# Binary jumps in continuum. II. Non-equilibrium process and a Vlasov-type scaling limit 

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#### Abstract

Let $\Gamma$ denote the space of all locally finite subsets (configurations) in $\mathbb{R}^{d}$. A stochastic dynamics of binary jumps in continuum is a Markov process on $\Gamma$ in which pairs of particles simultaneously hop over $\mathbb{R}^{d}$. We discuss a non-equilibrium dynamics of binary jumps. We prove the existence of an evolution of correlation functions on a finite time interval. We also show that a Vlasov-type mesoscopic scaling for such a dynamics leads to a generalized Boltzmann non-linear equation for the particle density.


Key words: continuous system, binary jumps, non-equilibrium dynamics, correlation functions, scaling limit, Vlasov scaling, Poisson measure
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## 1 Introduction

Let $\Gamma=\Gamma\left(\mathbb{R}^{d}\right)$ denote the space of all locally finite subsets (configurations) in $\mathbb{R}^{d}, d \in \mathbb{N}$. A stochastic dynamics of jumps in continuum is a Markov process on $\Gamma$ in which groups of particles simultaneously hop over $\mathbb{R}^{d}$, i.e., at each jump time several points of the configuration change their positions.

The simplest case corresponds to the so-called Kawasaki-type dynamics in continuum. This dynamics is a Markov process on $\Gamma$ in which particles hop over $\mathbb{R}^{d}$ so that, at each jump time, only one particle changes its position. For a study of an equilibrium Kawasaki dynamics in continuum, we refer the reader

[^0]to the papers $[11-13,17,18,20,21]$ and the references therein. Under the socalled balance condition on the jump rate, a proper Gibbs distribution is an invariant, and even symmetrizing measure for such a dynamics, see [13]. To obtain a simpler measure, e.g. Poissonian, as a symmetrizing measure for a Kawasaki-type dynamics, we should either suppose quite unnatural conditions on the jump rate, or consider a free dynamics. In the free Kawasaki dynamics, in the course of a random evolution, each particle of the configuration randomly hops over $\mathbb{R}^{d}$ without any interaction with other particles (see [19] for details).

In the dynamics of binary jumps, at each jump time two points of the (infinite) configuration change their positions in $\mathbb{R}^{d}$. A randomness for choosing a pair of points provides a random interaction between particles of the system, even in the case where the jump rate only depends on the hopping points (and does not depend on the other points of the configuration). A Poisson measure may be invariant, or even symmetrizing for such a dynamics. In the first part of the present paper [10], we considered such a process with generator

$$
\begin{align*}
(L F)(\gamma)= & \sum_{\left\{x_{1}, x_{2}\right\} \subset \gamma}
\end{align*} \int_{\left(\mathbb{R}^{d}\right)^{2}} Q\left(x_{1}, x_{2}, d h_{1} \times d h_{2}\right) .
$$

Here, the measure $Q\left(x_{1}, x_{2}, d h_{1} \times d h_{2}\right)$ describes the rate at which two particles, $x_{1}$ and $x_{2}$, of configuration $\gamma$ simultaneously hop to $x_{1}+h_{1}$ and $x_{2}+h_{2}$, respectively. Under some additional conditions on the measure $Q$, we studied a corresponding equilibrium dynamics for which a Poisson measure is a symmetrizing measure. We also considered two different scalings of the rate measure $Q$, which led us to a diffusive dynamics and to a birth-and-death dynamics, respectively.

In the present paper, we restrict our attention to the special case of the generator (1.1). We denote the arrival points by $y_{i}=x_{i}+h_{i}, i=1,2$. An (informal) generator of a binary jump process of our interest is given by

$$
\begin{align*}
(L F)(\gamma)=\sum_{\left\{x_{1}, x_{2}\right\} \subset \gamma} \int_{\mathbb{R}^{d}} & \int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) \\
& \times\left(F\left(\gamma \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{y_{1}, y_{2}\right\}\right)-F(\gamma)\right) d y_{1} d y_{2} . \tag{1.2}
\end{align*}
$$

Here

$$
\begin{equation*}
c\left(x_{1}, x_{2}, y_{1}, y_{2}\right)=c\left(\left\{x_{1}, x_{2}\right\},\left\{y_{1}, y_{2}\right\}\right) \tag{1.3}
\end{equation*}
$$

is a non-negative measurable function. Our aim is to study a non-equilibrium dynamics corresponding to (1.2). Note that the Poisson measure with any positive constant intensity will be invariant for this dynamics. If, additionally,

$$
\begin{equation*}
c\left(\left\{x_{1}, x_{2}\right\},\left\{y_{1}, y_{2}\right\}\right)=c\left(\left\{y_{1}, y_{2}\right\},\left\{x_{1}, x_{2}\right\}\right), \tag{1.4}
\end{equation*}
$$

then any such measure will even be symmetrizing.

It should be noted that similar dynamics of finite particle systems were studied in $[2,3,14]$. In particular, in [3], the authors studied a non-equilibrium dynamics of velocities of particles, such that the law of conservation of momentum is satisfied for this system. However, the methods applied to finite particle systems seem not to be applicable to infinite systems of our interest.

Let us also note an essential difference between lattice and continuous systems. An important example of a Markov dynamics on lattice configurations is the so-called exclusion process. In this process, particles randomly hop over the lattice under the only restriction to have no more than one particle at each site of the lattice. This process may have a Bernoulli measure as an invariant (and even symmetrizing) measure, but a corresponding stochastic dynamics has non-trivial properties and possesses an interesting and reach scaling limit behavior. A straightforward generalization of the exclusion process to the continuum gives a free Kawasaki dynamics, because the exclusion restriction (yielding an interaction between particles) obviously disappears for configurations in continuum. To introduce the simplest (in certain sense) interaction, we consider the generator above.

The generator (1.2) informally provides a functional evolution via the (backward) Kolmogorov equation

$$
\begin{equation*}
\frac{\partial F_{t}}{\partial t}=L F_{t},\left.\quad F_{t}\right|_{t=0}=F_{0} \tag{1.5}
\end{equation*}
$$

However, the problem of existence of a solution to (1.5) in some functional space seems to be a very difficult problem. Fortunately, in applications, we usually need only an information about a mean value of a function on $\Gamma$ with respect to some probability measure on $\Gamma$, rather than a full point-wise information about this function. Therefore, we turn to a weak evolution of probability measures (states) on $\Gamma$. This evolution of states is informally given as a solution to the initial value problem:

$$
\begin{equation*}
\frac{d}{d t} \int_{\Gamma} F d \mu_{t}=\int_{\Gamma} L F d \mu_{t},\left.\quad \mu_{t}\right|_{t=0}=\mu_{0} \tag{1.6}
\end{equation*}
$$

provided, of course, that a solution exists. The problem (1.6) may be rewritten in terms of correlation functionals $k_{t}$ of states $\mu_{t}$ :

$$
\begin{equation*}
\frac{d}{d t} \int_{\Gamma_{0}} G \cdot k_{t} d \lambda=\int_{\Gamma_{0}} \widehat{L} G \cdot k_{t} d \lambda,\left.\quad k_{t}\right|_{t=0}=k_{0} \tag{1.7}
\end{equation*}
$$

Here $\Gamma_{0}$ is the space of all finite configurations in $\mathbb{R}^{d}$ and $\lambda$ is the LebesguePoisson measure on $\Gamma_{0}$. Functions $G$ on $\Gamma_{0}$ are called quasi-observables. Note that every correlation functional $k_{t}$ can be considered as an infinite vector $\left(k_{t}^{(n)}\right)_{n=0}^{\infty}$, where $k_{t}^{(n)}$ is a function on $n$-point configurations in $\mathbb{R}^{d}$, called the $n$-th correlation function.

The dynamics of correlation functions corresponding to (1.7) has a chain structure which is similar to the BBGKY-hierarchy for Hamiltonian dynamics.

The corresponding generator of this dynamics has an upper triangular, twodiagonal structure in a Fock-type space. Note that the most important case from the point of view of applications, being also the most interesting from the mathematical point of view, is the case of bounded (non-integrable) correlation functions. Because of (1.7), the dynamics of correlation functions should be treated in a weak form. Therefore, we consider a pre-dual evolution of quasiobservables, whose generator has a lower triangular, two-diagonal structure in a Fock-type space of integrable functions.

The paper is organized as follows. In Section 2, we describe the model and give some necessary preliminary information. Section 3 is devoted to the functional evolution of both quasi-observables and correlation functions. We derive some information about the structure and properties of the lower triangular, two-diagonal generator of the dynamics of quasi-observables (Propositions 3.1 and 3.3). We further prove a general result for a lower triangular, two-diagonal generator, which shows that the evolution may be obtained, for all times, recursively in a scale of Banach spaces whose norms depend on time (Theorem 3.5). Next, we construct a dual dynamics on a finite time interval (Theorem 3.6). We show that this dynamics is indeed an evolution of correlation functions, which in turn leads to an evolution of probability measures on $\Gamma$ (Theorem 3.9). In Section 4 we consider a Vlasov-type scaling for our dynamics. The limiting evolution (which exists by Proposition 4.4) has a chaos preservation property. This means that the corresponding dynamics of states transfers a Poisson measure with a non-homogeneous intensity into a Poisson measure whose nonhomogeneous intensity satisfies a non-linear evolution (kinetic) equation. We finally present sufficient conditions for the existence and uniqueness of a solution to this equation (Proposition 4.6).

It is worth noting that we rigorously prove the convergence of the scaled evolution of the infinite particle system to the limiting evolution (which in turn leads to the kinetic equation). This seems to be a new step even for finite particle systems.

## 2 Preliminaries

Let $\mathcal{B}\left(\mathbb{R}^{d}\right)$ be the Borel $\sigma$-algebra on $\mathbb{R}^{d}, d \in \mathbb{N}$, and let $\mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$ denote the system of all bounded sets from $\mathcal{B}\left(\mathbb{R}^{d}\right)$. The configuration space over $\mathbb{R}^{d}$ is defined as the set of all locally finite subsets of $\mathbb{R}^{d}$ :

$$
\Gamma:=\left\{\gamma \subset \mathbb{R}^{d}| | \gamma_{\Lambda} \mid<\infty \text { for any } \Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)\right\} .
$$

Here $|\cdot|$ denotes the cardinality of a set and $\gamma_{\Lambda}:=\gamma \cap \Lambda$. One can identify any $\gamma \in \Gamma$ with the positive Radon measure $\sum_{x \in \gamma} \delta_{x} \in \mathcal{M}\left(\mathbb{R}^{d}\right)$, where $\delta_{x}$ is the Dirac measure with mass at $x$, and $\mathcal{M}\left(\mathbb{R}^{d}\right)$ stands for the set of all positive Radon measures on $\mathcal{B}\left(\mathbb{R}^{d}\right)$. The space $\Gamma$ can be endowed with the relative topology as a subset of the space $\mathcal{M}\left(\mathbb{R}^{d}\right)$ with the vague topology, i.e., the weakest topology on $\Gamma$ with respect to which all maps $\Gamma \ni \gamma \mapsto\langle f, \gamma\rangle:=$ $\int_{\mathbb{R}^{d}} f(x) \gamma(d x)=\sum_{x \in \gamma} f(x), f \in C_{0}\left(\mathbb{R}^{d}\right)$, are continuous. Here, $C_{0}\left(\mathbb{R}^{d}\right)$ is the
space of all continuous functions on $\mathbb{R}^{d}$ with compact support. The corresponding Borel $\sigma$-algebra $\mathcal{B}(\Gamma)$ coincides with the smallest $\sigma$-algebra on $\Gamma$ for which all mappings $\Gamma \ni \gamma \mapsto\left|\gamma_{\Lambda}\right| \in \mathbb{N}_{0}:=\mathbb{N} \cup\{0\}$ are measurable for any $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$, see e.g. [1]. It is worth noting that $\Gamma$ is a Polish space (see e.g. [16]).

Let $\mathcal{F}_{\text {cyl }}(\Gamma)$ denote the set of all measurable cylinder functions on $\Gamma$. Each $F \in \mathcal{F}_{\text {cyl }}(\Gamma)$ is characterized by the following property: $F(\gamma)=F\left(\gamma_{\Lambda}\right)$ for some $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$ and for any $\gamma \in \Gamma$.

A stochastic dynamics of binary jumps in continuum is a Markov process on $\Gamma$ in which pairs of particles simultaneously hop over $\mathbb{R}^{d}$, i.e., at each jump time two points of the configuration change their positions. Thus, an (informal) generator of such a process has the form (1.2), where $c\left(x_{1}, x_{2}, y_{1}, y_{2}\right)$ is a nonnegative measurable function which satisfies (1.3) and

$$
c\left(x_{1}, x_{2}, \cdot, \cdot\right) \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{d} \times \mathbb{R}^{d}\right) \quad \text { for a.a. } x_{1}, x_{2} \in \mathbb{R}^{d}
$$

The function $c$ describes the rate at which pairs of particles hop over $\mathbb{R}^{d}$.
Remark 2.1. Note that, in general, the expression on the right hand side of (1.2) is not necessarily well defined for all $\gamma \in \Gamma$, even if $F \in \mathcal{F}_{\text {cyl }}(\Gamma)$. Nevertheless, for such $F,(L F)(\gamma)$ has sense at least for all $\gamma \in \Gamma$ with $|\gamma|<\infty$.

In various applications, the evolution of states of the system (i.e., measures on the configuration space $\Gamma$ ) helps one to understand the behavior of the process and gives possible candidates for invariant states. In fact, various properties of such an evolution form the main information needed in applications. Using the duality between functions and measures, this evolution may be considered in a weak form, given, as usual, by the expression

$$
\begin{equation*}
\langle F, \mu\rangle=\int_{\Gamma} F(\gamma) d \mu(\gamma) \tag{2.1}
\end{equation*}
$$

Therefore, the evolution of states is informally given as a solution to the initial value problem:

$$
\begin{equation*}
\frac{d}{d t}\left\langle F, \mu_{t}\right\rangle=\left\langle L F, \mu_{t}\right\rangle,\left.\quad \mu_{t}\right|_{t=0}=\mu_{0} \tag{2.2}
\end{equation*}
$$

provided, of course, that a solution exists. For a wide class of probability measures on $\Gamma$, one can consider a corresponding evolution of their correlation functionals, see below.

The space of $n$-point configurations in an arbitrary $Y \in \mathcal{B}\left(\mathbb{R}^{d}\right)$ is defined by

$$
\Gamma^{(n)}(Y):=\{\eta \subset Y| | \eta \mid=n\}, \quad n \in \mathbb{N}
$$

We set $\Gamma^{(0)}(Y):=\{\emptyset\}$. As a set, $\Gamma^{(n)}(Y)$ may be identified with the quotient of $\widetilde{Y^{n}}:=\left\{\left(x_{1}, \ldots, x_{n}\right) \in Y^{n} \mid x_{k} \neq x_{l}\right.$ if $\left.k \neq l\right\}$ with respect to the natural action of the permutation group $S_{n}$ on $\widetilde{Y^{n}}$. Hence, one can introduce the corresponding Borel $\sigma$-algebra, which will be denoted by $\mathcal{B}\left(\Gamma^{(n)}(Y)\right)$. The space of finite configurations in an arbitrary $Y \in \mathcal{B}\left(\mathbb{R}^{d}\right)$ is defined by

$$
\Gamma_{0}(Y):=\bigsqcup_{n \in \mathbb{N}_{0}} \Gamma^{(n)}(Y)
$$

This space is equipped with the disjoint union topology. Therefore, we consider the corresponding Borel $\sigma$-algebra $\mathcal{B}\left(\Gamma_{0}(Y)\right)$. In the case where $Y=\mathbb{R}^{d}$, we will omit $Y$ in the notation, namely, $\Gamma_{0}=\Gamma_{0}\left(\mathbb{R}^{d}\right), \Gamma^{(n)}=\Gamma^{(n)}\left(\mathbb{R}^{d}\right)$.

