten's first deformation of su(2), Quesne and Beckers non-linear deformation of su(2) etc. Examples were studied in connection with different quantum deformations of algebras, such as Quantum Unit Disc (Klimek and Lesnievski), one-dimensional q-CCR and their non linear transformation, ets see [7, 8].

Thus, for example, for one-parameter Quantum Unit Disc there corresponds the dynamical system: $f(\lambda) = \frac{(q+\mu)\lambda-\mu}{\mu\lambda+1-\mu}$, where μ is a parameter of deformation, for two-parameter Quantum Unit Disc there corresponds $f(\lambda) = \frac{(q+\mu)\lambda+1-q-\mu}{\mu\lambda+1-\mu}$, for Witten's first deformation of su(2) there corresponds two-dimensional quadratic map $f(x, y) = (p^{-1}(1+p^{-1}x), g(gy-x+(p-p^{-1})x^2))$, where $g = \pm 1$ depending on the chosen real form and p is a parameter of deformation.

In the present paper we will deal with a one-dimensional polynomial map $f : \mathbb{R} \to \mathbb{R}$ and consider *-algebra $\mathcal{A}_f = \mathbb{C}\langle X, X^* | XX^* = f(X^*X) \rangle$. Under condition of simplicity of the dynamical system (f, \mathbb{R}) we prove that the enveloping C^* -algebra is GCR (type I) C^* -algebra and investigate some other properties. We also discuss the question what relation between dynamical systems (f_1, \mathbb{R}) and (f_2, \mathbb{R}) corresponds to the isomorphism of enveloping C^* -algebras of *-algebras \mathcal{A}_{f_1} and \mathcal{A}_{f_2} .

In the last section we consider an example: Unharmonical Quantum Oscillator, i.e. the twoparametric family of *-algebras $\mathcal{A}_{a,b} = \mathbb{C}\langle X, X^* | XX^* = 1 + aX^*X - b(X^*X)^2 \rangle$, where a and b are real parameters with b > 0. Some partitioning of parametric domain into parts depending on isomorphism class of C*-enveloping algebra are given.

2 Simple dynamical systems

For the convenience of the reader we repeat the relevant material from [12, 8] without proofs, thus making our exposition self-contained. By the dynamical system we mean a continuous map $f: \mathbb{R} \to \mathbb{R}$ or $f: I \to I$, where $I \subset \mathbb{R}$ is a closed bounded interval. By the orbit of dynamical system (f, \mathbb{R}) we mean a sequence $\delta = (x_k)_{k \in P}$, where P is one of the sets \mathbb{Z} , \mathbb{N} , such that $f(x_k) = x_{k+1}$. But sometimes we will consider orbit as the set $\{x_k | k \in P\}$. The set of all orbits will be denoted by $\operatorname{Orb}(f)$. For $x \in \mathbb{R}$ denote by $\mathcal{O}_+(x)$ the forward orbit, i.e. $(f^k(x))_{k\geq 0}$. For every orbit $\delta \in \operatorname{Orb}(f)$ define $\omega(\delta)$ be the set of accumulation points of forward half-orbit and $\alpha(\delta)$ be the set of accumulation points of backward half-orbit.

By the positive orbit of $(f(), \mathbb{R})$ we mean a sequence $\omega = (x_k)_{k \in \mathbb{Z}}$ such that $f(x_k) = x_{k+1}$ and $x_k > 0$ for all integer k. Unilateral positive orbit is a sequence $\omega = (x_k)_{k \in \mathbb{N}}$ (Fock-orbit) such that $x_1 = 0$ and $f(x_k) = x_{k+1}, x_k > 0$ for k > 1 or $\omega = (x_{-k})_{k \in \mathbb{N}}$ (anti-Fock-orbit) such that $x_{-1} = 0$ and $f(x_k) = x_{k+1}, x_k > 0$ for k < -1. Define $\operatorname{Orb}_+(f)$ be the set of all positive orbits which are either periodic (cycles) or contain no cycles. Note that $\omega(\delta) = \emptyset$ for any anti-Fock orbit δ and $\alpha(\delta_1) = \emptyset$ for the Fock orbit δ_1 .

Cycle $\beta = \{\beta_1, \ldots, \beta_m\}$ is called attractive if there is a neighborhood U of β such that $f(U) \subseteq U$ and $\bigcap_{i>0} f^i(U) = \beta$.

