## Some exact solutions of the many-dimensional sine-Gordon equation

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In the present paper we construct the multiparametrical families of exact solutions of the many-dimensional nonlinear d'Alembert equation

$$\Box U = \sinh U,\tag{1}$$

where  $\Box = \partial^2 / \partial x_0^2 - \cdots - \partial^2 / \partial x_n^2$ . This equation is concerned with some problems of field theory [1]. In the case of n = 1 the analysis and the physical interpretation of solutions of this equations are given in [2].

Up to date the inverse-scattering method is applied for solving two-dimensional nonlinear equations (KdV, sine-Gordon, nonlinear Schrödinger and some others) mainly and the attempts to extend this method for solving many-dimensional equations are not so successful.

To construct some classes of exact solutions of the many-dimensional equation (1), we use group-theoretical ideas of Lie which were applied fruitfully by Birkhoff [3], Sedov [4] and Ovsyannikov [5] to nonlinear equations of hydrodynamics.

The maximal local invariance group of eq.(1) is the Poincaré group P(1,n) of rotations and translations of the (1+n)-dimensional space  $R^{1,n}$ .

We look for the solutions to eq.(1) of the form

$$U(x) = \varphi(\omega), \tag{2}$$

where  $\varphi$  is a function of the invariant variable  $\omega$  only (for more details see [6]). We use the following set of invariants which were presented in [7]. (Below the summation convention is employed. The parameters  $\alpha_{\nu}, \beta_{\nu}, \ldots$  are arbitrary real constants.)

$$\omega = (x_\nu x^\nu)^{1/2},\tag{3a}$$

$$\omega = \left[ (\beta_{\nu} y^{\nu})^2 + y_{\nu} y^{\nu} \right]^{1/2}, \qquad \beta_{\nu} \beta^{\nu} = -1,$$
(3b)

$$\omega = \left[ (\beta_{\nu} y^{\nu})^2 - y_{\nu} y^{\nu} \right]^{1/2}, \qquad \beta_{\nu} \beta^{\nu} = 1,$$
(3c)

$$\omega = \alpha_{\nu} x^{\nu}, \qquad \alpha_{\nu} \alpha^{\nu} = l = \pm 1,$$
(3d)

$$\omega = \frac{1}{2} (\alpha_{\nu} y^{\nu})^2 + a\beta_{\nu} y^{\nu}, \qquad \alpha_{\nu} \alpha^{\nu} = \alpha_{\nu} \beta^{\nu} = 0, \qquad \beta_{\nu} \beta^{\nu} = l = \pm 1,$$
(3e)

$$\omega = \beta_{\nu} y^{\nu} + a \ln \alpha_{\nu} y^{\nu}, \qquad a \neq 0,$$
  

$$\alpha_{\nu} \alpha^{\nu} = \alpha_{\nu} \beta^{\nu} = 0, \qquad \beta_{\nu} \beta^{\nu} = l = \pm 1, \qquad y^{\nu} = x^{\nu} + a^{\nu},$$
(3f)

 $a, a^{\nu}$  are arbitrary constants.

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Substituting (2) into the many-dimensional partial differential eq.(1) we reduce it to the ordinary differential equations

$$U' + \frac{N_1}{\omega}U' = \sinh U \tag{4a}$$

(cases (3a), (3b))

$$-U' - \frac{N_2}{\omega}U' = \sinh U,\tag{4b}$$

(case (3c))

$$U' = l \sinh U, \qquad l = \pm 1 \tag{4c}$$

(cases (3d), (3e)).

Here  $N_1$  and  $N_2$  are natural numbers depending on the value of the space dimension n.

When  $N_1 \neq 0$  and  $N_2 \neq 0$  eqs.(4a) and (4b) cannot be solved in explicit form. Taking l = 1 we have from (4c)

$$w = \int \frac{du}{\sqrt{2\cosh U + C}},$$

C is an arbitrary constant. The solutions of eq.(1) are found by in inversion of elliptic integrals [8]:

$$U = 2 \operatorname{tgh}^{-1} \{ \operatorname{sn}(z,k) \}, \qquad z = \frac{\sqrt{C+2}}{2} \omega, \quad k^2 = \frac{C-2}{C+2}, \quad C > 2, \tag{5a}$$

$$U = 2 \operatorname{tgh}^{-1} \{ \sin z \}, \qquad z = \sqrt{2}\omega, \quad C = 2.$$
 (5b)

Writing the integration constant in the form

$$w = \int \frac{du}{\sqrt{2\cosh U - C}}$$

we have analogously

$$U = \cosh^{-1} \frac{2 - C \operatorname{sn}^2(\omega, k)}{2 \operatorname{cn}^2(\omega, k)}, \qquad k^2 = \frac{C+2}{4}, \quad 0 < C < 2,$$
(5c)

$$U = \cosh^{-1} \frac{C/2 - \operatorname{sn}^2(z, k)}{\operatorname{cn}^2(z, k)}, \qquad z = \frac{\sqrt{C+2}}{2}\omega, \quad k^2 = \frac{4}{C+2}, \quad C > 2, \quad (5d)$$

$$U = 4 \operatorname{tgh}^{-1} \exp[\omega], \qquad C = 2, \tag{5e}$$

$$U = \cosh^{-1} \{ \operatorname{cn}(\omega, k) \}^{-1}, \qquad C = 0, \quad k^2 = \frac{1}{2}.$$
(5f)

When l = -1 we have the solution

$$U = \cosh^{-1}\left\{\frac{C}{2}\operatorname{cn}^{2}(z,k) + \operatorname{sn}^{2}(z,k)\right\},\$$

$$z = \frac{\sqrt{C+2}}{2}\omega, \quad k^{2} = \frac{C-2}{C+2}, \quad C > 2.$$
(5g)

Here  $\operatorname{sn} z$  and  $\operatorname{cn} z$  are Jacobi elliptic functions.

We add that in ref. [6, 7] the same way some solutions of the many-dimensional equation  $\Box U = \sin U$  are found. In [9] the solutions of this equation were obtained by symmetry reduction.

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