The image of the Lebesgue product measure $(d x)^{n}$ under the mapping

$$
\widetilde{\left(\mathbb{R}^{d}\right)^{n}} \ni\left(x_{1}, \ldots, x_{n}\right) \mapsto\left\{x_{1}, \ldots, x_{n}\right\} \in \Gamma^{(n)}
$$

will be denoted by $m^{(n)}$. We set $m^{(0)}:=\delta_{\{\emptyset\}}$. Let $z>0$ be fixed. The Lebesgue-Poisson measure $\lambda_{z}$ on $\Gamma_{0}$ is defined by

$$
\begin{equation*}
\lambda_{z}:=\sum_{n=0}^{\infty} \frac{z^{n}}{n!} m^{(n)} \tag{2.3}
\end{equation*}
$$

For any $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$, the restriction of $\lambda_{z}$ to $\Gamma(\Lambda):=\Gamma_{0}(\Lambda)$ will also be denoted by $\lambda_{z}$. The space $(\Gamma, \mathcal{B}(\Gamma))$ is the projective limit of the family of spaces $\{(\Gamma(\Lambda), \mathcal{B}(\Gamma(\Lambda)))\}_{\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)}$. The Poisson measure $\pi_{z}$ on $(\Gamma, \mathcal{B}(\Gamma))$ is given as the projective limit of the family of measures $\left\{\pi_{z}^{\Lambda}\right\}_{\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)}$, where $\pi_{z}^{\Lambda}:=e^{-z m(\Lambda)} \lambda_{z}$ is a probability measure on $(\Gamma(\Lambda), \mathcal{B}(\Gamma(\Lambda)))$ and $m(\Lambda)$ is the Lebesgue measure of $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$; for details see e.g. [1]. We will mostly use the measure $\lambda:=\lambda_{1}$.

A set $M \in \mathcal{B}\left(\Gamma_{0}\right)$ is called bounded if there exist $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$ and $N \in \mathbb{N}$ such that $M \subset \bigsqcup_{n=0}^{N} \Gamma^{(n)}(\Lambda)$. The set of bounded measurable functions with bounded support will be denoted by $B_{\mathrm{bs}}\left(\Gamma_{0}\right)$, i.e., $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ if $G \Gamma_{\Gamma_{0} \backslash M}=0$ for some bounded $M \in \mathcal{B}\left(\Gamma_{0}\right)$. We also consider the larger set $L_{\mathrm{ls}}^{0}\left(\Gamma_{0}\right)$ of all measurable functions on $\Gamma_{0}$ with local support, which means: $G \in L_{\mathrm{ls}}^{0}\left(\Gamma_{0}\right)$ if $G \upharpoonright_{\Gamma_{0} \backslash \Gamma(\Lambda)}=0$ for some $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$. Any $\mathcal{B}\left(\Gamma_{0}\right)$-measurable function $G$ on $\Gamma_{0}$ is, in fact, defined by a sequence of functions $\left\{G^{(n)}\right\}_{n \in \mathbb{N}_{0}}$, where $G^{(n)}$ is a $\mathcal{B}\left(\Gamma^{(n)}\right)$-measurable function on $\Gamma^{(n)}$. Functions on $\Gamma$ and $\Gamma_{0}$ will be called observables and quasi-observables, respectively.

We consider the following mapping from $L_{\mathrm{ls}}^{0}\left(\Gamma_{0}\right)$ into $\mathcal{F}_{\mathrm{cyl}}(\Gamma)$ :

$$
\begin{equation*}
(K G)(\gamma):=\sum_{\eta \Subset \gamma} G(\eta), \quad \gamma \in \Gamma \tag{2.4}
\end{equation*}
$$

where $G \in L_{\mathrm{ls}}^{0}\left(\Gamma_{0}\right)$, see, e.g., $[15,23,24]$. The summation in (2.4) is taken over all finite subconfigurations $\eta \in \Gamma_{0}$ of the (infinite) configuration $\gamma \in \Gamma$; we denote this by the symbol $\eta \Subset \gamma$. The mapping $K$ is linear, positivity preserving, and invertible with

$$
\begin{equation*}
\left(K^{-1} F\right)(\eta)=\sum_{\xi \subset \eta}(-1)^{|\eta \backslash \xi|} F(\xi), \quad \eta \in \Gamma_{0} \tag{2.5}
\end{equation*}
$$

Remark 2.2. Note that, using formula (2.5), we can extend the mapping $K^{-1}$ to functions $F$ which are well-defined, at least, on $\Gamma_{0}$.

Let $\mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ denote the set of all probability measures $\mu$ on $(\Gamma, \mathcal{B}(\Gamma))$ which have finite local moments of all orders, i.e., $\int_{\Gamma}\left|\gamma_{\Lambda}\right|^{n} \mu(d \gamma)<+\infty$ for all $\Lambda \in$ $\mathcal{B}_{b}\left(\mathbb{R}^{d}\right)$ and $n \in \mathbb{N}_{0}$. A measure $\rho$ on $\left(\Gamma_{0}, \mathcal{B}\left(\Gamma_{0}\right)\right)$ is called locally finite if
$\rho(A)<\infty$ for all bounded sets $A$ from $\mathcal{B}\left(\Gamma_{0}\right)$. The set of all such measures is denoted by $\mathcal{M}_{\mathrm{lf}}\left(\Gamma_{0}\right)$.

One can define a transform $K^{*}: \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma) \rightarrow \mathcal{M}_{\mathrm{lf}}\left(\Gamma_{0}\right)$, which is dual to the $K$-transform, i.e., for every $\mu \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ and $G \in \mathcal{B}_{\mathrm{bs}}\left(\Gamma_{0}\right)$, we have

$$
\int_{\Gamma}(K G)(\gamma) \mu(d \gamma)=\int_{\Gamma_{0}} G(\eta)\left(K^{*} \mu\right)(d \eta)
$$

The measure $\rho_{\mu}:=K^{*} \mu$ is called the correlation measure of $\mu$.
As shown in [15], for any $\mu \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ and $G \in L^{1}\left(\Gamma_{0}, \rho_{\mu}\right)$, the series (2.4) is $\mu$-a.s. absolutely convergent. Furthermore, $K G \in L^{1}(\Gamma, \mu)$ and

$$
\begin{equation*}
\int_{\Gamma_{0}} G(\eta) \rho_{\mu}(d \eta)=\int_{\Gamma}(K G)(\gamma) \mu(d \gamma) \tag{2.6}
\end{equation*}
$$

A measure $\mu \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ is called locally absolutely continuous with respect to $\pi:=\pi_{1}$ if $\mu^{\Lambda}:=\mu \circ p_{\Lambda}^{-1}$ is absolutely continuous with respect to $\pi^{\Lambda}:=\pi_{1}^{\Lambda}$ for all $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$. Here $\Gamma \ni \gamma \mapsto p_{\Lambda}(\gamma):=\gamma \cap \Lambda \in \Gamma(\Lambda)$. In this case, the correlation measure $\rho_{\mu}$ is absolutely continuous with respect to $\lambda=\lambda_{1}$. A correlation functional of $\mu$ is then defined by

$$
k_{\mu}(\eta):=\frac{d \rho_{\mu}}{d \lambda}(\eta), \quad \eta \in \Gamma_{0}
$$

The functions $k_{\mu}^{(0)}:=1$ and

$$
\begin{gather*}
k_{\mu}^{(n)}:\left(\mathbb{R}^{d}\right)^{n} \longrightarrow \mathbb{R}_{+}, \quad n \in \mathbb{N},  \tag{2.7}\\
k_{\mu}^{(n)}\left(x_{1}, \ldots, x_{n}\right):= \begin{cases}k_{\mu}\left(\left\{x_{1}, \ldots, x_{n}\right\}\right), & \text { if }\left(x_{1}, \ldots, x_{n}\right) \in \widetilde{\left(\mathbb{R}^{d}\right)^{n}} \\
0, & \text { otherwise }\end{cases}
\end{gather*}
$$

are called correlation functions of $\mu$, and they are well known in statistical physics, see e.g [26], [27].

In view of Remark 2.1 and (2.5), the mapping

$$
(\widehat{L} G)(\eta):=\left(K^{-1} L K G\right)(\eta), \quad \eta \in \Gamma_{0}
$$

is well defined for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$, where $K^{-1}$ is understood in the sense of Remark 2.2.

Let $k$ be a measurable function on $\Gamma_{0}$ such that $\int_{B} k d \lambda<\infty$ for any bounded $B \in \mathcal{B}\left(\Gamma_{0}\right)$. Then, for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$, we can consider an analog of pairing (2.1),

$$
\begin{equation*}
\langle\langle G, k\rangle\rangle:=\int_{\Gamma_{0}} G \cdot k d \lambda \tag{2.8}
\end{equation*}
$$

Using this duality, we may consider a dual mapping $\widehat{L}^{*}$ of $\widehat{L}$. As a result, we obtain two initial value problems:

$$
\begin{align*}
\frac{\partial G_{t}}{\partial t} & =\widehat{L} G_{t}, & & \left.G_{t}\right|_{t=0}=G_{0}  \tag{2.9}\\
\frac{d}{d t}\left\langle\left\langle G, k_{t}\right\rangle\right\rangle=\left\langle\left\langle G, \widehat{L}^{*} k_{t}\right\rangle\right\rangle & =\left\langle\left\langle\widehat{L} G, k_{t}\right\rangle\right\rangle, & & \left.k_{t}\right|_{t=0}=k_{0} \tag{2.10}
\end{align*}
$$

We are going to solve the first problem in some functional space over $\Gamma_{0}$. The second problem, corresponding to (2.2), can be realized by means of (2.8), or solved independently.

The next section is devoted to a (rigorous) solution of these two problems.

## 3 Functional evolution

We denote, for any $n \in \mathbb{N}$,

$$
\begin{equation*}
X_{n}:=L^{1}\left(\left(\mathbb{R}^{d}\right)^{n}, d x^{(n)}\right), \tag{3.1}
\end{equation*}
$$

where we set $d x^{(n)}=d x_{1} \cdots d x_{n}$ and $X_{0}:=\mathbb{R}$. The symbol $\|\cdot\|_{X_{n}}$ stands for the norm of the space (3.1).

For an arbitrary $C>0$, we consider the functional Banach space

$$
\begin{equation*}
\mathcal{L}_{C}:=L^{1}\left(\Gamma_{0}, C^{|\eta|} d \lambda(\eta)\right) \tag{3.2}
\end{equation*}
$$

Throughout this paper, the symbol $\|\cdot\|_{\mathcal{L}_{C}}$ denotes for the norm of the space (3.2). Then, for any $G \in \mathcal{L}_{C}$, we may identify $G$ with the sequence $\left(G^{(n)}\right)_{n \geq 0}$, where $G^{(n)}$ is a symmetric function on $\left(\mathbb{R}^{d}\right)^{n}$ (defined almost everywhere) such that

$$
\begin{equation*}
\|G\|_{\mathcal{L}_{C}}=\sum_{n=0}^{\infty} \frac{1}{n!} \int_{\left(\mathbb{R}^{d}\right)^{n}}\left|G^{(n)}\left(x^{(n)}\right)\right| C^{n} d x^{(n)}=\sum_{n=0}^{\infty} \frac{C^{n}}{n!}\left\|G^{(n)}\right\|_{X_{n}}<\infty \tag{3.3}
\end{equation*}
$$

In particular, $G^{(n)} \in X_{n}, n \in \mathbb{N}_{0}$. Here and below we set $x^{(n)}=\left(x_{1}, \ldots, x_{n}\right)$.
We consider the dual space $\left(\mathcal{L}_{C}\right)^{\prime}$ of $\mathcal{L}_{C}$. It is evident that this space can be realized as the Banach space

$$
\mathcal{K}_{C}:=\left\{k: \Gamma_{0} \rightarrow \mathbb{R} \mid k \cdot C^{-|\cdot|} \in L^{\infty}\left(\Gamma_{0}, d \lambda\right)\right\}
$$

with the norm $\|k\|_{\mathcal{K}_{C}}:=\left\|k(\cdot) C^{-|\cdot|}\right\|_{L^{\infty}\left(\Gamma_{0}, \lambda\right)}$, where the pairing between any $G \in \mathcal{L}_{C}$ and $k \in \mathcal{K}_{C}$ is given by (2.8). In particular,

$$
|\langle\langle G, k\rangle\rangle| \leq\|G\|_{C} \cdot\|k\|_{\mathcal{K}_{C}}
$$

Clearly, $k \in \mathcal{K}_{C}$ implies

$$
\begin{equation*}
|k(\eta)| \leq\|k\|_{\mathcal{K}_{C}} C^{|\eta|} \quad \text { for } \lambda \text {-a.a. } \eta \in \Gamma_{0} \tag{3.4}
\end{equation*}
$$

Proposition 3.1. Let for a.a. $x_{1}, x_{2}, y_{1} \in \mathbb{R}^{d}$

$$
\begin{equation*}
\tilde{c}\left(x_{1}, x_{2}, y_{1}\right)=\tilde{c}\left(\left\{x_{1}, x_{2}\right\}, y_{1}\right):=\int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) d y_{2}<\infty . \tag{3.5}
\end{equation*}
$$

Then, for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$, the following formula holds

$$
\begin{equation*}
(\widehat{L} G)(\eta)=\left(L_{0} G\right)(\eta)+(W G)(\eta) \tag{3.6}
\end{equation*}
$$

where

$$
\begin{align*}
\left(L_{0} G\right)(\eta)= & \sum_{\left\{x_{1}, x_{2}\right\} \subset \eta} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) \\
& \times\left(G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{y_{1}, y_{2}\right\}\right)-G(\eta)\right) d y_{1} d y_{2}  \tag{3.7}\\
(W G)(\eta)= & \sum_{x_{2} \in \eta} \sum_{x_{1} \in \eta \backslash x_{2}} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) \\
& \times\left(G\left(\left(\eta \backslash x_{2}\right) \backslash x_{1} \cup y_{1}\right)-G\left(\eta \backslash x_{2}\right)\right) d y_{1} \tag{3.8}
\end{align*}
$$

Proof. We have, for $F=K G$,

$$
\begin{aligned}
& (K G)\left(\gamma \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{y_{1}, y_{2}\right\}\right)-(K G)(\gamma) \\
& =\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{y_{1}, y_{2}\right\}} G(\eta)-\sum_{\eta \Subset \gamma} G(\eta) \\
& =\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup y_{1}\right)+\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup y_{2}\right) \\
& \quad+\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup y_{1} \cup y_{2}\right)-\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup x_{1}\right) \\
& \quad-\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup x_{2}\right)-\sum_{\eta \Subset \gamma \backslash\left\{x_{1}, x_{2}\right\}} G\left(\eta \cup x_{1} \cup x_{2}\right) \\
& = \\
& \quad\left(K \left(G\left(\cdot \cup y_{1}\right)+G\left(\cdot \cup y_{2}\right)+G\left(\cdot \cup\left\{y_{1}, y_{2}\right\}\right)\right.\right. \\
& \left.\left.\quad-G\left(\cdot \cup x_{1}\right)-G\left(\cdot \cup x_{2}\right)-G\left(\cdot \cup\left\{x_{1}, x_{2}\right\}\right)\right)\right)\left(\gamma \backslash\left\{x_{1}, x_{2}\right\}\right) .
\end{aligned}
$$