Point $x \in \mathbb{R}$ is called non-wandering if for every its neighborhood U there exists a positive integer m such that $f^m(U) \cap U \neq \emptyset$.

Since we will consider only bounded from above functions f and positive orbits we can always consider our dynamical system on closed interval $[0, \sup f]$.

In this paper we will deal with a simple dynamical system which possesses one of the equivalent properties listed in the following theorem:

Theorem 1 ([12], 3.14]). Let (f(), I) be dynamical system with $f \in C(I)$, $(I \subset \mathbb{R}$ is closed bounded interval). The following conditions are equivalent:

1) for every $x \in I$ $\omega(x) = \omega(\mathcal{O}_+(x))$ is cycle;

2) Per(f) is closed;

3) every non-wandering point is periodic.

f is called partially monotone, if I decomposes into a finite union of sub intervals, on which f is monotone.

Let us mention the following statement from [8].

Theorem 2. Let (f, I) be a dynamical system with partially monotone and continuous f. Then the following conditions are equivalent:

1) Per(f) is closed;

2) for some positive integer m the relation $Fix(f^{2^{m+1}}) = Fix(f^{2^m})$ holds;

3) any quasi-invariant ergodic measure is concentrated on a single element of the trajectory decomposition.

The class of dynamical systems which satisfies equivalent conditions 1–3 of Theorem 2 is denoted by \mathcal{F}_{2^m} . Let us note that when $\operatorname{Per}(f)$ is closed Theorem ([12], 3.12) implies that the length of every cycle is a power of 2 and there no homoclinical orbits (i.e. orbit δ such that $\alpha(\delta) = \omega(\delta)$ is a cycle).

We will need the following lemma:

Lemma 1. Let (f, \mathbb{R}) be simple dynamical system with bounded f and the set of periodic points which are not the points of attractive cycles, i.e. the set $[0, \sup f] \cap \operatorname{Per}(f) \setminus \bigcup_{\beta \text{-attractive cycle } \beta}$ be finite then for every orbit $\delta \in \operatorname{Orb}_+(f)$ the α -boundary $\alpha(\delta)$ is cycle which is not attractive. **Proof.** 1. Let us show that every α -boundary point is non-wandering: if $x \in \alpha(\delta)$ then for arbitrary $\epsilon > 0$ and positive integer n there is $y \in B_{\epsilon}(x)$ and integer $l \ge n$ such that $f^{l}(y) \in B_{\epsilon}(x)$. Indeed, if $\delta = (x_{k})_{k \in \mathbb{Z}}$ then there is subsequence $x_{-n_{k}} \to x$. For a given $\epsilon > 0$ we can find an integer k_{0} such that $x_{-n_{k}} \in B_{\epsilon}(x)$ for all $k \ge k_{0}$. Take $k_{1} > k_{0}$ such that $n_{k_{1}} \ge n_{k_{0}} + n$ and put $y = x_{-n_{k_{1}}} \in B_{\epsilon}(x)$ and $l = n_{k_{1}} - n_{k_{0}} \ge n$. Then $f^{l}(y) = x_{-n_{k_{0}}} \in B_{\epsilon}(x)$.

2. Since for a simple dynamical system every non-wandering point is periodic we obtain that $\alpha(f) = \bigcup_{\delta \in \operatorname{Orb}_+(f)} \alpha(\delta)$ is contained in $\operatorname{Per}(f) \cap [0, \sup f]$.