Hence, for any measurable $h$ on $\Gamma \times \mathbb{R}^{d} \times \mathbb{R}^{d}$,

$$
\begin{aligned}
& K^{-1}\left(\sum_{\left\{x_{1}, x_{2}\right\} \subset} h\left(\cdot \backslash\left\{x_{1}, x_{2}\right\}, x_{1}, x_{2}\right)\right)(\eta) \\
& =\sum_{\xi \subset \eta}(-1)^{|\eta \backslash \xi|} \sum_{\left\{x_{1}, x_{2}\right\} \subset \xi} h\left(\xi \backslash\left\{x_{1}, x_{2}\right\}, x_{1}, x_{2}\right) \\
& =\sum_{\left\{x_{1}, x_{2}\right\} \subset \eta} \sum_{\xi \subset \eta \backslash\left\{x_{1}, x_{2}\right\}}(-1)^{\left|\eta \backslash\left\{x_{1}, x_{2}\right\} \backslash \xi\right|} h\left(\xi, x_{1}, x_{2}\right) \\
& =\sum_{\left\{x_{1}, x_{2}\right\} \subset \eta}\left(K^{-1} h\left(\cdot, x_{1}, x_{2}\right)\right)\left(\eta \backslash\left\{x_{1}, x_{2}\right\}\right) .
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
& (\widehat{L} G)(\eta)=\left(K^{-1} L K G\right)(\eta) \\
& =\sum_{\left\{x_{1}, x_{2}\right\} \subset \eta} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(\left\{x_{1}, x_{2}\right\},\left\{y_{1}, y_{2}\right\}\right)\left(G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup y_{1}\right)\right. \\
& \left.\quad+G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup y_{2}\right)+G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{y_{1}, y_{2}\right\}\right)\right) d y_{1} d y_{2}
\end{aligned}
$$

$$
\begin{aligned}
& -\sum_{\left\{x_{1}, x_{2}\right\} \subset \eta} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(\left\{x_{1}, x_{2}\right\},\left\{y_{1}, y_{2}\right\}\right)\left(G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup x_{1}\right)\right. \\
& \left.\quad+G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup x_{2}\right)+G\left(\eta \backslash\left\{x_{1}, x_{2}\right\} \cup\left\{x_{1}, x_{2}\right\}\right)\right) d y_{1} d y_{2},
\end{aligned}
$$

from where the statement follows.
As we noted before, any function $G$ on $\Gamma_{0}$ may be identified with the infinite vector $\left(G^{(n)}\right)_{n \geq 0}$ of symmetric functions. Due to this identification, any operator on functions on $\Gamma_{0}$ may be considered as an infinite operator matrix. By Proposition 3.1, the operator matrix $\widehat{L}$ has a two-diagonal structure. More precisely, on the main diagonal, we have operators $L_{0}^{(n)}, n \in \mathbb{N}_{0}$, where $L_{0}^{(0)}=L_{0}^{(1)}=0$ and for $n \geq 2$

$$
\begin{align*}
\left(L_{0}^{(n)} G^{(n)}\right)\left(x^{(n)}\right) & =\sum_{i=1}^{n} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{i}, x_{j}, y_{1}, y_{2},\right) \\
& \times\left(G ^ { ( n ) } \left(x_{1}, \ldots, \underset{\hat{i}}{\left.\left.y_{1}, \ldots, y_{2}, \ldots, x_{n}\right)-G^{(n)}\left(x^{(n)}\right)\right) d y_{1} d y_{2}}\right.\right. \tag{3.9}
\end{align*}
$$

On the lower diagonal, we have operators $W^{(n)}, n \in \mathbb{N}$, where $W^{(1)}=0$ and for $n \geq 2$

$$
\begin{align*}
\left(W^{(n)} G^{(n-1)}\right)\left(x^{(n)}\right)= & 2 \sum_{i=1}^{n}
\end{align*} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{i}, x_{j}, y_{1}\right) .
$$

Let us formulate our main conditions on the rate $c$ :

$$
\begin{align*}
& c_{1}:=\underset{x_{1}, x_{2} \in \mathbb{R}^{d}}{\operatorname{ess} \sup } \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) d y_{1} d y_{2}<\infty,  \tag{3.11}\\
& c_{2}:=\underset{x_{1}, x_{2} \in \mathbb{R}^{d}}{\operatorname{ess} \sup _{\mathbb{R}^{d}}} \int_{\mathbb{R}^{d}} c\left(y_{1}, y_{2}, x_{1}, x_{2}\right) d y_{1} d y_{2}<\infty,  \tag{3.12}\\
& c_{3}:=\underset{x_{1} \in \mathbb{R}^{d}}{\operatorname{essssup}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) d y_{1} d y_{2} d x_{2}<\infty,  \tag{3.13}\\
& c_{4}:=\underset{x_{1} \in \mathbb{R}^{d}}{\operatorname{esssup}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(y_{1}, y_{2}, x_{1}, x_{2}\right) d y_{1} d y_{2} d x_{2}<\infty . \tag{3.14}
\end{align*}
$$

Under conditions (3.11)-(3.12), we define the following functions

$$
\begin{align*}
& a_{1}\left(x_{1}, x_{2}\right):=\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{1}, x_{2}, y_{1}, y_{2}\right) d y_{1} d y_{2} \in[0, \infty),  \tag{3.15}\\
& a_{2}\left(x_{1}, x_{2}\right):=\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(y_{1}, y_{2}, x_{1}, x_{2}\right) d y_{1} d y_{2} \in[0, \infty) . \tag{3.16}
\end{align*}
$$

Remark 3.2. Note that, if the function $c$ satisfies the symmetry condition (1.4), then conditions (3.12), (3.14) follow from (3.11), (3.13), respectively, and

$$
\begin{equation*}
a_{1}\left(x_{1}, x_{2}\right)=a_{2}\left(x_{1}, x_{2}\right) \tag{3.17}
\end{equation*}
$$

In this case, the operator $L_{0}^{(n)}$ is symmetric in $L^{2}\left(\left(\mathbb{R}^{d}\right)^{n}, d x^{(n)}\right)$. Moreover, the operator $L$ given by (1.2) is (informally) symmetric in $L^{2}\left(\Gamma, \pi_{z}\right)$ for any $z>0$.
Proposition 3.3. (i) Let (3.11), (3.12) hold. Then, for any $G^{(n)} \in X_{n}$,

$$
\begin{equation*}
\left\|L_{0}^{(n)} G^{(n)}\right\|_{X_{n}} \leq \frac{n(n-1)}{2}\left(c_{1}+c_{2}\right)\left\|G^{(n)}\right\|_{X_{n}} \tag{3.18}
\end{equation*}
$$

Moreover, if additionally

$$
\begin{equation*}
a_{2}\left(x_{1}, x_{2}\right) \leq a_{1}\left(x_{1}, x_{2}\right) \quad \text { for a.a. } x_{1}, x_{2} \in \mathbb{R}^{d}, \tag{3.19}
\end{equation*}
$$

then $L_{0}^{(n)}$ is the generator of a contraction semigroup in $X_{n}$.
(ii) Let (3.13), (3.14) hold. Then, for any $G^{(n-1)} \in X_{n-1}$,

$$
\begin{equation*}
\left\|W^{(n)} G^{(n-1)}\right\|_{X_{n}} \leq n(n-1)\left(c_{3}+c_{4}\right)\left\|G^{(n-1)}\right\|_{X_{n-1}} \tag{3.20}
\end{equation*}
$$

Proof. The estimates (3.18), (3.20) follow directly from (3.9), (3.10) and (3.11)(3.14). To prove that the bounded operator $L_{0}^{(n)}$ is the generator of a contraction semigroup in $X_{n}$, it is enough to show that $L_{0}^{(n)}$ is dissipative (see e.g. [25]). For any $\kappa>0$ and $G^{(n)} \in X_{n}$,

$$
\begin{aligned}
&\left\|L_{0}^{(n)} G^{(n)}-\kappa G^{(n)}\right\|_{X_{n}} \\
&=\left.\int_{\left(\mathbb{R}^{d}\right)^{n}}\right|_{i=1} ^{n} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{i}, x_{j}, y_{1}, y_{2},\right) G^{(n)}\left(x_{1}, \ldots, \underset{\hat{i}}{\left.y_{1}, \ldots,{\underset{\hat{j}}{ }}^{y_{2}}, \ldots, x_{n}\right) d y_{1} d y_{2}}\right. \\
& \quad-\sum_{i=1}^{n} \sum_{j=i+1}^{n} a_{1}\left(x_{i}, x_{j}\right) G^{(n)}\left(x^{(n)}\right)-\kappa G^{(n)}\left(x^{(n)}\right) \mid d x^{(n)},
\end{aligned}
$$

and, using the obvious inequality $\|f-g\|_{X_{n}} \geq\left|\|f\|_{X_{n}}-\|g\|_{X_{n}}\right|$, we continue

$$
\begin{aligned}
& \geq\left|\int_{\left(\mathbb{R}^{d}\right)^{n}}\right| \sum_{i=1}^{n} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{i}, x_{j}, y_{1}, y_{2},\right) \\
& \quad \times G^{(n)}\left(x_{1}, \ldots, y_{1}, \ldots, y_{2}, \ldots, x_{n}\right) d y_{1} d y_{2} \mid d x^{(n)} \\
& \quad
\end{aligned}
$$

Since $G^{(n)}$ is a symmetric function, we get

$$
\begin{aligned}
& \quad \int_{\left(\mathbb{R}^{d}\right)^{n}} \mid \sum_{i=1}^{n} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x_{i}, x_{j}, y_{1}, y_{2},\right) \\
& \quad \times G^{(n)}\left(x_{1}, \ldots, y_{1}, \ldots, y_{2}, \ldots, x_{n}\right) d y_{1} d y_{2} \mid d x^{(n)} \\
& \leq \frac{n(n-1)}{2} \int_{\left(\mathbb{R}^{d}\right)^{n}} a_{2}\left(y_{1}, y_{2}\right)\left|G^{(n)}\left(y_{1}, y_{2}, x_{1}, \ldots, x_{n-2}\right)\right| d y_{1} d y_{2} d x_{1} \ldots d x_{n-2} .
\end{aligned}
$$

Therefore, if (3.19) holds, then

$$
\left\|L_{0}^{(n)} G^{(n)}-\kappa G^{(n)}\right\|_{X_{n}} \geq \kappa\left\|G^{(n)}\right\|_{X_{n}}
$$

which proves the dissipativity of $L_{0}^{(n)}$, see e.g. [5, Proposition 3.23].
Remark 3.4. If the function $c$ satisfies the symmetry condition (1.4), then (3.19) trivially follows from (3.17).

In Theorems 3.5 and 3.6 below, we formulate general results which are applicable to our dynamics under assumptions (3.11)-(3.14), (3.19).
Theorem 3.5. Consider the initial value problem

$$
\begin{align*}
\frac{\partial}{\partial t} G_{t}(\eta) & =\left(L_{0} G_{t}\right)(\eta)+\left(W G_{t}\right)(\eta), \quad t>0, \eta \in \Gamma_{0}  \tag{3.21}\\
\left.G_{t}\right|_{t=0} & =G_{0}
\end{align*}
$$

Here, for any $G=\left(G^{(n)}\right)_{n \geq 0}$,

$$
\begin{aligned}
\left(L_{0} G\right)^{(n)} & =L_{0}^{(n)} G^{(n)}, n \geq 1 \\
(W G)^{(n)} & =W^{(n)} G^{(n-1)}, n \geq 2 \\
\left(L_{0} G\right)^{(0)} & =(W G)^{(0)}=(W G)^{(1)}=0
\end{aligned}
$$

Further suppose that $L_{0}^{(n)}$ is a bounded generator of a strongly continuous contraction semigroup $e^{t L_{0}^{(n)}}$ in $X_{n}$, while $W^{(n)}$ is a bounded operator from $X_{n-1}$ into $X_{n}$ whose norm satisfies

$$
\begin{equation*}
\left\|W^{(n)}\right\|_{X_{n-1} \rightarrow X_{n}} \leq \operatorname{Bn}(n-1), \quad n \geq 1 \tag{3.22}
\end{equation*}
$$

for some $B \geq 1$ which is independent of $n$.
Let $C>0$ and $G_{0} \in \mathcal{L}_{C}$. Then the initial value problem (3.21) has a unique solution $G_{t} \in \mathcal{L}_{\rho(t, C)}$, where

$$
\begin{equation*}
\rho(t, C):=\frac{C}{1+B C t} . \tag{3.23}
\end{equation*}
$$

Furthermore,

$$
\begin{equation*}
\left\|G_{t}\right\|_{\mathcal{L}_{\rho(t, C)}} \leq\left\|G_{0}\right\|_{\mathcal{L}_{C}} \tag{3.24}
\end{equation*}
$$

Proof. Let us rewrite the initial value problem (3.21) as an infinite system of differential equations. Namely, for any $n \geq 1$,

$$
\begin{equation*}
\frac{\partial}{\partial t} G_{t}^{(n)}\left(x^{(n)}\right)=\left(L_{0}^{(n)} G_{t}^{(n)}\right)\left(x^{(n)}\right)+\left(W^{(n)} G_{t}^{(n-1)}\right)\left(x^{(n)}\right) \tag{3.25}
\end{equation*}
$$

with $G_{t}^{(0)}=G_{0}^{(0)}$. This system may be solved recurrently: for each $n \geq 1$

$$
\begin{align*}
G_{t}^{(n)}\left(x^{(n)}\right)= & \left(e^{t L_{0}^{(n)}} G_{0}^{(n)}\right)\left(x^{(n)}\right) \\
& +\int_{0}^{t}\left(e^{(t-s) L_{0}^{(n)}} W^{(n)} G_{s}^{(n-1)}\right)\left(x^{(n)}\right) d s \tag{3.26}
\end{align*}
$$

Iterating (3.26), we obtain

$$
G_{t}^{(n)}\left(x^{(n)}\right)=\sum_{k=0}^{n}\left(V_{k, n}(t) G_{0}^{(n-k)}\right)\left(x^{(n)}\right)
$$

where $V_{k, n}(t): X_{n-k} \rightarrow X_{n}$ is given by

$$
\begin{aligned}
V_{k, n}(t):=\int_{0}^{t} & \int_{0}^{s_{1}} \cdots \int_{0}^{s_{k-1}} e^{\left(t-s_{1}\right) L_{0}^{(n)}} W^{(n)} e^{\left(s_{1}-s_{2}\right) L_{0}^{(n-1)}} W^{(n-1)} \cdots \\
& \times e^{\left(s_{k-1}-s_{k}\right) L_{0}^{(n-k+1)}} W^{(n-k+1)} e^{s_{k} L_{0}^{(n-k)}} d s_{k} \cdots d s_{1}
\end{aligned}
$$

for $2 \leq k \leq n-1$, and

$$
\begin{aligned}
V_{1, n}(t) G^{(n-1)} & :=\int_{0}^{t} e^{\left(t-s_{1}\right) L_{0}^{(n)}} W^{(n)} e^{s_{1} L_{0}^{(n-1)}} G^{(n-1)} d s_{1} \\
V_{0, n}(t) G^{(n)} & :=e^{t L_{0}^{(n)}} G^{(n)} \\
V_{n, n}(t) G^{(n)} & :=\chi_{\{n=0\}} G^{(0)}
\end{aligned}
$$

Hence, since (3.22) holds and since each operator $e^{s L_{0}^{(i)}}$, with $s \geq 0$, is a contraction in $X_{i}$, we get, for $1 \leq k \leq n-1, n \geq 2$,

$$
\begin{aligned}
\left\|V_{k, n}(t) G^{(n-k)}\right\|_{X_{n}} & \leq \frac{t^{k}}{k!} B^{k} n(n-1)(n-1)(n-2) \cdots(n-k+1)(n-k) \\
& =(t B)^{k} \frac{n!}{k!(n-k)!} \frac{(n-1)!}{(n-k-1)!}\left\|G^{(n-k)}\right\|_{X_{n-k}},
\end{aligned}
$$

and $\left\|V_{0, n}(t) G^{(n)}\right\|_{X_{n}} \leq\left\|G^{(n)}\right\|_{X_{n}}$ for $n \geq 1$. Therefore, for $n \geq 1$,

$$
\begin{align*}
\left\|G_{t}^{(n)}\right\|_{X_{n}} & \leq \sum_{k=0}^{n-1}(t B)^{k} \frac{n!}{k!(n-k)!} \frac{(n-1)!}{(n-k-1)!}\left\|G^{(n-k)}\right\|_{X_{n-k}} \\
& =\sum_{k=1}^{n}(t B)^{(n-k)} \frac{n!}{(n-k)!k!} \frac{(n-1)!}{(k-1)!}\left\|G_{0}^{(k)}\right\|_{X_{k}} \tag{3.27}
\end{align*}
$$

Then, for any $q(t)>0$,

$$
\begin{aligned}
& \left\|G_{t}\right\|_{\mathcal{L}_{C q(t)}} \\
& =\left|G_{t}^{(0)}\right|+\sum_{n=1}^{\infty} \frac{C^{n} q^{n}(t)}{n!}\left\|G_{t}^{(n)}\right\|_{X_{n}} \\
& \leq\left|G_{0}^{(0)}\right|+\sum_{n=1}^{\infty} \frac{C^{n} q^{n}(t)}{n!} \sum_{k=1}^{n}\left\|G_{0}^{(k)}\right\|_{X_{k}}(t B)^{n-k} \frac{n!}{k!(n-k)!} \frac{(n-1)!}{(k-1)!} \\
& =\left|G_{0}^{(0)}\right|+\sum_{k=1}^{\infty}\left\|G_{0}^{(k)}\right\|_{X_{k}} \frac{1}{k!} \sum_{n=k}^{\infty} C^{n} q^{n}(t)(t B)^{n-k} \frac{1}{(n-k)!} \frac{(n-1)!}{(k-1)!} \\
& =\left|G_{0}^{(0)}\right|+\sum_{k=1}^{\infty}\left\|G_{0}^{(k)}\right\|_{X_{k}} \frac{C^{k} q^{k}(t)}{k!} \sum_{n=0}^{\infty} q^{n}(t)(t B C)^{n} \frac{(n+k-1)!}{n!(k-1)!}
\end{aligned}
$$

Now, let $q(t)=\frac{1}{1+B C t}$. For any $x \in[0,1)$ and $m \in \mathbb{N}$,

$$
\left(\frac{1}{1-x}\right)^{m+1}=\sum_{n=0}^{\infty} x^{n} \frac{(n+m)!}{n!m!}
$$

Applying this equality to $x=q(t) B C t<1$ and $m=k-1$, we obtain

$$
q^{k}(t) \sum_{n=0}^{\infty} q^{n}(t)(t B C)^{n} \frac{(n+k-1)!}{n!(k-1)!}=\left(\frac{q(t)}{1-q(t) B C t}\right)^{k}=1
$$

Therefore,

$$
\left\|G_{t}\right\|_{\mathcal{L}_{C q(t)}} \leq\left|G_{0}^{(0)}\right|+\sum_{k=1}^{\infty}\left\|G_{0}^{(k)}\right\|_{X_{k}} \frac{C^{k}}{k!}=\left\|G_{0}\right\|_{\mathcal{L}_{C}}
$$

which proves the statement.
In fact, we have a linear evolution operator

$$
V(t): \mathcal{L}_{C} \rightarrow \mathcal{L}_{\rho(t, C)}
$$

satisfying $G_{t}=V(t) G_{0}$ and

$$
\begin{equation*}
\|V(t)\|_{\mathcal{L}_{C} \rightarrow \mathcal{L}_{\rho(t, C)}} \leq 1 \tag{3.28}
\end{equation*}
$$

Theorem 3.6. Let the conditions of Theorem 3.5 be satisfied. Further suppose that there exists $A>0$ such that $\left\|L_{0}^{(n)}\right\|_{X_{n} \rightarrow X_{n}} \leq A n(n-1), n \geq 1$. Let $C_{0}>0$, $k_{0} \in \mathcal{K}_{C_{0}}, T=\frac{1}{B C_{0}}$. Then for any $t \in(0, T)$, there exists $k_{t} \in \mathcal{K}_{C_{t}}$ with

$$
\begin{equation*}
C_{t}=\frac{C_{0}}{1-B C_{0} t} \tag{3.29}
\end{equation*}
$$

such that, for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$,

$$
\frac{d}{d t}\left\langle\left\langle G, k_{t}\right\rangle\right\rangle=\left\langle\left\langle\left(L_{0}+W\right) G, k_{t}\right\rangle\right\rangle, \quad t \in(0, T)
$$

Moreover, for any $t \in(0, T)$,

$$
\begin{equation*}
\left\|k_{t}\right\|_{\mathcal{K}_{C_{t}}} \leq\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} \tag{3.30}
\end{equation*}
$$

Proof. Let $t \in(0, T)$ be arbitrary. The function $f(x)=\rho(t, x)=\frac{x}{1+x B t}, x \geq 0$, increases to $\frac{1}{t B}$ as $x \rightarrow+\infty$. Since $C_{0}<\frac{1}{t B}$, there exists a unique solution to $f(x)=C_{0}$, namely, $x=C_{t}$, given by (3.29). Take any $G_{0} \in \mathcal{L}_{C_{t}}$. By Theorem 3.5, there exists an evolution $G_{0} \mapsto G_{\tau}$ for any $\tau>0$ such that $G_{\tau} \in \mathcal{L}_{\rho\left(\tau, C_{t}\right)}$. Consider this evolution at the moment $\tau=t$. Since

$$
\begin{equation*}
\rho\left(t, C_{t}\right)=\frac{C_{t}}{1+B C_{t} t}=C_{0} \tag{3.31}
\end{equation*}
$$

we have $G_{t} \in \mathcal{L}_{C_{0}}$. Therefore, $\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle$ is well-defined. Moreover, by (3.24),

$$
\begin{align*}
\left|\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle\right| & \leq\left\|G_{t}\right\|_{\mathcal{L}_{C_{0}}}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}}=\left\|G_{t}\right\|_{\mathcal{L}_{\rho\left(t, C_{t}\right)}}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} \\
& \leq\left\|G_{0}\right\|_{\mathcal{L}_{C_{t}}}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} . \tag{3.32}
\end{align*}
$$

Therefore, the mapping $G_{0} \mapsto\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle$ is a linear continuous functional on the space $\mathcal{L}_{C_{t}}$. Hence, there exists $k_{t} \in \mathcal{K}_{C_{t}}$ such that, for any $G_{0} \in \mathcal{L}_{C_{t}}$,

$$
\begin{equation*}
\left\langle\left\langle G_{0}, k_{t}\right\rangle\right\rangle=\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle V(t) G_{0}, k_{0}\right\rangle\right\rangle . \tag{3.33}
\end{equation*}
$$

We note that $k_{t}$ depends on $k_{0}$ and does not depend on $G_{0}$. Further, (3.32) implies (3.30).

Let now $G_{0} \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$. Consider a function $g=g_{G_{0}, k_{0}}:[0, T) \rightarrow \mathbb{R}$, $g(t):=\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle G_{0}, k_{t}\right\rangle\right\rangle$. We have

$$
\begin{equation*}
g(t)=\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle=\sum_{n=0}^{\infty} \frac{1}{n!} \int_{\left(\mathbb{R}^{d}\right)^{n}} G_{t}^{(n)}\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} \tag{3.34}
\end{equation*}
$$

By (3.32), for any $\left[0, T^{\prime}\right] \subset[0, T)$,

$$
\left|\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle\right| \leq\left\|G_{0}\right\|_{\mathcal{L}_{C^{\prime}}}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}}, \quad t \in\left[0, T^{\prime}\right] .
$$

Hence, the series (3.34) converges on $\left[0, T^{\prime}\right]$. Using the well-known representation

$$
\begin{equation*}
e^{t L_{0}^{(n)}} G_{0}=G_{0}+\int_{0}^{t} e^{s L_{0}^{(n)}} L_{0}^{(n)} G_{0} d s \tag{3.35}
\end{equation*}
$$

(see e.g. [5, Lemma 1.3 (iv)]), we derive from (3.26) and Fubini's theorem:

$$
\begin{align*}
g_{n}(t):= & \int_{\left(\mathbb{R}^{d}\right)^{n}} G_{t}^{(n)}\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} \\
= & \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(e^{t L_{0}^{(n)}} G_{0}^{(n)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} \\
& +\int_{0}^{t} \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(e^{(t-s) L_{0}^{(n)}} W^{(n)} G_{s}^{(n-1)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \\
= & \int_{\left(\mathbb{R}^{d}\right)^{n}} G_{0}^{(n)}\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} \\
& +\int_{0}^{t} \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(e^{s L_{0}^{(n)}} L_{0}^{(n)} G_{0}^{(n)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \\
& +\int_{0}^{t} \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(e^{(t-s) L_{0}^{(n)}} W^{(n)} G_{s}^{(n-1)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s . \tag{3.36}
\end{align*}
$$

Since $L_{0}^{(n)}: X_{n} \rightarrow X_{n}$ and $W^{(n)}: X_{n-1} \rightarrow X_{n}$ are bounded, the functions inside the time integrals are continuous in $s$. Therefore, by (3.36) and (3.26), $g_{n}(t)$ is differentiable on $(0, T)$ and

$$
\begin{align*}
g_{n}^{\prime}(t)= & \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(L_{0}^{(n)} e^{t L_{0}^{(n)}} G_{0}^{(n)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} \\
& +\int_{0}^{t} \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(L_{0}^{(n)} e^{(t-s) L_{0}^{(n)}} W^{(n)} G_{s}^{(n-1)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \\
& +\int_{\left(\mathbb{R}^{d}\right)^{n}}\left(W^{(n)} G_{t}^{(n-1)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \\
= & \int_{\left(\mathbb{R}^{d}\right)^{n}}\left(L_{0}^{(n)} G_{t}^{(n)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \\
& +\int_{\left(\mathbb{R}^{d}\right)^{n}}\left(W^{(n)} G_{t}^{(n-1)}\right)\left(x^{(n)}\right) k_{0}^{(n)}\left(x^{(n)}\right) d x^{(n)} d s \tag{3.37}
\end{align*}
$$

Hence, for any $n \geq 2$,

$$
\left|g_{n}^{\prime}(t)\right| \leq\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} C_{0}^{n} n(n-1)\left(A\left\|G_{t}\right\|_{X_{n}}+B\left\|G_{t}\right\|_{X_{n-1}}\right)
$$

Analogously to the proof of Theorem 3.5, we obtain, for all $t \in\left[0, T^{\prime}\right] \subset[0, T)$,

$$
\begin{aligned}
& \sum_{n=1}^{\infty} \frac{1}{n!}\left|g_{n}^{\prime}(t)\right| \\
\leq & \text { const } \cdot \sum_{n=1}^{\infty} \frac{1}{n!} C_{0}^{n} n(n-1) \sum_{k=1}^{n}(t B)^{(n-k)} \frac{n!}{(n-k)!k!} \frac{(n-1)!}{(k-1)!}\left\|G_{0}^{(k)}\right\|_{X_{k}} \\
\leq & \text { const } \cdot \sum_{k=1}^{\infty} \frac{1}{k!}\left\|G_{0}^{(k)}\right\|_{X_{k}} \sum_{n=k}^{\infty} C_{0}^{n}\left(T^{\prime} B\right)^{(n-k)} \frac{n(n-1)}{(n-k)!} \frac{(n-1)!}{(k-1)!}
\end{aligned}
$$

$$
\begin{equation*}
=\text { const } \cdot \sum_{k=1}^{\infty} \frac{C_{0}^{k}}{k!}\left\|G_{0}^{(k)}\right\|_{X_{k}} \sum_{n=0}^{\infty} C_{0}^{n}\left(T^{\prime} B\right)^{n} \frac{(n+k-1)}{n!} \frac{(n+k)!}{(k-1)!}<\infty \tag{3.38}
\end{equation*}
$$

since $G_{0} \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ (and so there exists $K \in \mathbb{N}$ such that $G_{0}^{(k)}=0$ for all $k \geq K$ ) and the inner series converges as $C_{0} T^{\prime} B<1$.

Hence, $g(t)$ is differentiable on any $\left[0, T^{\prime}\right] \subset[0, T)$. Next, (3.34), (3.37), and (3.38) imply that

$$
\begin{equation*}
g^{\prime}(t)=\frac{d}{d t}\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle\left(L_{0}+W\right) G_{t}, k_{0}\right\rangle\right\rangle \tag{3.39}
\end{equation*}
$$

and, moreover, $\left(L_{0}+W\right) G_{t} \in \mathcal{L}_{C_{0}}$. Therefore, using (3.33), (3.39) and the obvious inclusion $\left(L_{0}+W\right) G_{0} \in \mathcal{L}_{C_{t}}$, we obtain

$$
\begin{aligned}
\frac{d}{d t}\left\langle\left\langle G_{0}, k_{t}\right\rangle\right\rangle & =\frac{d}{d t}\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle\left(L_{0}+W\right) G_{t}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle\left(L_{0}+W\right) V(t) G_{0}, k_{0}\right\rangle\right\rangle \\
& =\left\langle\left\langle V(t)\left(L_{0}+W\right) G_{0}, k_{0}\right\rangle\right\rangle=\left\langle\left\langle\left(L_{0}+W\right) G_{0}, k_{t}\right\rangle\right\rangle
\end{aligned}
$$

provided

$$
\begin{equation*}
\left(L_{0}+W\right) V(t) G_{0}=V(t)\left(L_{0}+W\right) G_{0} \tag{3.40}
\end{equation*}
$$

To prove (3.40), we consider, for each $N \in \mathbb{N}$, the space $\mathbb{X}_{N}:=\bigoplus_{n=0}^{N} X_{n}$ with the norm $\|\cdot\|_{\mathbb{X}_{N}}:=\sum_{n=0}^{N}\|\cdot\|_{X_{n}}$. For any $\mathbb{G} \in \mathbb{X}_{N}, \mathbb{G}=\left(\mathbb{G}^{(0)}, \ldots, \mathbb{G}^{(N)}\right)$, we define the following function on $\Gamma_{0}$ :

$$
\mathbb{I}_{N} \mathbb{G}:=\left(\mathbb{G}^{(0)}, \ldots, \mathbb{G}^{(N)}, 0,0, \ldots\right)
$$

For any function $G$ on $\Gamma_{0}, G=\left(G^{(0)}, \ldots, G^{(n)}, \ldots\right)$, with $G^{(n)} \in X_{n}$, we define the following element of $\mathbb{X}_{N}$ :

$$
\mathbb{P}_{N} G:=\left(G^{(0)}, \ldots, G^{(n)}\right)
$$

The system of differential equations (3.25) for $1 \leq n \leq N$ can be considered as one equation $\frac{\partial}{\partial t} \mathbb{G}_{t}=\mathbb{L}_{N} \mathbb{G}_{t}$ in $\mathbb{X}_{N}$ with

$$
\mathbb{L}_{N}:=\mathbb{P}_{N}\left(L_{0}+W\right) \mathbb{I}_{N}
$$

Clearly, $\mathbb{L}_{N}$ is a bounded operator in $\mathbb{X}_{N}$. Hence, there exists a unique vectorvalued solution of this equation, $\mathbb{G}_{t}=e^{t \mathbb{L}_{N}} \mathbb{G}_{0}$. The $n$-th component of $\mathbb{G}_{t}$, i.e., $\mathbb{G}_{t}^{(n)}$, coincides with the $G_{t}^{(n)}$ obtained in Theorem 3.5, for each $0 \leq n \leq N$, where $G_{0}=\mathbb{I}_{N} \mathbb{G}_{0}$. More precisely, for $0 \leq n \leq N$,

$$
\begin{equation*}
\left(V(t) G_{0}\right)^{(n)}=\left(\mathbb{I}_{N} e^{t \mathbb{L}_{N}} \mathbb{G}_{0}\right)^{(n)}=\left(e^{t \mathbb{L}_{N}} \mathbb{G}_{0}\right)^{(n)}=\left(e^{t \mathbb{L}_{N}} \mathbb{P}_{N} G_{0}\right)^{(n)} \tag{3.41}
\end{equation*}
$$

It is well known that a bounded operator $\mathbb{L}_{N}$ commutes with its semigroup $e^{t \mathbb{L}_{N}}$. Note also that, for $0 \leq n \leq N, G^{(n)}=\left(\mathbb{P}_{N} G\right)^{(n)}$. Therefore, for all $N \geq 1$,
$0 \leq n \leq N$, and for $G_{0}=\mathbb{I}_{N} \mathbb{G}_{0}$, we obtain

$$
\begin{aligned}
\left(\left(L_{0}+W\right) V(t) G_{0}\right)^{(n)} & =\left(\mathbb{P}_{N}\left(L_{0}+W\right) V(t) G_{0}\right)^{(n)}=\left(\mathbb{L}_{N} e^{t \mathbb{L}_{N}} \mathbb{G}_{0}\right)^{(n)} \\
& =\left(e^{t \mathbb{L}_{N}} \mathbb{L}_{N} \mathbb{G}_{0}\right)^{(n)}=\left(e^{t \mathbb{L}_{N}} \mathbb{P}_{N}\left(L_{0}+W\right) \mathbb{I}_{N} \mathbb{G}_{0}\right)^{(n)} \\
& =\left(V(t)\left(L_{0}+W\right) G_{0}\right)^{(n)},
\end{aligned}
$$

where in the last equality we applied (3.41) for $\left(L_{0}+W\right) G_{0}$ instead of $G_{0}$. Hence, (3.40) holds.
Remark 3.7. Note that the initial value problem $\frac{\partial}{\partial t} k_{t}=\widehat{L}^{*} k_{t},\left.k_{t}\right|_{t=t_{1}}=k_{t_{1}} \in$ $\mathcal{K}_{C_{t_{1}}}$ for some $t_{1}<T=\frac{1}{B C_{0}}$, has a solution only on the time interval $\left[t_{1}, t_{1}+\right.$ $\left.T_{1}\right)=\left[t_{1}, T\right)$, since $T_{1}=\frac{1}{B C_{t_{1}}}=\frac{1-B C_{0} t_{1}}{B C_{0}}=T-t_{1}$.
Remark 3.8. Using an estimate analogous to (3.38), one can show that $\frac{\partial G_{t}}{\partial t} \in$ $\mathcal{L}_{\rho(t, C)}$ if $G_{0}$ belongs to $B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ (or even to a larger subset of $\mathcal{L}_{C}$ ).