3. Let β be an attractive cycle and assume that $\beta \subseteq \alpha(\delta)$, then $\alpha(\delta) = \beta$. Indeed, let $\beta_1 \in \alpha(\delta)$ be another cycle, then there is $\epsilon > 0$ such that $\beta_1 \cap B_{\epsilon}(\beta) = \emptyset$. Since β is an attractive cycle there is $\eta > 0$ and $\eta < \epsilon$ such that for arbitrary $y \in B_{\eta}(\beta)$ we have $\mathcal{O}_+(y) \subseteq B_{\epsilon}(\beta)$. Let $\delta = (x_k)_{k \in \mathbb{Z}}$. Since $\beta_1 \in \alpha(\delta)$ there is a positive integer k_0 such that $x_{-k_0} \in B_{\epsilon_1}(\beta_1)$, where $\epsilon_1 > 0$ chosen such that $B_{\epsilon_1}(\beta_1) \cap B_{\epsilon}(\beta) = \emptyset$. But $\beta \in \alpha(\delta)$ so there is a positive integer $k_1 > k_0$ with the property $x_{-k_1} \in B_{\eta}(\beta)$ then $\mathcal{O}_+(x_{-k_1}) \subseteq B_{\epsilon}(\beta)$ and, obviously, $x_{-k_0} \in \mathcal{O}_+(x_{-k_1})$. This is a contradiction. Thus we have proved that $\alpha(\delta) = \beta$. But this implies that $\alpha(\delta) = \omega(\delta) = \beta$. So δ is a homoclinical orbit and so (f, \mathbb{R}) is not simple. Which is a contradiction. Thus $\alpha(\delta)$ has no attractive cycles.

4. Let us prove that $\alpha(\delta)$ is a single cycle for every orbit δ . We already know that $\alpha(\delta) \subset [0, \sup f] \cap \operatorname{Per}(f) \setminus \bigcup_{\beta \text{-attractive cycle}} \beta$. Since the last one is a finite set there is $\epsilon > 0$ such that for arbitrary distinct cycles $\beta_1, \beta_2 \in \alpha(\delta)$ we have $B_{\epsilon}(\beta_1) \cap B_{\epsilon}(\beta_2) = \emptyset$ and $f(B_{\epsilon}(\beta_1)) \cap B_{\epsilon}(\beta_2) = \emptyset$ (the last is a possible since $f(\beta_1) = \beta_1$ and so $f(\beta_1) \cap \beta_2 = \emptyset$). There is a positive integer n such that $x_{-k} \in B_{\epsilon}(\alpha(\delta))$ for every k > n. Thus for every cycle $\beta \in \alpha(\delta)$ there is positive $\epsilon_1 < \epsilon$ such that $f(y) \in B_{\epsilon_1}(\beta_1)$ for every $y \in B_{\epsilon_1}(\beta)$, where $\beta_1 \neq \beta$ (in the opposite case would be $\alpha(\delta) = \beta$). This contradicts $f(B_{\epsilon}(\beta)) \cap B_{\epsilon}(\beta_1) = \emptyset$.

3 Representation theory of *-algebras associated with \mathcal{F}_{2^m} dynamical systems

The following theorem is due to Samoilenko and Ostrovskii [8].

Theorem 3. Let f be partially monotone continuous map and (f, \mathbb{R}) be \mathcal{F}_{2^m} dynamical system. Let $A = \mathbb{C}\langle X, X^* | XX^* = f(X^*X) \rangle$ be corresponding *-algebra.

1. To every positive non-cyclic orbit $\omega(x_k)_{k\in \mathbb{Z}}$ there corresponds an irreducible representation π_{ω} in Hilbert space $l_2(\mathbb{Z})$ given by the formulae: $Ue_k = e_{k-1}$, $Ce_k = \sqrt{x_k}e_k$ for $k \in \mathbb{Z}$ and X = UC is a polar decomposition.

2. To positive Fock-orbit $\omega = (x_k)_{k \in N}$ there corresponds an irreducible representation π_{ω} in Hilbert space $l_2(N)$ given by the formulae: $Ue_0 = 0$, $Ue_k = e_{k-1}$, $Ce_k = \sqrt{x_k}e_k$ for k > 1 and X = UC.

3. To positive anti-Fock-orbit $\omega = (x_{-k})_{k \in N}$ there corresponds an irreducible representation π_{ω} in Hilbert space $l_2(N)$ given by the formulae: $Ue_k = e_{k-1}$, $Ce_k = \sqrt{x_k}e_k$ for k > 1 and X = UC.

4. To cyclic positive orbit $\omega = (x_k)_{k \in N}$ of length m there corresponds a family of m-dimensional irreducible representation $\pi_{\omega,\phi}$ in Hilbert space $l_2(\{1,\ldots,m\})$ given by the formulae: $Ue_0 = e^{i\phi}e_{m-1}, Ue_k = e_{k-1}, Ce_k = \sqrt{x_k}e_k$ for $k = 1, \ldots, m; 0 \le \phi \le 2\pi$ and X = UC.

This is a complete list of unequivalent irreducible representation of a given *-algebra.

4 Enveloping C^* -algebra

Let f be a bounded from above Hermitian polynomial (hence f is always partially monotone and continuous). Let $A_f = \mathbb{C}\langle X, X^* | XX^* = f(X^*X) \rangle$ be *-algebra given by generators and relations which has at least one representation. Let $C = \sup f$. Then for any representation π of *-algebra A_f we have $||X|| \leq \sqrt{C}$. Thus there is an enveloping C^* -algebra, which we denote by \mathcal{E}_f . Let us note that by Theorem 3.3 [12] for $f \in C^1(I, I)$ simplicity of a dynamical system is equivalent to $(f, I) \in \mathcal{F}_{2^m}$ for some integer m.

Theorem 4. Let a dynamical system (f, \mathbb{R}) be simple and $\delta \in Orb_+(f)$.

1. If δ is non-cyclic bilateral orbit than $C^*(\pi_{\delta}) = Z \times_{\delta} C(\overline{\delta})$ is a cross-product of C^* -algebra, where $\overline{\delta} = \delta \cup \omega(\delta) \cup \alpha(\delta)$.

The set of irreducible representation $\operatorname{Irr}(C^*(\pi_{\delta}))$ is $\pi_{\delta}, \pi_{\omega(\delta),\phi}, \pi_{\alpha(\delta),\phi}, \text{ where } 0 \leq \phi \leq 2\pi$.

2. Assume that 0 is not a periodic point. If δ is a Fock-orbit then $C^*(\pi_{\delta}) \cong M_m(\mathcal{T}(C(\mathbb{T})))$ is a matrix algebra of dimension $m = |\omega(\delta)|$ over C^* -algebra $\mathcal{T}(C(\mathbb{T}))$ of the Toeplitz operators.

The same is true for anti-Fock orbit with $m = |\alpha(\delta)|$.

Proof. Let $\pi \in \operatorname{Irr}(C^*(\pi_{\delta}))$ then $\sigma(\pi(C^2)) \subseteq \sigma(\pi_{\delta}(C^2)) = \delta \cup \omega(\delta) \cup \alpha(\delta)$. Since every irreducible representation of $C^*(\pi_{\delta})$ is also an irreducible representation of *-algebra \mathcal{A} there is orbit δ' such that $\pi = \pi_{\delta'}$. Then we will have $\delta' \subseteq \delta \cup \omega(\delta) \cup \alpha(\delta)$. Since no cycle can be properly contained in an orbit δ' we conclude that $\delta' = \delta$ or $\delta' = \omega(\delta)$ or $\delta' = \alpha(\delta)$. So π must be one of the representations listed in the theorem. Let us prove that all of them are actually representations of the algebra $C^*(\pi_{\delta})$. Let $A = \operatorname{diag}(\ldots, x_{-1}, x_0, x_1, \ldots)$ be diagonal operator in Hilbert space $l_2(\mathbb{Z})$ with orthonormal basis $\{e_k\}$, where $\delta = (\ldots, x_1, x_0, x_1, \ldots)$. Let U be a bilateral shift operator $Ue_k = e_{k-1}$. The equality $UAU^* = f(A)$ implies $UAU^* \in C^*(A)$. Since $B = U^*AU$ is a diagonal operator $Be_k = x_{k+1}e_k$ and the mapping $x_k \to x_{k+1}$ is mutually continuous on the closure of $\delta($ which is $\sigma(A)) B \in C^*(A)$. Thus the mapping $\rho(D) = UDU^*$ is a automorphism of $C^*(A)$. Let us prove that $C^*(\pi(\delta))$ is a cross-product C^* -algebra. Consider the linear subspace $L_t = \{\sum_{-m}^n A_i u^i e_t | A_i \in C^*(A); m, n \ge 0 \}$. Then L_t is dense in H and