Thus, by Theorems 3.5 and 3.6, under conditions (3.11)-(3.14), (3.19) for the binary jumps dynamics with generator (1.2) we have the evolution of quasiobservables and the corresponding dual one. We will now show that the latter evolution generates an evolution of probability measures on $\Gamma$.

Theorem 3.9. Let (3.11)-(3.14) and (3.19) hold. Fix a measure $\mu \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ which has a correlation functional $k_{\mu} \in \mathcal{K}_{C_{0}}, C_{0}>0$. Consider the evolution $k_{\mu} \mapsto k_{t} \in \mathcal{K}_{C_{t}}, t \in(0, T)$, where $T=1 /\left(c_{3}+c_{4}\right) C_{0}$. Then, for any $t \in$ $(0, T)$, there exists a unique measure $\mu_{t} \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ such that $k_{t}$ is the correlation functional of $\mu_{t}$.

Proof. We first recall the following definition. Let a measurable, non-negative function $k$ on $\Gamma_{0}$ be such that $\int_{M} k(\eta) d \lambda(\eta)<\infty$ for any bounded $M \in \mathcal{B}\left(\Gamma_{0}\right)$. The function $k$ is said to be Lenard positive definite if $\langle\langle G, k\rangle\rangle \geq 0$ for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ such that $K G \geq 0$. It was shown in [24] that any such $k$ is the correlation functional of some probability measure on $\Gamma$. If, additionally, $k \in \mathcal{K}_{C}$ for some $C>0$, then this measure is uniquely defined (cf. [22]) and belongs to $\mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ (cf. [15]). Therefore, to prove the theorem, it is enough to show that $k_{t}$ is Lenard positive definite for any $t \in(0, T)$.

Since the measure $\mu \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ has the correlation functional $k_{\mu}, \mu$ is locally absolutely continuous with respect to $\pi$, and for any $\Lambda \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ and $\lambda$-a.a $\eta \in \Gamma(\Lambda)$

$$
\begin{equation*}
\frac{d \mu^{\Lambda}}{d \lambda}(\eta)=\int_{\Gamma(\Lambda)}(-1)^{|\xi|} k_{\mu}(\eta \cup \xi) d \lambda(\xi) \tag{3.42}
\end{equation*}
$$

see [15, Proposition 4.3]. Since $k_{\mu} \in \mathcal{K}_{C_{0}}$, we have, by (3.42), (3.4), and (2.3),

$$
\begin{equation*}
\frac{d \mu^{\Lambda}}{d \lambda}(\eta) \leq\left\|k_{\mu}\right\|_{\mathcal{K}_{C_{0}}} e^{C_{0} m(\Lambda)} C_{0}^{|\eta|} \tag{3.43}
\end{equation*}
$$

for $\lambda$-a.a $\eta \in \Gamma(\Lambda)$.
We fix $\Lambda_{0} \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$ and consider the projection $\mu_{0}:=\mu^{\Lambda_{0}}$ on $\Gamma\left(\Lambda_{0}\right)$. By (3.43), for $\lambda$-a.a. $\eta \in \Gamma\left(\Lambda_{0}\right)$

$$
\begin{equation*}
R_{0}(\eta):=\frac{d \mu_{0}}{d \lambda}(\eta) \leq A_{0} C_{0}^{|\eta|} \tag{3.44}
\end{equation*}
$$

where $A_{0}:=\left\|k_{\mu}\right\|_{\mathcal{K}_{C_{0}}} e^{C_{0} m\left(\Lambda_{0}\right)}$. Clearly, $\mu_{0}$ may be considered as a measure on the whole of $\Gamma$ if we set $R_{0}$ to be equal to 0 outside of $\Gamma\left(\Lambda_{0}\right)$. Hence,

$$
\mu_{0}(A)=\int_{\Gamma\left(\Lambda_{0}\right) \cap A} R_{0}(\eta) d \lambda(\eta), \quad A \in \mathcal{B}(\Gamma)
$$

On the other hand, $R_{0}$ being extended by zero outside of $\Gamma\left(\Lambda_{0}\right)$ can also be regarded as a $\mathcal{B}\left(\Gamma_{0}\right)$-measurable function. Evidently, that in this case $0 \leq R_{0} \in$ $L^{1}\left(\Gamma_{0}, d \lambda\right)$ with $\int_{\Gamma_{0}} R_{0} d \lambda=1$. Note that

$$
\begin{equation*}
k_{0}:=\mathbb{1}_{\Gamma\left(\Lambda_{0}\right)} k_{\mu} \in \mathcal{K}_{C_{0}} \tag{3.45}
\end{equation*}
$$

is the correlation functional of $\mu_{0}$. Here and below, $\mathbb{1}_{\Delta}$ stands for the indicator function of a set $\Delta$. By [15, Proposition 4.2], for $\lambda$-a.a. $\eta \in \Gamma\left(\Lambda_{0}\right)$

$$
\begin{equation*}
k_{0}(\eta)=\int_{\Gamma\left(\Lambda_{0}\right)} R_{0}(\eta \cup \xi) d \lambda(\xi) \tag{3.46}
\end{equation*}
$$

There exists an $N_{0} \in \mathbb{N}$ such that $\int_{\left(\mathbb{R}^{d}\right)^{N_{0}}} R_{0}^{\left(N_{0}\right)} d x^{\left(N_{0}\right)}>0$ (otherwise $R_{0}=0$ $\lambda$-a.e.). We set

$$
r:=\int_{\bigsqcup_{n=0}^{N_{0}} \Gamma^{(n)}} R_{0}(\eta) d \lambda(\eta) \in(0,1]
$$

For each $N \geq N_{0}$, we define

$$
\begin{equation*}
R_{0, N}(\eta)=\mathbb{1}_{\{|\eta| \leq N\}}(\eta) R_{0}(\eta)\left(\int_{\bigsqcup_{n=0}^{N} \Gamma^{(n)}} R_{0}(\eta) d \lambda(\eta)\right)^{-1} \tag{3.47}
\end{equation*}
$$

Then, clearly, $0 \leq R_{0, N} \in L^{1}\left(\Gamma_{0}, d \lambda\right)$, with $\int_{\Gamma_{0}} R_{0, N} d \lambda=1$. Moreover, $R_{0, N}$ has a bounded support on $\Gamma_{0}$. By (3.47) and (3.44) we have

$$
\begin{equation*}
R_{0, N}(\eta) \leq r^{-1} R_{0}(\eta) \leq r^{-1} A_{0} C_{0}^{|\eta|} \tag{3.48}
\end{equation*}
$$

for $\lambda$-a.a. $\eta \in \Gamma_{0}$.
We define a probability measure $\mu_{0, N} \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$, concentrated on $\Gamma_{0}$, by $d \mu_{0, N}=R_{0, N} d \lambda$. By [15, Proposition 4.2], the correlation functional $k_{0, N}$ of $\mu_{0, N}$ has the following representation

$$
\begin{equation*}
k_{0, N}(\eta)=\int_{\Gamma\left(\Lambda_{0}\right)} R_{0, N}(\eta \cup \xi) d \lambda(\xi) \tag{3.49}
\end{equation*}
$$

for $\lambda$-a.a. $\eta \in \Gamma\left(\Lambda_{0}\right)$. It is evident now that $k_{0, N}$ has a bounded support on $\Gamma_{0}$. Moreover, by (3.46), (3.49), and the first inequality in (3.48), we get $k_{0, N} \leq \frac{1}{r} k_{0} \in \mathcal{K}_{C_{0}}$ and

$$
\begin{equation*}
\left\|k_{0, N}\right\|_{\mathcal{K}_{C_{0}}} \leq \frac{1}{r}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} \tag{3.50}
\end{equation*}
$$

By the definition of a correlation functional, for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$

$$
\begin{align*}
\left|\left\langle\left\langle G, k_{0}\right\rangle\right\rangle-\left\langle\left\langle G, k_{0, N}\right\rangle\right\rangle\right| & =\left|\int_{\Gamma_{0}}(K G)(\eta)\left(R_{0}(\eta)-R_{0, N}(\eta)\right) d \lambda(\eta)\right| \\
& \leq D \int_{\Gamma\left(\Lambda_{0}\right)}(1+|\eta|)^{M}\left|R_{0}(\eta)-R_{0, N}(\eta)\right| d \lambda(\eta) \tag{3.51}
\end{align*}
$$

for some $D=D(G)>0$ and $M=M(G) \in \mathbb{N}$ (see [15, Proposition 3.1]). By (3.47), $R_{0, N}(\eta) \rightarrow R_{0}(\eta)$ for $\lambda$-a.a. $\eta \in \Gamma\left(\Lambda_{0}\right)$. Furthermore, by (3.44) and (3.48),

$$
\left|R_{0}(\eta)-R_{0, N}(\eta)\right| \leq A_{0}\left(1+r^{-1}\right) C_{0}^{|\eta|}
$$

By (2.3),

$$
\int_{\Gamma\left(\Lambda_{0}\right)}(1+|\eta|)^{M} C_{0}^{|\eta|} d \lambda(\eta)<\infty
$$

Therefore, by the dominated convergence theorem, (3.51) yields

$$
\begin{equation*}
\lim _{N \rightarrow \infty}\left\langle\left\langle G, k_{0, N}\right\rangle\right\rangle=\left\langle\left\langle G, k_{0}\right\rangle\right\rangle . \tag{3.52}
\end{equation*}
$$

As before, we identify a function $F$ on $\Gamma_{0}$ with a sequence of symmetric functions $F^{(n)}$ on $\left(\mathbb{R}^{d}\right)^{n}, n \in \mathbb{N}_{0}$. Fix any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ and let $F$ be the restriction of $K G$ to $\Gamma_{0}$. Then there exist $\Lambda=\Lambda_{F} \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right), M=M_{F} \in \mathbb{N}$, and $D=D_{F}>0$ such that for all $\eta \in \Gamma_{0}$,

$$
|F(\eta)|=|F(\eta \cap \Lambda)| \leq D(1+|\eta \cap \Lambda|)^{M}
$$

(see [15, Proposition 3.1]). In particular, $F^{(n)}$ is bounded on $\left(\mathbb{R}^{d}\right)^{n}$ for each $n$. We restrict the operator $L$ given by (1.2) to functions on $\Gamma_{0}$. This restriction, $L_{0}$, is given by (3.7).

We define, for any $R \in L^{1}\left(\Gamma_{0}, \lambda\right)$, the function $L_{0}^{*} R$ on $\Gamma_{0}$ by $\left(L_{0}^{*} R\right)^{(n)}:=$ $\left(L_{0}^{(n)}\right)^{*} R^{(n)}$, where $\left(L_{0}^{(n)}\right)^{*}$ is given by the right hand side of (3.9) in which $c\left(x_{i}, x_{j}, y_{1}, y_{2}\right)$ is replaced by $c\left(y_{1}, y_{2}, x_{i}, x_{j}\right)$. Analogously to the proof of Proposition 3.3, we conclude that $\left(L_{0}^{(n)}\right)^{*}$ is a bounded generator of a strongly continuous semigroup on $X_{n}=L^{1}\left(\left(\mathbb{R}^{d}\right)^{n}, d x^{(n)}\right)$. In the dual space $X_{n}^{*}:=$ $L^{\infty}\left(\left(\mathbb{R}^{d}\right)^{n}, d x^{(n)}\right)$, we consider the dual operator to $\left(L_{0}^{(n)}\right)^{*}$, which is just the $L_{0}^{(n)}$ given by (3.9). It is easy to see that, under condition $(3.11), L_{0}^{(n)}$ is a bounded operator on $X_{n}^{*}$. Note that $L_{0}^{(n)} 1=0$ implies

$$
\int_{\left(\mathbb{R}^{d}\right)^{n}} e^{t\left(L_{0}^{(n)}\right)^{*}} R^{(n)} d x^{(n)}=\int_{\left(\mathbb{R}^{d}\right)^{n}}\left(e^{t L_{0}^{(n)}} 1\right) R^{(n)} d x^{(n)}=\int_{\left(\mathbb{R}^{d}\right)^{n}} R^{(n)} d x^{(n)}
$$

To show that $e^{t\left(L_{0}^{(n)}\right)^{*}}$ preserves the cone $X_{n}^{+}$of all positive functions in $X_{n}$, we write $\left(L_{0}^{(n)}\right)^{*}=L_{1}+L_{2}$, where

$$
\begin{aligned}
&\left(L_{1} R^{(n)}\right)\left(x^{(n)}\right)=\sum_{i=1}^{n} \sum_{j=i+1}^{n} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(y_{1}, y_{2}, x_{i}, x_{j}\right) \\
& \times R^{(n)}\left(x_{1}, \ldots, y_{1}, \ldots, y_{2}, \ldots, x_{n}\right) d y_{1} d y_{2} \\
& \hat{i}
\end{aligned}
$$

and

$$
\left(L_{2} R^{(n)}\right)\left(x^{(n)}\right)=-\left(\sum_{i=1}^{n} \sum_{j=i+1}^{n} a_{2}\left(x_{i}, x_{j}\right)\right) R^{(n)}\left(x^{(n)}\right)
$$

Clearly, $L_{1}$ and $L_{2}$ are bounded operators on $X_{n}$. Since $c \geq 0, L_{1}$ preserves the cone $X_{n}^{+}$, hence so does the semigroup $e^{t L_{1}}$. Since $L_{2}$ is a bounded multiplication operator, the semigroup $e^{t L_{2}}$ is a positive multiplication operator in $X_{n}$. Therefore, $e^{t\left(L_{0}^{(n)}\right)^{*}}$ preserves $X_{n}^{+}$by the Lie-Trotter product formula.

Therefore, if we define functions $R_{t, N}, N \geq N_{0}$ (cf. (3.47)) on $\Gamma_{0}$ by

$$
R_{t, N}^{(n)}:=e^{t\left(L_{0}^{(n)}\right)^{*}} R_{0, N}^{(n)}
$$

then $0 \leq R_{t, N} \in L^{1}\left(\Gamma_{0}, d \lambda\right)$ with $\int_{\Gamma_{0}} R_{t, N} d \lambda=1$. Note that $R_{t, N}^{(n)} \equiv 0$ for $n>N$. Therefore, we can define a measure $\tilde{\mu}_{t, N} \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$, concentrated on $\Gamma_{0}$, by $d \tilde{\mu}_{t, N}=R_{t, N} d \lambda$ (in fact, the measure $\tilde{\mu}_{t, N}$ is concentrated on $\bigsqcup_{n=0}^{N} \Gamma^{(n)}$ ). We denote by $\tilde{k}_{t, N}$ the correlation functional of $\tilde{\mu}_{t, N}$.