 L_t is isomorphic to a dense subspace in $L_2(\mathbb{Z}, \mathbb{C})$ via isomorphism $\sum A_i u^i e_t \to f()$, where $f(i) = (A_i e_{t+i}, e_{t+i})$. Direct computations show that $A(\sum \alpha_i e_i) = \sum \alpha_i \phi_t (U^{-i}AU^{*-i})$, where $\phi_t(D) = (De_t, e_t)$ for all $D \in C^*(A)$ is a one-dimensional representation of $C^*(A)$. Thus π_{δ} is a regular representation, λ_{ϕ} , associated with the representation ϕ_t . Since $\phi = \bigoplus_{t \in \mathbb{Z}} \phi_t$ is a faithful representation of $C^*(A)$ we conclude that λ_{ϕ} is faithful on the cross-product $\mathbb{Z} \times_{\rho} C^*(A)$ (see [9], Theorems 7.7.5 and 7.7.7). Since all representations λ_{ϕ_t} are isomorphic to π_{δ} we conclude that $C^*(\pi_{\delta})$ is $\mathbb{Z} \times_{\rho} C^*(A)$.

Consider the case of unilateral orbits. Since point 0 is not periodic and the dynamical system is simple we conclude that for every orbit δ there is $\eta > 0$ such that $\delta \subseteq [\eta, \sup f]$, i.e. δ is separeted from zero. Let $\delta = (x_k)_{k \in \mathbb{N}}$ be the Fock orbit. Then $\pi_{\delta}(X)$ is a weighted shift operator with all weights separated from zero. If X = UC is polar decomposition of X then $U, C \in C^*(X)$ and the algebra of compact operators $\mathcal{K} \subseteq C^*(X)$ (see [1], Lemma 2.1). We know that $\omega(\delta) = (y_k)_{k \in \mathbb{N}}$ is periodic orbit and $x_k - y_k \to 0$. By Theorem 3 $C = \operatorname{diag}(\sqrt{x_k})$. Let us put $C_1 = \operatorname{diag}(\sqrt{y_k})$ then since $C - C_1 \in \mathcal{K}$ and $\mathcal{K} \subseteq C^*(U, C)$ and $\mathcal{K} \subseteq C^*(U, C_1)$ we conclude that $C^*(U, C) = C^*(U, C_1)$ as operator algebras. If $m = |\omega(\delta)|$ then $C^*(U, C_1) = C^*(UC_1)$ is an algebra generated by *m*-periodic weighted shift. It is known that C^* -algebra generated by all *m*periodic weighted shifts in a given separable Hilbert space is isomorphic to $\mathcal{T}(C(\mathbb{T}))$ (the Toeplitz operators) and there is *m*-periodic shift which generate this algebra. But if $D = \operatorname{diag}(d_k)$ and $D_1 = \operatorname{diag}(d_k^1)$ are diagonal operators with *m*-periodic coefficients and *m* is least possible, then $C^*(U,D) = C^*(U,D_1)$ (since the map $g: d_k \to d_k^1$ is continuous on $\sigma(D)$ and by functional calculus $f(D) = D_1$ and obviously, $f(UDU^*) = Uf(D)U^*$). From these facts follows that $C^*(X) = \mathcal{T}(C(\mathbb{T}))$. For anti-Fock orbits arguments are the same.

Define support of the dynamical system (f, \mathbb{R}) to be the union $X = \bigcup_{\delta \in \text{Orb}_+(f)} \delta$ and the finite support X_{fin} to be union of positive cycles.

Theorem 5. If a dynamical system $(f(), \mathbb{R})$ is simple then the C^{*}-algebra \mathcal{E}_f is GCR (type I C^{*}-algebra), and the finite spectrum is homeomorphic to X_{fin}/\sim , where \sim is an orbit equivalence relation.

Proof. First let us show that the finite-dimensional spectrum $\operatorname{Irr}(\mathcal{E}_f) \simeq (X_{\operatorname{fin}}/) \times S^1(X_{\operatorname{fin}})$ is finite and so compact set). Indeed, it is obvious that $f: X_{\operatorname{fin}} \to X_{\operatorname{fin}}$ is one-to-one map. Thus we can apply the results from [5]. If $\delta \in \operatorname{Orb}_+(f)$ is a bilateral non-periodic orbit then by previous theorem, the set of irreducible representation $\operatorname{Irr}(C^*(\pi_{\delta}))$ is $\pi_{\delta}, \pi_{\omega(\delta),\phi}, \pi_{\alpha(\delta),\phi}$, where $0 \leq \phi \leq 2\pi$. Since we know the topology on finite dimensional representations we conclude that $\operatorname{Irr}(C^*(\pi_{\delta}))$ is T_0 space. Hence $C^*(\pi_{\delta})$ is GCR- C^* -algebra. It is known that $\mathcal{T}(C(\mathbb{T}))$ is also a GCR algebra. Cyclic orbits generate finite-dimensional and so GCR algebras. Hence \mathcal{E}_f is GCR C^* -algebra.

The question of isomorphism of enveloping C^* -algebras may turn to be very difficult. Even in one-to-one case there are only fragmentary results in this direction, for example it is known that for minimal dynamical systems on Cantor sets the isomorphism of cross-product C^* algebras equivalent to orbit-equivalence of corresponding dynamical systems (with condition on K_0 -groups) see [4], Theorem 4. However, in particular "discrete" case we have the following:

Theorem 6. If dynamical systems (f_1, \mathbb{R}) and (f_2, \mathbb{R}) are simple and for every non-cyclic orbit $\delta \in \operatorname{Orb}_+(f_k)$ there exists a point $x \in \delta$ isolated in the support space of a dynamical system (f_k, \mathbb{R}) , then $\mathcal{E}_{f_1} \cong \mathcal{E}_{f_2}$ if and only if there is a one-to-one map $\phi : \operatorname{Orb}_+(f_1)/\sim \to \operatorname{Orb}_+(f_2)/\sim$, such that $|\phi(\delta)| = |\delta|$ and $\phi(\omega(\delta)) = \omega(\phi(\delta))$, $\phi(\alpha(\delta)) = \alpha(\phi(\delta))$. Moreover in this case the topology on $\operatorname{Irr}(\mathcal{E}_{f_k})$ is given by its base consisting of closed sets $\{\pi_{\delta}, \pi_{\omega(\delta),\phi}, \pi_{\alpha(\delta),\phi} | \phi \in S^1\}$, where $\delta \in \operatorname{Orb}_+(f_k) \setminus \operatorname{Cyc}(f)$ and $\{\pi_{\beta,\phi} | \phi \in M\}$, where β is a positive cycle and M is a closed subset in S^1 . Thus $\mathcal{E}_{f_1} \cong \mathcal{E}_{f_2}$ if and only if their dual spaces are homeomorphic.

Proof. Let $\psi : \mathcal{E}(f_1) \to \mathcal{E}(f_2)$ be an isomorphism. Then ψ induces a homeomorphism of spectra spaces $\psi^* : \mathcal{E}(\hat{f}_2) \to \mathcal{E}(\hat{f}_1)$, i.e the spaces of irreducible representation with Jacobson's topology. We know that $\mathcal{E}(\hat{f}_1)$ can be identified with $\operatorname{Orb}_+(f_1)/$. With this identification we have one-to-one map $\psi^* : \operatorname{Orb}_+(f_2) \to \operatorname{Orb}_(f_1)$. Let $\delta \in \operatorname{Orb}_+(f_2)$. Since ψ^* is homeomorphism $\psi^*(\overline{\pi_\delta}) = \overline{\psi^*(\pi)}$. As we know that $\overline{\pi_\delta} = \{\pi_{\omega(\delta),\phi}, \pi_{\alpha(\delta),\phi}\}$ we have proved necessity of conditions of the theorem.