For each function $G$ on $\Gamma_{0}$ we define $K_{0} G:=(K G) \Gamma_{\Gamma_{0}}$. Take any $G_{0} \in$ $B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ such that $K G_{0} \geq 0$ on $\Gamma$. We denote $F_{0}:=K_{0} G_{0} \geq 0$ on $\Gamma_{0}$. We have, by the definition of a correlation functional,

$$
\begin{equation*}
\left\langle\left\langle G_{0}, k_{t, N}\right\rangle\right\rangle=\left\langle K_{0} G_{0}, R_{t, N}\right\rangle \geq 0 \tag{3.53}
\end{equation*}
$$

On the other hand, if we define a function $U(t) F_{0}$ on $\Gamma_{0}$ by $\left(U(t) F_{0}\right)^{(n)}=$ $e^{t L_{0}^{(n)}} F_{0}^{(n)}$, we obtain

$$
\begin{align*}
& \left\langle F_{0}, R_{t, N}\right\rangle=\sum_{n=0}^{N} \frac{1}{n!}\left\langle F_{0}^{(n)}, R_{t, N}^{(n)}\right\rangle=\sum_{n=0}^{N} \frac{1}{n!}\left\langle F_{0}^{(n)}, e^{t\left(L_{0}^{(n)}\right)^{*}} R_{0, N}^{(n)}\right\rangle \\
= & \sum_{n=0}^{N} \frac{1}{n!}\left\langle e^{t L_{0}^{(n)}} F_{0}^{(n)}, R_{0, N}^{(n)}\right\rangle=\left\langle U(t) F_{0}, R_{0, N}\right\rangle \\
= & \left\langle\left\langle K_{0}^{-1} U(t) K_{0} G_{0}, k_{0, N}\right\rangle\right\rangle . \tag{3.54}
\end{align*}
$$

It is evident, by Proposition 3.1, that

$$
\begin{equation*}
\left(K_{0}^{-1} U(t) K_{0} G_{0}\right)^{(n)}=e^{t\left(L_{0}^{(n)}+W^{(n)}\right)} G_{0}^{(n)}=\left(V(t) G_{0}\right)^{(n)}, \tag{3.55}
\end{equation*}
$$

where $V(t)$ is as in (3.28).

As a result, from (3.53)-(3.55), we get

$$
\begin{equation*}
\left\langle\left\langle G_{0}, k_{t, N}\right\rangle\right\rangle=\left\langle\left\langle G_{t}, k_{0, N}\right\rangle\right\rangle, \tag{3.56}
\end{equation*}
$$

where $G_{t}=V(t) G_{0}$. $\underset{\sim}{\text { By }}$ y $(3.45), k_{0} \in \mathcal{K}_{C_{0}}$. Hence, by Theorem 3.6, for any $t \in(0, T)$, there exists $\tilde{k}_{t} \in \mathcal{K}_{C_{t}}$ such that

$$
\begin{equation*}
\left\langle\left\langle G_{0}, \tilde{k}_{t}\right\rangle\right\rangle=\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle . \tag{3.57}
\end{equation*}
$$

Note that here, for a given $t \in(0, T)$, we may consider $G_{0} \in B_{\mathrm{bs}}\left(\Gamma_{0}\right) \subset \mathcal{L}_{C_{t}}$, where $C_{t}$ is given by (3.29). Then, by the proof of Theorem 3.6, $G_{t}=V(t) G_{0} \in$ $\mathcal{L}_{C_{0}}$. By (3.56) and (3.57), to prove that

$$
\begin{equation*}
\left\langle\left\langle G_{0}, \tilde{k}_{t}\right\rangle\right\rangle=\lim _{N \rightarrow \infty}\left\langle\left\langle G_{0}, k_{t, N}\right\rangle\right\rangle \tag{3.58}
\end{equation*}
$$

we only need to show that

$$
\begin{equation*}
\lim _{N \rightarrow \infty}\left\langle\left\langle G_{t}, k_{0, N}\right\rangle\right\rangle=\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle \tag{3.59}
\end{equation*}
$$

The latter fact is a direct consequence of (3.52) if we take into account that $G_{t} \in \mathcal{L}_{C_{0}}$ and that the set $B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ is dense in $\mathcal{L}_{C_{0}}$. Indeed, let us consider, for a fixed $t \in(0, T)$ and for any $\varepsilon>0$, a function $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ such that $\left\|G-G_{t}\right\|_{\mathcal{L}_{C_{0}}}<\varepsilon$. Then, by (3.52), there exists an $N_{1} \geq N_{0}$, such that, for any $N \geq N_{1}$

$$
\left|\left\langle\left\langle G, k_{0, N}\right\rangle\right\rangle-\left\langle\left\langle G, k_{0}\right\rangle\right\rangle\right|<\varepsilon .
$$

Therefore, by (3.50), for any $N \geq N_{1}$

$$
\begin{aligned}
& \left|\left\langle\left\langle G_{t}, k_{0, N}\right\rangle\right\rangle-\left\langle\left\langle G_{t}, k_{0}\right\rangle\right\rangle\right| \\
\leq & \left\|G-G_{t}\right\|_{\mathcal{L}_{C_{0}}}\left\|k_{0, N}\right\|_{\mathcal{K}_{C_{0}}}+\left|\left\langle\left\langle G, k_{0, N}\right\rangle\right\rangle-\left\langle\left\langle G, k_{0}\right\rangle\right\rangle\right|+\left\|G-G_{t}\right\|_{\mathcal{L}_{C_{0}}}\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}} \\
< & \varepsilon\left(r^{-1}+1\right)\left\|k_{0}\right\|_{\mathcal{K}_{C_{0}}}+\varepsilon,
\end{aligned}
$$

which proves (3.59). Therefore, (3.58) holds. Hence, by (3.53), $\tilde{k}_{t}=: \tilde{k}_{t}^{\Lambda_{0}}$ is Lenard positive definite for any $t \in(0, T)$.

As a result, for each $\Lambda \in \mathcal{B}_{\mathrm{b}}\left(\mathbb{R}^{d}\right)$, the evolution $k_{\mu^{\Lambda}} \mapsto \tilde{k}_{t}^{\Lambda}, t \in(0, T)$, preserves positive-definiteness and $\left\langle\left\langle G_{0}, \tilde{k}_{t}^{\Lambda}\right\rangle\right\rangle=\left\langle\left\langle G_{t}, k_{\mu^{\wedge}}\right\rangle\right\rangle$. On the other hand, by Theorem 3.6, we have the evolution $k_{\mu} \mapsto k_{t}, t \in(0, T)$ satisfying $\left\langle\left\langle G_{0}, k_{t}\right\rangle\right\rangle=$ $\left\langle\left\langle G_{t}, k_{\mu}\right\rangle\right\rangle$.

Since $k_{\mu^{\Lambda}}=\mathbb{1}_{\Gamma(\Lambda)} k_{\mu}$, it is evident that, for any $t \in(0, T),\left\langle\left\langle G_{t}, k_{\mu^{\wedge}}\right\rangle\right\rangle \rightarrow$ $\left\langle\left\langle G_{t}, k_{\mu}\right\rangle\right\rangle$ as $\Lambda \nearrow \mathbb{R}^{d}$. Therefore,

$$
\left\langle\left\langle G_{0}, k_{t}\right\rangle\right\rangle=\lim _{\Lambda \nearrow \mathbb{R}^{d}}\left\langle\left\langle G_{0}, \tilde{k}_{t}^{\Lambda}\right\rangle\right\rangle \geq 0
$$

Hence, for each $t \in(0, T)$, there exists a unique measure $\mu_{t} \in \mathcal{M}_{\mathrm{fm}}^{1}(\Gamma)$ whose correlation functional is $k_{t}$.

## 4 Vlasov-type scaling

For the reader's convenience, we start with explaining the idea of the Vlasovtype scaling. A general scheme for both birth-and-death and conservative dynamics may be found in [7]. Certain realizations of this approach were studied in $[4,6,8,9]$.

We would like to construct a scaling of the generator $L$, say $L_{\varepsilon}, \varepsilon>0$, such that the following requirements are satisfied. Assume we have an evolution $V_{\varepsilon}(t)$ corresponding to the equation $\frac{\partial}{\partial t} G_{t, \varepsilon}=\widehat{L}_{\varepsilon} G_{t, \varepsilon}$. Assume that, in some functional space, we have the dual evolution $V_{\varepsilon}^{*}(t)$ with respect to the duality (2.8). Let us choose an initial function for this dual evolution with a big singularity in $\varepsilon$, namely, $k_{0}^{(\varepsilon)}(\eta) \sim \varepsilon^{-|\eta|} r_{0}(\eta)$ as $\varepsilon \rightarrow 0\left(\eta \in \Gamma_{0}\right)$, with some function $r_{0}$, being independent of $\varepsilon$. Our first requirement on the scaling $L \mapsto L_{\varepsilon}$ is that the evolution $V_{\varepsilon}^{*}(t)$ preserves the order of the singularity:

$$
\begin{equation*}
\left(V_{\varepsilon}^{*}(t) k_{0}^{(\varepsilon)}\right)(\eta) \sim \varepsilon^{-|\eta|} r_{t}(\eta) \quad \text { as } \varepsilon \rightarrow 0, \quad \eta \in \Gamma_{0} \tag{4.1}
\end{equation*}
$$

where $r_{t}$ is such that the dynamics $r_{0} \mapsto r_{t}$ preserves the so-called LebesguePoisson exponents. Namely, if $r_{0}(\eta)=e_{\lambda}\left(p_{0}, \eta\right):=\prod_{x \in \eta} p_{0}(x)$, then $r_{t}(\eta)=$ $e_{\lambda}\left(p_{t}, \eta\right)=\prod_{x \in \eta} p_{t}(x)$. (Here, $\prod_{x \in \emptyset}:=1$.) Furthermore, we require that the $p_{t}$ 's satisfy a (nonlinear, in general) differential equation

$$
\begin{equation*}
\frac{\partial}{\partial t} p_{t}(x)=v\left(p_{t}\right)(x) \tag{4.2}
\end{equation*}
$$

which will be called a Vlasov-type equation.
For any $c>0$, we set $\left(R_{c} G\right)(\eta)=c^{|\eta|} G(\eta)$. Roughly speaking, (4.1) means that

$$
R_{\varepsilon} V_{\varepsilon}^{*}(t) R_{\varepsilon^{-1}} r_{0} \sim r_{t} .
$$

This gives us a hint to consider the map $R_{\varepsilon} \widehat{L}_{\varepsilon}^{*} R_{\varepsilon^{-1}}$, which is dual to

$$
\widehat{L}_{\varepsilon, \text { ren }}=R_{\varepsilon^{-1}} \widehat{L}_{\varepsilon} R_{\varepsilon} .
$$

Remark 4.1. We expect that the limiting dynamics for $R_{\varepsilon} V_{\varepsilon}^{*}(t) R_{\varepsilon^{-1}}$ will preserve the Lebesgue-Poisson exponents. Note that the Lebesgue-Poisson exponent $e_{\lambda}\left(p_{t}\right), t \geq 0$, is the correlation functional of the Poisson measure $\pi_{p_{t}}$ on $\Gamma$ with the non-constant intensity $p_{t}$ (for a rigorous definition of such a Poisson measure, see e.g. [1]). Therefore, at least heuristically, we expect to have the limiting dynamics: $\pi_{p_{0}} \mapsto \pi_{p_{t}}$, where $p_{t}$ satisfies the equation (4.2).

Below we will realize this scheme in the case of the generator $L$ given by (1.2). We will consider, for any $\varepsilon>0$, the scaled operator $L_{\varepsilon}=\varepsilon L$. Then, obviously, $\widehat{L}_{\varepsilon} G=\varepsilon \widehat{L} G=\varepsilon L_{0} G+\varepsilon W G$, where $L_{0}$ and $W$ are given by Proposition 3.1.

Proposition 4.2. For any $\varepsilon>0$

$$
\begin{equation*}
\widehat{L}_{\varepsilon, \text { ren }}=\varepsilon L_{0}+W \text {. } \tag{4.3}
\end{equation*}
$$

Moreover, let (3.11)-(3.14), (3.19) hold. Then, for any $C>0$, the initial value problem

$$
\begin{align*}
\frac{\partial}{\partial t} G_{t, \varepsilon}(\eta) & =\left(\widehat{L}_{\varepsilon, \text { ren }} G_{t, \varepsilon}\right)(\eta)  \tag{4.4}\\
\left.G_{t, \varepsilon}\right|_{t=0} & =G_{0, \varepsilon} \in \mathcal{L}_{C}
\end{align*}
$$

has a unique solution $G_{t, \varepsilon} \in \mathcal{L}_{\rho(t, C)}$, where $\rho(t, C)$ is given by (3.23).
Proof. By Proposition 3.1, $L_{0} R_{\varepsilon}=R_{\varepsilon} L_{0}$ and $W R_{\varepsilon}=\varepsilon^{-1} R_{\varepsilon} W$. Therefore,

$$
\begin{aligned}
\widehat{L}_{\varepsilon, \mathrm{ren}} & =R_{\varepsilon^{-1}} \varepsilon L_{0} R_{\varepsilon}+R_{\varepsilon^{-1}} \varepsilon W R_{\varepsilon} \\
& =\varepsilon R_{\varepsilon^{-1}} R_{\varepsilon} L_{0}+\varepsilon^{-1} \varepsilon R_{\varepsilon^{-1}} R_{\varepsilon} W=\varepsilon L_{0}+W
\end{aligned}
$$

By Proposition 3.3, $\left(\varepsilon L_{0}\right)^{(n)}$ is a generator of a contraction semigroup in $X_{n}$ for any $n \geq 1$. Hence, the statement is a direct consequence of Theorem 3.5. Note that the solution of (4.4) can be found recursively:

$$
\begin{align*}
G_{t, \varepsilon}^{(n)}\left(x^{(n)}\right)=( & \left.e^{\varepsilon t L_{0}^{(n)}} G_{0, \varepsilon}^{(n)}\right)\left(x^{(n)}\right) \\
& +\int_{0}^{t}\left(e^{\varepsilon(t-s) L_{0}^{(n)}} W^{(n)} G_{s, \varepsilon}^{(n-1)}\right)\left(x^{(n)}\right) d s, n \geq 1 \tag{4.5}
\end{align*}
$$

and $G_{t, \varepsilon}^{(0)}=G_{0, \varepsilon}^{(0)}$.
By (4.3), $\widehat{L}_{\varepsilon, \text { ren }} G(\eta) \rightarrow W G(\eta)$ as $\varepsilon \rightarrow 0\left(\eta \in \Gamma_{0}\right)$. Let (3.13)-(3.14) hold. By Theorem 3.5, for any $C>0$, the initial value problem

$$
\begin{align*}
\frac{\partial}{\partial t} G_{t, V}(\eta) & =\left(W G_{t, V}\right)(\eta)  \tag{4.6}\\
\left.G_{t, V}\right|_{t=0} & =G_{0, V} \in \mathcal{L}_{C}
\end{align*}
$$

has a unique solution $G_{t, V} \in \mathcal{L}_{\rho(t, C)}$, with $\rho(t, C)$ given by (3.23). This solution can be constructed recursively, namely, $G_{t, V}^{(0)}=G_{0, V}^{(0)}$ and

$$
\begin{equation*}
G_{t, V}^{(n)}\left(x^{(n)}\right)=G_{0, V}^{(n)}\left(x^{(n)}\right)+\int_{0}^{t}\left(W^{(n)} G_{s, V}^{(n-1)}\right)\left(x^{(n)}\right) d s, n \geq 1 \tag{4.7}
\end{equation*}
$$

Theorem 4.3. Suppose that conditions (3.11)-(3.14) and (3.19) hold. Let $C>$ 0 and let $\left\{G_{0, V}, G_{0, \varepsilon}, \varepsilon>0\right\} \subset \mathcal{L}_{C}$ be such that

$$
\begin{equation*}
\left\|G_{0, \varepsilon}-G_{0, V}\right\|_{\mathcal{L}_{C}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.8}
\end{equation*}
$$

Then, for any $T>0$ and any $r<\rho(T, C)$,

$$
\begin{equation*}
\sup _{t \in[0, T]}\left\|G_{t, \varepsilon}-G_{t, V}\right\|_{\mathcal{L}_{r}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.9}
\end{equation*}
$$