Let ω_1 be a non-attractive cycle or an empty set and ω_2 be non-repellent cycles or empty sets. Denote $\Omega_{\omega_1}^{\omega_2}(f_1) = \{\delta \in \operatorname{Orb}_+(f_1) | \alpha(\delta) = \omega_1, \ \omega(\delta) = \omega_2\}$. Then $\operatorname{Orb}_+(f_1)$ is the disjoint union of these sets. For all $\delta \in \Omega_{\omega_1}^{\omega_2}(f_1)$ we will realize the corresponding representation π_{δ} in the same Hilbert space $H_{\omega_1}^{\omega_2}$. Consider an atomic representation of $\mathcal{E}(f_1)$ which is realized in Hilbert space $H = \otimes_{\omega_1,\omega_2}(H_{\omega_1}^{\omega_2})^{\otimes n(\omega_1,\omega_2)}$, where $n(\omega_1,\omega_2) = |\Omega_{\omega_1}^{\omega_2}(f_1)|$. Then $\mathcal{E}(f_1)$ is isomorphic to the algebra generated by the diagonal operator $C = \operatorname{diag}(C_{\omega_1,\omega_2})$, where $C_{\omega_1,\omega_2} = \operatorname{diag}(C_{\delta}|\delta \in \Omega_{\omega_1,\omega_2})$ and $C_{\delta} = \pi_{\delta}((XX^*)^{1/2})$ and block-diagonal with respect to direct sum decomposition of H operator $U = \operatorname{diag}(I \otimes U_{\omega_1,\omega_2})$. Discreteness of the dynamical system implies that all block-diagonal with respect to an expanded direct sum decomposition $H = \bigoplus_{\omega_1,\omega_2} \bigoplus_{n(\omega_1,\omega_2)} H_{\omega_1}^{\omega_2}$ compact operators belong to $\mathcal{E}(f_1)$. We will denote this subalgebra of compact operators by \mathcal{K}_1 . Modulo this compact operators C is $C_{\operatorname{normal}} = \operatorname{diag}(I \otimes A_{\omega_1,\omega_2})$, where $A_{\omega_1,\omega_2}e_k = x_k^1$ if k < 0 and $A_{\omega_1,\omega_2}e_k = x_k^2$ if k > 0 and $\omega_1 = (x_{-k}^1)_{k\in\mathbb{N}}$ and $\omega_2 = (x_k^2)_{k\in\mathbb{N}}$ regarded as periodic orbit. Moreover, it is obvious that $C^*(C_{\text{normal}}, U, \mathcal{K}_1) = C^*(C, U)$. If there is ϕ which satisfies all conditions of the theorem then we can consider $\mathcal{E}(f_1)$ and $\mathcal{E}(f_2)$ in the same Hilbert space and using functional calculus obtain $\phi^*(C_{\text{normal}}(f_1)) = C_{\text{normal}}(f_2)$ (where ϕ^* is continuous map $\operatorname{Per}(f_1)_+ \to \operatorname{Per}(f_2)_+$ which is lifting of ϕ). and so $\mathcal{E}(f_1) = \mathcal{E}(f_2)$ as operator algebras. This completes the proof.

5 Quadratic dynamical system

In this section we consider an example of one-dimensional quadratic dynamical system. Let $f_{a,b}(x) = 1 + ax - bx^2$ with $\{a, b\} \in \mathbb{R}$ and b > 0 to provide boundedness. Since when a < 0 dynamical system is one-to-one on \mathbb{R}_+ (and so all irreducible representations are one-dimensional) we assume that a > 0. This dynamical system is conjugated to $f_{\mu}(x) = \mu x(1 - x)$, where $\mu = 1 + \sqrt{a^2 - 2a + 1 + 4b}$. The values of parameter μ when bifurcations of cycles of one parametric family $\{f_{\mu}\}$ occur are given in [12]. However a conjugacy relation does not preserve positiveness, i.e. $\operatorname{Orb}_+(f_{a,b})$ may not map into $\operatorname{Orb}_+(f_{\mu})$.

If (a, b) belong to domain $D = \{(a, b) | b < \frac{1}{2} - \frac{a^2}{4} + \frac{a}{2} + \frac{\sqrt{1+2a}}{2}\}$ bounded by curve G (see Fig. 1) then for every $x \in [0; \sup f_{a,b}] \mathcal{O}_+(x) \subset [0; \sup f_{a,b}]$. Thus for such (a, b) algebra $A_{a,b}$ has Fock representation and as it easily can be shown has no anti-Fock representations. In the complement of D algebra $A_{a,b}$ has anti-Fock representations.



Proposition 1. If (a,b) belong to domain $P_1 = \{(a,b) | b < 1 - \frac{(a-1)^2}{4}\}$ bounded by curve Γ_2 (see picture) then $\mathcal{E}_{a,b}$ has one dimensional and Fock irreducible representations only. Moreover $\mathcal{E}_{a,b} \simeq \mathcal{T}(C(\mathbb{T})).$

Proof. For $(a,b) \in P_1$ dynamical system has two fix point $\beta_+ > 0$, $\beta_- < 0$ but has no other cycles.

Let us show that $\operatorname{Orb}_+(f) = \{\beta_+, \delta_1\}$, where δ_1 is the Fock orbit. If $\delta \in \operatorname{Orb}_+(f)$ and $\delta \neq \beta_+, \delta \neq \delta_1$ then $\alpha(\delta)$ is a cycle which cannot be the attractive point β_+ (see Lemma 1). Hence $\alpha(\delta) = \beta_- < 0$ which is contradiction.

Since $P_1 \subset D$ then δ_1 is positive orbit. And Theorem 4 implies that $\mathcal{E}_{a,b} \simeq \mathcal{T}(C(\mathbb{T}))$.

Proposition 2. Let $P_2 = \{(a, b) | 1 - \frac{(a-1)^2}{4} < b < -\frac{a^2}{4} + \frac{a}{2} + \frac{5}{4}\}$ bounded by curves Γ_2 and Γ_4 . Domain P_2 is divided into three domains P_2^1 , P_2^2 , P_2^3 : $P_2^1 = P_2 \cap D$; $P_2^2 = \{(a, b) \in P_2 \setminus D | b < a+1\}$; $P_2^3 = \{(a, b) \in P_2 \setminus D | b > a+1\}$. Then for $(a,b) \in P_2^1 C^*$ -algebra $\mathcal{E}_{a,b}$ has the family of one-, two-dimensional and the Fock irreducible representations, but has no anti-Fock representation. For $(a,b) \in P_2^2 C^*$ -algebra $\mathcal{E}_{a,b}$ has one-, two-dimensional and anti-Fock irreducible representations, but has no Fock representation. For $(a,b) \in P_2^3 C^*$ -algebra $\mathcal{E}_{a,b}$ has the family of one-dimensional representations and anti-Fock irreducible representations, but has no two-dimensional and Fock representations.

For (a, b) from domain bounded by curves Γ_4 and Γ_8 the dynamical system has 4-cycle, 2-cycle, two fix points and no other cycles.

The curve $\Gamma_{2^{\infty}} \approx 1.651225 - \frac{(a-1)^2}{4}$ separates the domain, where the dynamical system is simple from the one (including $\Gamma_{2^{\infty}}$), where the dynamical system is not simple.

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References

- Bunce J.W. and Deddens J.A., C^{*}-algebras generated by Weighted shifts, Indiana University Math. J., 1973, V.23, N 3.
- [2] Dixmier J., C*-Algebars and Their Representations, Moscow, Nauka, 1974.
- [3] Fairlie D.B., Quantum deformations of SU(2), J. Phys. A: Math. and Gen., 1990, V.23, 183–186.
- [4] Herman R.H., Putnam I.F. and Skau C.F., Ordered Bratteli diagrams, dimension groups and topological dynamics, *Intern. J. Math.*, 1992, V.3, 827–864.
- [5] Kawamura S., Tomiyama J. and Watatani Y., Finite-dimensional irreducible representations of C^{*}-algebras associated with topological dynamical systems, *Math. Scand.*, 1985, V.56, 241–248.
- [6] Klimek S. and Lesnevski A., Quantum Riemann surfaces I. The unit disc, Commun. Math. Phys., 1992, V.146, 103–122.
- [7] Klimyk A. and Schmüdgen K., Quantum Groups and Their Representations, Berlin, Heidelberg, New York, Springer, 1997.
- [8] Ostrovskyi V. and Samoilenko Yu., Introduction to the Theory of Representations of Finitely-Presented *-Algebras. I. Representations by Bounded Operators, Amsterdam, The Gordon and Breach Publ., 1999.
- [9] Pedersen G.K., C*-Algebras and Their Automorphism Groups, London Math. Soc. Monographs, Vol.14, London, Academic Press, 1979.
- [10] Popovych S., Representation of real forms of Witten's first deformation, in Proc. of the Second International Conference "Symmetry in Nonlinear Mathematical Physics", 1997, V.2, 393–396.
- [11] Samoilenko Yu.S., Turowska L.B. and Popovych S., Representations of a cubic deformation of su(2) and parasupersymmetric commutation relations, in Proc. of the Second International Conference "Symmetry in Nonlinear Mathematical Physics", 1997, V.2, 372–383.
- [12] Sharkovskii A.N., Maistrenko Yu.L. and Romanenko J.Yu., Deference Equations and Their Aplications, Kyiv, Naukova dumka, 1986 (in Russian).