Proof. By (4.5) and (4.7), we have

$$
\begin{aligned}
& \left\|G_{t, \varepsilon}^{(n)}-G_{t V}^{(n)}\right\|_{X_{n}} \\
\leq & \left\|e^{\varepsilon \varepsilon L_{0}^{(n)}} G_{0, \varepsilon}^{(n)}-G_{0, V}^{(n)}\right\|_{X_{n}} \\
& +\int_{0}^{t}\left\|e^{\varepsilon(t-s) L_{0}^{(n)}} W^{(n)} G_{s, \varepsilon}^{(n-1)}-W^{(n)} G_{s, V}^{(n-1)}\right\|_{X_{n}} d s \\
\leq & \left\|e^{\varepsilon t L_{0}^{(n)}}\left(G_{0, \varepsilon}^{(n)}-G_{0, V}^{(n)}\right)\right\|_{X_{n}}+\left\|\left(e^{\varepsilon t L_{0}^{(n)}}-1\right) G_{0, V}^{(n)}\right\|_{X_{n}} \\
& +\int_{0}^{t}\left\|e^{\varepsilon(t-s) L_{0}^{(n)}} W^{(n)}\left(G_{s, \varepsilon}^{(n-1)}-G_{s, V}^{(n-1)}\right)\right\|_{X_{n}} d s \\
& +\int_{0}^{t}\left\|\left(e^{\varepsilon(t-s) L_{0}^{(n)}}-1\right) W^{(n)} G_{s, V}^{(n-1)}\right\|_{X_{n}} d s
\end{aligned}
$$

By Proposition 3.3, $L_{0}^{(n)}$ is the generator of a contraction semigroup in $X_{n}$, hence we continue

$$
\begin{align*}
\leq & \left\|G_{0, \varepsilon}^{(n)}-G_{0, V}^{(n)}\right\|_{X_{n}}+\left\|\left(e^{\varepsilon t L_{0}^{(n)}}-1\right) G_{0, V}^{(n)}\right\|_{X_{n}} \\
& +\int_{0}^{t}\left\|W^{(n)}\left(G_{s, \varepsilon}^{(n-1)}-G_{s, V}^{(n-1)}\right)\right\|_{X_{n}} d s \\
& +\int_{0}^{t}\left\|\left(e^{\varepsilon(t-s) L_{0}^{(n)}}-1\right) W^{(n)} G_{s, V}^{(n-1)}\right\|_{X_{n}} d s . \tag{4.10}
\end{align*}
$$

By (3.3), for any $n \in \mathbb{N}_{0}$

$$
\begin{equation*}
\frac{C^{n}}{n!}\left\|G_{0, \varepsilon}^{(n)}-G_{0, V}^{(n)}\right\|_{X_{n}} \leq\left\|G_{0, \varepsilon}-G_{0, V}\right\|_{\mathcal{L}_{C}} \tag{4.11}
\end{equation*}
$$

Hence, condition (4.8) implies that $\left\|G_{0, \varepsilon}^{(n)}-G_{0, V}^{(n)}\right\|_{X_{n}} \rightarrow 0$ as $\varepsilon \rightarrow 0$. Since $L_{0}^{(n)}$ is a bounded generator of a strongly continuous contraction semigroup on $X_{n}$ and since $G_{0, V}^{(n)} \in X_{n}$, we get from (3.35):

$$
\sup _{t \in[0, T]}\left\|\left(e^{\varepsilon t L_{0}^{(n)}}-1\right) G_{0, V}^{(n)}\right\|_{X_{n}} \leq \varepsilon T\left\|L_{0}^{(n)} G_{0, V}^{(n)}\right\|_{X_{n}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0
$$

Next, by the inclusion $W^{(n)} G_{s, V}^{(n-1)} \in X_{n}$, formula (3.35) also yields that, for any fixed $t>0$ and any $s \in[0, t]$,

$$
\left\|\left(e^{\varepsilon(t-s) L_{0}^{(n)}}-1\right) W^{(n)} G_{s, V}^{(n-1)}\right\|_{X_{n}} \leq \int_{0}^{\varepsilon(t-s)}\left\|L_{0}^{(n)} W^{(n)} G_{\tau, V}^{(n)}\right\|_{X_{n}} d \tau
$$

By Proposition 3.3, we continue with $A:=\frac{c_{1}+c_{2}}{2}$ and $B:=c_{3}+c_{4}$ :

$$
\leq A B n^{2}(n-1)^{2} \int_{0}^{\varepsilon(t-s)}\left\|G_{\tau, V}^{(n-1)}\right\|_{X_{n-1}} d \tau
$$

By (3.3), similarly to (4.11), we estimate:

$$
\leq A B n^{2}(n-1)^{2} \int_{0}^{\varepsilon(t-s)}(\rho(\tau, C))^{-(n-1)}(n-1)!\left\|G_{\tau, V}\right\|_{\mathcal{L}_{\rho(\tau, C)}} d \tau
$$

For $0 \leq s \leq t \leq T$ and $0<\varepsilon<1$, by (3.23) and (3.24), we continue:

$$
\leq \varepsilon T A B n^{2}(n-1)^{2}(\rho(T, C))^{-(n-1)}(n-1)!\left\|G_{0, V}\right\|_{\mathcal{L}_{C}}
$$

Therefore,

$$
\sup _{t \in[0, T]} \int_{0}^{t}\left\|\left(e^{\varepsilon(t-s) L_{0}^{(n)}}-1\right) W^{(n)} G_{s, V}^{(n-1)}\right\|_{X_{n}} d s \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0
$$

Suppose now that

$$
\begin{equation*}
\sup _{t \in[0, T]}\left\|G_{t, \varepsilon}^{(n-1)}-G_{t, V}^{(n-1)}\right\|_{X_{n-1}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.12}
\end{equation*}
$$

Then, by (3.20),

$$
\begin{aligned}
& \sup _{t \in[0, T]} \int_{0}^{t}\left\|W^{(n)}\left(G_{s, \varepsilon}^{(n-1)}-G_{s, V}^{(n-1)}\right)\right\|_{X_{n}} d s \\
\leq & B n(n-1) T \sup _{t \in[0, T]}\left\|G_{t, \varepsilon}^{(n-1)}-G_{t, V}^{(n-1)}\right\|_{X_{n-1}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 .
\end{aligned}
$$

Note that, for any $t \in[0, T]$, using (4.11), we obtain

$$
\left|G_{t, \varepsilon}^{(0)}-G_{t, V}^{(0)}\right|=\left|G_{0, \varepsilon}^{(0)}-G_{0, V}^{(0)}\right| \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0
$$

As a result, by the induction principle, we conclude from (4.10) that, for any $n \in \mathbb{N}_{0}$,

$$
\begin{equation*}
\sup _{t \in[0, T]}\left\|G_{t, \varepsilon}^{(n)}-G_{t, V}^{(n)}\right\|_{X_{n}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.13}
\end{equation*}
$$

Let now $0<r<\rho(T, C)$. Then

$$
\begin{equation*}
\sup _{t \in[0, T]}\left\|G_{t, \varepsilon}-G_{t, V}\right\|_{\mathcal{L}_{r}} \leq \sum_{n=0}^{\infty} \frac{r^{n}}{n!} \sup _{t \in[0, T]}\left\|G_{t, \varepsilon}^{(n)}-G_{t, V}^{(n)}\right\|_{X_{n}} \tag{4.14}
\end{equation*}
$$

Hence, by (4.13) and (4.14), to prove the theorem it suffices to show that the series (4.14) converges uniformly in $\varepsilon$. But the latter series may be estimated, similarly to the considerations above:

$$
\begin{aligned}
& \sum_{n=0}^{\infty} \frac{r^{n}}{n!} \sup _{t \in[0, T]}\left(\left\|G_{t, \varepsilon}^{(n)}\right\|_{X_{n}}+\left\|G_{t, V}^{(n)}\right\|_{X_{n}}\right) \\
\leq & \sum_{n=0}^{\infty} \frac{r^{n}}{n!} \sup _{t \in[0, T]}\left((\rho(t, C))^{-n} n!\left\|G_{t, \varepsilon}\right\|_{\mathcal{L}_{\rho(t, C)}}+(\rho(t, C))^{-n} n!\left\|G_{t, V}\right\|_{\mathcal{L}_{\rho(t, C)}}\right) \\
\leq & \left(\left\|G_{0, V}\right\|_{\mathcal{L}_{C}}+\sup _{\varepsilon>0}\left\|G_{0, \varepsilon}\right\|_{\mathcal{L}_{C}}\right) \sum_{n=0}^{\infty}\left(\frac{r}{\rho(T, C)}\right)^{n}<\infty
\end{aligned}
$$

since $G_{0, \varepsilon} \rightarrow G_{0, V}$ in $\mathcal{L}_{C}$ yields $\sup _{\varepsilon>0}\left\|G_{0, \varepsilon}\right\|_{\mathcal{L}_{C}}<\infty$.

Proposition 4.4. Let the conditions of Theorem 3.6 hold. Let $C_{0}>0, T=$ $\frac{1}{B C_{0}}$, and $\left\{k_{0, V}, k_{0, \varepsilon}, \varepsilon>0\right\} \subset \mathcal{K}_{C_{0}}$. Then, for any $t \in(0, T)$, there exist functions $\left\{k_{t, V}, k_{t, \varepsilon}, \varepsilon>0\right\} \subset \mathcal{K}_{C_{t}}$, with $C_{t}$ given by (3.29), such that, for any $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$,

$$
\frac{\partial}{\partial t}\left\langle\left\langle G, k_{t, \varepsilon}\right\rangle\right\rangle=\left\langle\left\langle\left(\varepsilon L_{0}+W\right) G, k_{t, \varepsilon}\right\rangle\right\rangle, \quad \frac{\partial}{\partial t}\left\langle\left\langle G, k_{t, V}\right\rangle\right\rangle=\left\langle\left\langle W G, k_{t, V}\right\rangle\right\rangle .
$$

Moreover, $\left\|k_{t, \varepsilon}\right\|_{\mathcal{K}_{C_{t}}} \leq\left\|k_{0, \varepsilon}\right\|_{\mathcal{K}_{C_{0}}},\left\|k_{t, V}\right\|_{\mathcal{K}_{C_{t}}} \leq\left\|k_{0, V}\right\|_{\mathcal{K}_{C_{0}}}$. If, additionally,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left\|k_{0, \varepsilon}-k_{0, V}\right\|_{\mathcal{K}_{C_{0}}}=0 \tag{4.15}
\end{equation*}
$$

then for any $T^{\prime} \in(0, T), r_{0}>C_{T^{\prime}}$, and $G_{0} \in \mathcal{L}_{r_{0}}$,

$$
\begin{equation*}
\sup _{t \in\left[0, T^{\prime}\right]}\left|\left\langle\left\langle G_{0}, k_{t, \varepsilon}-k_{t, V}\right\rangle\right\rangle\right| \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.16}
\end{equation*}
$$

Proof. The first part of the statement follows from Theorem 3.6. Since the function $\left[0, T^{\prime}\right] \ni t \mapsto C_{t}$ is (strictly) increasing, we have $\left\{k_{t, V}, k_{t, \varepsilon}, \varepsilon>0\right\} \subset$ $\mathcal{K}_{C_{T^{\prime}}} \subset \mathcal{K}_{r_{0}}$. Moreover, by the proof of Theorem 3.6, for any $G_{0} \in \mathcal{L}_{r_{0}} \subset \mathcal{L}_{C_{t}}$,

$$
\begin{equation*}
\left\langle\left\langle G_{0}, k_{t, \varepsilon}\right\rangle\right\rangle=\left\langle\left\langle G_{t, \varepsilon}, k_{0, \varepsilon}\right\rangle\right\rangle, \quad\left\langle\left\langle G_{0}, k_{t, V}\right\rangle\right\rangle=\left\langle\left\langle G_{t, V}, k_{0, V}\right\rangle\right\rangle \tag{4.17}
\end{equation*}
$$

where $G_{t, \varepsilon}$ and $G_{t, V}$ are solutions to (4.4) and (4.6), respectively. Therefore, by Theorem 3.5, $\left\{G_{t, \varepsilon}, G_{t, V}\right\} \subset \mathcal{L}_{\rho\left(t, r_{0}\right)}$, where the function $\rho$ is given by (3.23). Since $\rho$ is (strictly) increasing in the second variable and since $r_{0}>C_{T^{\prime}}$, we have

$$
\rho\left(T^{\prime}, r_{0}\right)>\rho\left(T^{\prime}, C_{T^{\prime}}\right)=C_{0}
$$

Applying Theorem 4.3 with $C=r_{0}, r=C_{0}, T=T^{\prime}$, and $G_{0, \varepsilon}=G_{0, V}=G_{0}$ we obtain

$$
\begin{equation*}
\sup _{t \in\left[0, T^{\prime}\right]}\left\|G_{t, \varepsilon}-G_{t, V}\right\|_{\mathcal{L}_{C_{0}}} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.18}
\end{equation*}
$$

Then, by (4.17),

$$
\begin{align*}
& \left.\sup _{t \in\left[0, T^{\prime}\right]}\left|\left\langle G_{0}, k_{t, \varepsilon}-k_{t, V}\right\rangle\right\rangle\left|=\sup _{t \in\left[0, T^{\prime}\right]}\right|\left\langle G_{t, \varepsilon}, k_{0, \varepsilon}\right\rangle\right\rangle-\left\langle\left\langle G_{t, V}, k_{0, V}\right\rangle\right\rangle \mid \\
& \left.\leq \sup _{t \in\left[0, T^{\prime}\right]}\left|\left\langle G_{t, \varepsilon}-G_{t, V}, k_{0, \varepsilon}\right\rangle\right\rangle\left|+\sup _{t \in\left[0, T^{\prime}\right]}\right|\left\langle G_{t, V}, k_{0, \varepsilon}-k_{0, V}\right\rangle\right\rangle \mid \\
& \leq \sup _{t \in\left[0, T^{\prime}\right]}\left\|G_{t, \varepsilon}-G_{t, V}\right\|_{\mathcal{L}_{C_{0}}}\left\|k_{0, \varepsilon}\right\| \|_{\mathcal{C}_{0}} \\
& \quad+\sup _{t \in\left[0, T^{\prime}\right]}\left\|G_{t, V}\right\|_{\mathcal{L}_{C_{0}}}\left\|k_{0, \varepsilon}-k_{0, V}\right\|_{\mathcal{K}_{C_{0}}} . \tag{4.19}
\end{align*}
$$

Since $k_{0, \varepsilon} \rightarrow k_{0, V}$ in $\mathcal{K}_{C_{0}}$, we get $\sup _{\varepsilon>0}\left\|k_{0, \varepsilon}\right\|_{\mathcal{K}_{C_{0}}}<\infty$. Then, by (4.18), the first summand in (4.19) converges to 0 as $\varepsilon \rightarrow 0$. Next, by (3.24),

$$
\left\|G_{t, V}\right\|_{\mathcal{L}_{C_{0}}}=\left\|G_{t, V}\right\|_{\mathcal{L}_{\rho\left(t, C_{t}\right)}} \leq\left\|G_{0}\right\|_{\mathcal{L}_{C_{t}}} \leq\left\|G_{0}\right\|_{\mathcal{L}_{C_{T^{\prime}}}} \leq\left\|G_{0}\right\|_{\mathcal{L}_{r_{0}}}
$$

Hence, by (4.15), the second summand in (4.19) also converges to 0 as $\varepsilon \rightarrow 0$, which proves the second part of the statement.

We will now show that the evolution $k_{0, V} \mapsto k_{t, V}$ satisfies our second requirement on the initial scaling $L \mapsto L_{\varepsilon}$, namely, $e_{\lambda}\left(p_{0}, \eta\right) \mapsto e_{\lambda}\left(p_{t}, \eta\right)$.
Proposition 4.5. Let the conditions of Theorem 3.6 hold. Then for any $G \in$ $B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ and any $k \in \mathcal{K}_{C}$ with $C>0$,

$$
\begin{equation*}
\int_{\Gamma_{0}}(W G)(\eta) k(\eta) d \lambda(\eta)=\int_{\Gamma_{0}} G(\eta)\left(W^{*} k\right)(\eta) d \lambda(\eta) \tag{4.20}
\end{equation*}
$$

where

$$
\begin{align*}
\left(W^{*} k\right)(\eta)= & \sum_{y_{1} \in \eta} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) k\left(\eta \cup x_{2} \cup x_{1} \backslash y_{1}\right) d x_{2} d x_{1} \\
& -\sum_{x_{1} \in \eta} \int_{\mathbb{R}^{d}} a_{1}\left(x_{1}, x_{2}\right) k\left(\eta \cup x_{2}\right) d x_{2} . \tag{4.21}
\end{align*}
$$

Here the functions $\tilde{c}$ and $a_{1}$ are defined by (3.5) and (3.15), respectively.
Proof. First, we note that, under conditions (3.11)-(3.14), for $G \in B_{\mathrm{bs}}\left(\Gamma_{0}\right)$ and $k \in \mathcal{K}_{C}$, both integrals in (4.20) are well defined. Then, using (3.8) and e.g. [7, Lemma 1], we have

$$
\begin{aligned}
& \int_{\Gamma_{0}}(W G)(\eta) k(\eta) d \lambda(\eta) \\
= & \int_{\Gamma_{0}} \int_{\mathbb{R}^{d}} \sum_{x_{1} \in \eta} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) G\left(\eta \backslash x_{1} \cup y_{1}\right) d y_{1} k\left(\eta \cup x_{2}\right) d x_{2} d \lambda(\eta) \\
& -\int_{\Gamma_{0}} \int_{\mathbb{R}^{d}} \sum_{x_{1} \in \eta} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) G(\eta) d y_{1} k\left(\eta \cup x_{2}\right) d x_{2} d \lambda(\eta) \\
= & \int_{\Gamma_{0}} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \sum_{y_{1} \in \eta} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) G(\eta) k\left(\eta \cup x_{2} \cup x_{1} \backslash y_{1}\right) d x_{2} d x_{1} d \lambda(\eta) \\
& -\int_{\Gamma_{0}} \int_{\mathbb{R}^{d}} \sum_{x_{1} \in \eta} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y_{1}\right) G(\eta) d y_{1} k\left(\eta \cup x_{2}\right) d x_{2} d \lambda(\eta),
\end{aligned}
$$

which proves the statement.
Thus, for any $p_{t} \in L^{\infty}\left(\mathbb{R}^{d}\right)$,

$$
\begin{aligned}
\left(W^{*} e_{\lambda}\left(p_{t}\right)\right)(\eta)= & \sum_{y \in \eta} e_{\lambda}\left(p_{t}, \eta \backslash y\right) \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(x_{1}, x_{2}, y\right) p_{t}\left(x_{1}\right) p_{t}\left(x_{2}\right) d x_{2} d x_{1} \\
& -\sum_{y \in \eta} e_{\lambda}\left(p_{t}, \eta \backslash y\right) p_{t}(y) \int_{\mathbb{R}^{d}} a_{1}\left(y, x_{2}\right) p_{t}\left(x_{2}\right) d x_{2}
\end{aligned}
$$

On the other hand, if $\frac{d}{d t} p_{t}$ exists, then

$$
\frac{\partial}{\partial t} e_{\lambda}\left(p_{t}, \eta\right)=\sum_{y \in \eta} e_{\lambda}\left(p_{t}, \eta \backslash y\right) \frac{\partial}{\partial t} p_{t}(y)
$$

Therefore, there exists a (point-wise) solution $k_{t}=e_{\lambda}\left(p_{t}, \eta\right)$ of the initial value problem

$$
\begin{equation*}
\frac{\partial k_{t}}{\partial t}=W^{*} k_{t},\left.\quad k_{t}\right|_{t=0}=e_{\lambda}\left(p_{0}, \eta\right) \tag{4.22}
\end{equation*}
$$

provided $p_{t}$ satisfies the non-linear Vlasov-type equation

$$
\begin{align*}
\frac{\partial}{\partial t} p_{t}(x)= & \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) p_{t}\left(y_{1}\right) p_{t}\left(y_{2}\right) d y_{1} d y_{2} \\
& -p_{t}(x) \int_{\mathbb{R}^{d}} a_{1}\left(x, x_{2}\right) p_{t}\left(x_{2}\right) d x_{2} \tag{4.23}
\end{align*}
$$

If the symmetry condition (1.4) holds, then we may rewrite (4.23) in the Boltzmanntype form

$$
\begin{align*}
\frac{\partial}{\partial t} p_{t}(x)=\int_{\mathbb{R}^{d}} & \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x, x_{2}, y_{1}, y_{2}\right) \\
& \times\left[p_{t}\left(y_{1}\right) p_{t}\left(y_{2}\right)-p_{t}(x) p_{t}\left(x_{2}\right)\right] d y_{1} d y_{2} d x_{2} \tag{4.24}
\end{align*}
$$

We are interested in positive bounded solutions of (4.23).
Proposition 4.6. Let $C>0$ and let $0 \leq p_{0} \in L^{\infty}\left(\mathbb{R}^{d}\right)$ with $\left\|p_{0}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq C$. Assume that (3.11) and (3.14) hold and, moreover,

$$
\begin{equation*}
\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(y, u_{1}, x, u_{2}\right) d u_{1} d u_{2} \leq \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x, y, u_{1}, u_{2}\right) d u_{1} d u_{2} \tag{4.25}
\end{equation*}
$$

Then, for any $T>0$, there exists a function $0 \leq p_{t} \in L^{\infty}\left(\mathbb{R}^{d}\right), t \in[0, T]$, which solves (4.23) and, moreover,

$$
\begin{equation*}
\max _{t \in[0, T]}\left\|p_{t}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq C \tag{4.26}
\end{equation*}
$$

This functions is a unique non-negative solution to (4.23) which satisfies (4.26).
Proof. Let us fix an arbitrary $T>0$ and define the Banach space $X_{T}:=$ $C\left([0, T], L^{\infty}\left(\mathbb{R}^{d}\right)\right)$ of all continuous functions on $[0, T]$ with values in $L^{\infty}\left(\mathbb{R}^{d}\right)$; the norm on $X_{T}$ is given by

$$
\|u\|_{T}:=\max _{t \in[0, T]}\left\|u_{t}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}
$$

We denote by $X_{T}^{+}$the cone of all nonnegative functions from $X_{T}$. For a given $C>0$, denote by $B_{T, C}^{+}$the set of all functions $u$ from $X_{T}^{+}$with $\|u\|_{T} \leq C$.

Let $\Phi$ be a mapping which assigns to any $v \in X_{T}$ the solution $u_{t}$ of the linear Cauchy problem

$$
\begin{align*}
\frac{\partial}{\partial t} u_{t}(x)= & -u_{t}(x) \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{t}(y) d y \\
& +\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) v_{t}\left(y_{1}\right) v_{t}\left(y_{2}\right) d y_{1} d y_{2}  \tag{4.27}\\
\left.u_{t}(x)\right|_{t=0}= & p_{0}(x)
\end{align*}
$$

namely,

$$
\begin{align*}
(\Phi v)_{t}(x)= & \exp \left\{-\int_{0}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{s}(y) d y d s\right\} p_{0}(x) \\
+ & \int_{0}^{t} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{\tau}(y) d y d \tau\right\} \\
& \times \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) v_{s}\left(y_{1}\right) v_{s}\left(y_{2}\right) d y_{1} d y_{2} d s . \tag{4.28}
\end{align*}
$$

Clearly, $v_{t} \geq 0$ implies $(\Phi v)_{t} \geq 0$. Moreover, $v \in X_{T}^{+}$yields

$$
\left|\left(\Phi v_{t}\right)(x)\right| \leq\left|p_{0}(x)\right|+c_{4} T\|v\|_{T}^{2}
$$

Therefore, $\Phi: X_{T}^{+} \rightarrow X_{T}^{+}$. Obviously, $u_{t}$ solves (4.23) if and only if $u$ is a fixed point of the map $\Phi$.

Since $0 \leq p_{0}(x) \leq C$ for a.a. $x \in \mathbb{R}^{d}$, we get from (4.28) and (4.25), for any $v \in B_{T, C}^{+}$,

$$
\begin{aligned}
0 \leq(\Phi v)_{t}(x) \leq & C \exp \left\{-\int_{0}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{s}(y) d y d s\right\} \\
& +C \int_{0}^{t} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{\tau}(y) d y d \tau\right\} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{s}(y) d y d s \\
\leq & C \exp \left\{-\int_{0}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{s}(y) d y d s\right\} \\
& +C \int_{0}^{t} \frac{\partial}{\partial s} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{\tau}(y) d y d \tau\right\} d s=C
\end{aligned}
$$

Therefore, $\Phi: B_{T, C}^{+} \rightarrow B_{T, C}^{+}$.
Let us choose any $\Upsilon \in(0, T)$ such that $2 C\left(c_{3}+c_{4}\right) \Upsilon<1$. Clearly, $v \in B_{T, C}^{+}$ implies $v \in B_{\Upsilon, C}^{+}$. By what we have shown above, $\Phi: B_{\Upsilon, C}^{+} \rightarrow B_{\Upsilon, C}^{+}$. Note the elementary inequalities $\left|e^{-a}-e^{-b}\right| \leq|a-b|$ and

$$
\left|p e^{-a}-q e^{-b}\right| \leq e^{-a}|p-q|+q e^{-b}\left|e^{-(a-b)}-1\right| \leq e^{-a}|p-q|+q e^{-b}|a-b|
$$

which hold for any $a, b, p, q \geq 0$. Hence, for any $v, w \in B_{\Upsilon, C}^{+}, t \in[0, \Upsilon]$, we get

$$
\begin{aligned}
& \left|(\Phi v)_{t}(x)-(\Phi w)_{t}(x)\right| \\
\leq & \left|\int_{0}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{s}(y) d y d s-\int_{0}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) w_{s}(y) d y d s\right| p_{0}(x) \\
& +\int_{0}^{t} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{\tau}(y) d y d \tau\right\} \\
& \times \mid \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) v_{s}\left(y_{1}\right) v_{s}\left(y_{2}\right) d y_{1} d y_{2}
\end{aligned}
$$

$$
\begin{aligned}
& \quad-\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) w_{s}\left(y_{1}\right) w_{s}\left(y_{2}\right) d y_{1} d y_{2} \mid d s \\
& +\int_{0}^{t} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) w_{\tau}(y) d y d \tau\right\} \\
& \times \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} \tilde{c}\left(y_{1}, y_{2}, x\right) w_{s}\left(y_{1}\right) w_{s}\left(y_{2}\right) d y_{1} d y_{2} \\
& \times \\
& =: \int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) v_{\tau}(y) d y d \tau-\int_{2}+I_{3} .
\end{aligned}
$$

By (3.11), (3.14), and (4.25), we have

$$
\begin{aligned}
I_{1} & \leq C c_{3}\|v-w\|_{\Upsilon \Upsilon}, \\
I_{3} & \leq C c_{3}\|v-w\|_{\Upsilon} t \int_{0}^{t} \frac{\partial}{\partial s} \exp \left\{-\int_{s}^{t} \int_{\mathbb{R}^{d}} a_{1}(x, y) w_{\tau}(y) d y d \tau\right\} d s \\
& \leq C c_{3}\|v-w\|_{\Upsilon} \Upsilon \\
I_{2} & \leq 2 C c_{4}\|v-w\|_{\Upsilon} \Upsilon .
\end{aligned}
$$

Therefore,

$$
\|\Phi v-\Phi w\|_{\Upsilon} \leq 2 C\left(c_{3}+c_{4}\right) \Upsilon\|v-w\|_{\Upsilon}
$$

for any $v, w \in B_{\Upsilon, C}^{+}$. Since $B_{\Upsilon, C}^{+}$is a metric space (with the metric induced by the norm $\left.\|\cdot\|_{\Upsilon}\right)$ and since $2 C\left(c_{3}+c_{4}\right) \Upsilon<1$, there exists a unique $u^{*} \in$ $B_{\Upsilon, C}^{+}$such that $u^{*}=\Phi\left(u^{*}\right)$. Hence, $u_{t}^{*}$ solves (4.23) for $t \in[0, \Upsilon]$. Since $0 \leq u_{\Upsilon}^{*}(x) \leq C$ for a.a. $x \in \mathbb{R}^{d}$, we can consider the equation (4.23) with the initial value $\left.p_{t}(x)\right|_{t=\Upsilon}=u_{\Upsilon}^{*}(x)$. Then we obtain a unique non-negative solution which satisfies $\max _{t \in[\Upsilon, 2 \Upsilon]}\left\|u_{t}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq C$, and so on. As a result, we obtain a solution of (4.23) on $[0, T]$. The uniqueness among all solutions from $B_{T, C}^{+}$is now obvious.

Remark 4.7. Note that, in the proof of Proposition 4.6, we did not use the property (1.3). On the other hand, all considerations remain true if, instead of (4.25), we assume that

$$
\begin{equation*}
\int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(u_{1}, y, x, u_{2}\right) d u_{1} d u_{2} \leq \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c\left(x, y, u_{1}, u_{2}\right) d u_{1} d u_{2} \tag{4.29}
\end{equation*}
$$

Suppose we have an expansion

$$
\begin{equation*}
c\left(\left\{x_{1}, x_{2}\right\},\left\{y_{1}, y_{2}\right\}\right)=c^{\prime}\left(x_{1}, x_{2}, y_{1}, y_{2}\right)+c^{\prime \prime}\left(x_{1}, x_{2}, y_{1}, y_{2}\right) \tag{4.30}
\end{equation*}
$$

where $c^{\prime}, c^{\prime \prime}$ are functions which satisfy conditions (4.25) and (4.29), respectively, as well as conditions (3.11) and (3.14) (but do not necessarily satisfy (1.3)). It is easy to see that all considerations in the proof of Proposition 4.6 remain true for this $c$. Let us give a quite natural example of functions $c^{\prime}, c^{\prime \prime}$.

Example 4.8 (cf. [10]). Let $c$ be given by (4.30) with

$$
c^{\prime}\left(x_{1}, x_{2}, y_{1}, y_{2}\right)=\varkappa a\left(x_{1}-y_{1}\right) a\left(x_{2}-y_{2}\right)\left[b\left(x_{1}-x_{2}\right)+b\left(y_{1}-y_{2}\right)\right] .
$$

and

$$
c^{\prime \prime}\left(x_{1}, x_{2}, y_{1}, y_{2}\right)=c^{\prime}\left(x_{2}, x_{1}, y_{1}, y_{2}\right)
$$

Here $\varkappa>0,0 \leq a, b \in L^{1}\left(\mathbb{R}^{d}\right) \cap L^{\infty}\left(\mathbb{R}^{d}\right),\|a\|_{L^{1}\left(\mathbb{R}^{d}\right)}=\|b\|_{L^{1}\left(\mathbb{R}^{d}\right)}=1$ and $b$ is an even function. Then the condition (4.29) for $c^{\prime \prime}$ coincides with the condition (4.25) for $c^{\prime}$. Note also that (3.11)-(3.14) are satisfied for $c$. For example, $c_{4}=4 \varkappa$ and

$$
c_{1} \leq 2 \varkappa\|b\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}+2 \varkappa\|a\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}\|b\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}<\infty .
$$

Next, let us check whether (4.25) holds for $c^{\prime}$. We have

$$
\begin{aligned}
& \varkappa \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c^{\prime}\left(y, u_{1}, x, u_{2}\right) d u_{1} d u_{2} \\
= & \varkappa \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} a(y-x) a\left(u_{1}-u_{2}\right)\left(b\left(y-u_{1}\right)+b\left(x-u_{2}\right)\right) d u_{1} d u_{2}=2 \varkappa a(y-x)
\end{aligned}
$$

and

$$
\begin{aligned}
& \varkappa \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} c^{\prime}\left(x, y, u_{1}, u_{2}\right) d u_{1} d u_{2} \\
= & \varkappa \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} a\left(x-u_{1}\right) a\left(y-u_{2}\right)\left(b(x-y)+b\left(u_{1}-u_{2}\right)\right) d u_{1} d u_{2} \\
= & \varkappa b(x-y)+\varkappa \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} a\left(x-u_{1}\right) a\left(y-u_{2}\right) b\left(u_{1}-u_{2}\right) d u_{1} d u_{2}
\end{aligned}
$$

Since $b$ is even, (4.25) holds if, for example, $2 a(x) \leq b(x), x \in \mathbb{R}^{d}$.

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