# Symmetric differential operators of the second order in Poisson spaces 

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#### Abstract

We study symmetric differential operators of the second order in Poisson spaces.


## 1 Introduction

Recently [4] we have proved an analogue of the classical Gauss formula for the configuration space $\Gamma_{\Lambda}$ over a domain $\Lambda$ of $\mathbb{R}^{d}$ and have studied symmetric extensions of the corresponding Laplacian on $\Gamma_{\Lambda}$. (We refer to [1], [2], 6], [8] for the basis notions of analysis and geometry on configuration spaces.) The purpose of this paper is to extend some of the results of [4] to the case of general second order differential operators.

The article is arranged as follows. In Section 2 we present some well-known facts on second order differential operators in $L^{2}\left(\mathbb{R}^{d}\right)$. In Section 3 we construct the second quantization of a closable operator and define symmetric second order differential operators $H_{A_{\sigma}}^{P}$ and $H_{A_{\sigma}}^{P, \Lambda}$ in $L^{2}\left(\Gamma, \pi_{\sigma}\right)$ and $L^{2}\left(\Gamma_{\Lambda}, \pi_{\sigma}^{\Lambda}\right)$ respectively. Here $\pi_{\sigma}$ is the Poisson measure with intensity $\sigma$ and $\pi_{\sigma}^{\Lambda}$ is the restriction of $\pi_{\sigma}$ to $\Gamma_{\Lambda}$. In Section 4 we prove an analogue of the classical Green formula for $H_{A_{\sigma}}^{P}$. It gives the possibility to describe symmetric extensions of the corresponding minimal operator in $L^{2}\left(\Gamma_{\Lambda}, \pi_{\sigma}^{\Lambda}\right)$ (see Section 5).

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## 2 Differential operators of the second order in $L^{2}\left(\mathbb{R}^{d}, \sigma\right)$.

Let $A$ be the second order symmetric differential expression of the form

$$
\begin{equation*}
(A f)(x):=\sum_{j, k=1}^{d} \frac{\partial}{\partial x_{j}}\left(a_{j k}(x) \frac{\partial f}{\partial x_{k}}(x)\right)+c(x) f(x), \tag{2.1}
\end{equation*}
$$

where $a_{j k} \in C^{2}\left(\mathbb{R}^{d}\right), c \in C\left(\mathbb{R}^{d}\right)$ and $a_{j k}=a_{k j}$ (all the functions $a_{j k}, c$ are assumed to be real).

Let $\Lambda$ be an open domain of $\mathbb{R}^{d}$. In the following we always suppose that the boundary $\partial \Lambda$ of $\Lambda$ is piecewise $C^{1}$. Let $\tilde{m}$ be the surface measure on $\partial \Lambda$ corresponding to Lebesgue measure $m$ on $\mathbb{R}^{d}$. By $n_{s}$ we denote the outer normal to $\partial \Lambda$ (at the point $s \in \partial \Lambda$ ) and for any $f \in C^{1}(\bar{\Lambda})$ we consider the co-normal derivative

$$
\begin{equation*}
\frac{\partial f}{\partial n^{a}}(s):=\sum_{j, k=1}^{d} a_{j k}(s) \frac{\partial f}{\partial s_{j}}(s) \cos \left(n_{s}, s_{k}\right), \quad s \in \partial \Lambda . \tag{2.2}
\end{equation*}
$$

Let us introduce also the bilinear form

$$
\begin{equation*}
a(u, v)(x):=\sum_{j, k=1}^{d} a_{j k}(x) \frac{\partial u}{\partial x_{j}}(x) \frac{\partial v}{\partial x_{k}}(x), \quad u, v \in C^{2}\left(\mathbb{R}^{d}\right) . \tag{2.3}
\end{equation*}
$$

Consider a $C^{1}$-density $\rho>0$ on $\mathbb{R}^{d}$. Set $\sigma(d x)=\rho(x) m(d x)$, then $\sigma$ is a nonatomic Radon measure on $\mathbb{R}^{d}$. Let $\tilde{\sigma}$ be the surface measure on $\partial \Lambda$ corresponding to $\sigma$

$$
\tilde{\sigma}(d s)=\rho(s) \tilde{m}(d s)
$$

It's easy to see, that for any $f, g \in C_{0}^{2}\left(\mathbb{R}^{d}\right)$

$$
\begin{aligned}
\int_{\mathbb{R}^{d}}((A f)(x) g(x) & -f(x)(A g)(x)) \sigma(d x) \\
& =\int_{\mathbb{R}^{d}}(f(x) a(g, \ln \rho)(x)-g(x) a(f, \ln \rho)(x)) \sigma(d x) .
\end{aligned}
$$

Then let us construct a new differential expression $A_{\sigma}$ which acts on the smooth function $f$ on $\mathbb{R}^{d}$ as follows

$$
\begin{equation*}
\left(A_{\sigma} f\right)(x):=(A f)(x)+a(f, \ln \rho)(x) . \tag{2.4}
\end{equation*}
$$

It defines the second order symmetric differential operator:

$$
C_{0}^{2}\left(\mathbb{R}^{d}\right) \ni f \mapsto\left(A_{\sigma} f\right) \in L^{2}(\sigma),
$$

where $L^{p}(\sigma):=L^{p}\left(\mathbb{R}^{d}, \sigma\right)$.
In the following we always suppose that

$$
\begin{equation*}
c \in L^{2}(\sigma) \cap L^{1}(\sigma) . \tag{2.5}
\end{equation*}
$$

Definition 2.1. Introduce a class of smooth functions, which are equal to constant outside some compact set:

$$
\begin{equation*}
\mathcal{K}_{2}:=\text { linear hull of }\left\{1, C_{0}^{2}\left(\mathbb{R}^{d}\right)\right\} . \tag{2.6}
\end{equation*}
$$

The following proposition is a variant of well-known integration by part formula.

Proposition 2.2. For any $f, g \in \mathcal{K}_{2}$ the following statements are true:

1. $\left\{A_{\sigma} f, A_{\sigma} f \cdot g\right\} \in L^{2}(\sigma) \cap L^{1}(\sigma)$.
2. $\left(A_{\sigma} f, g\right)_{L^{2}(\sigma)}=\left(f, A_{\sigma} g\right)_{L^{2}(\sigma)}$.
3. (The first Green formula)

$$
\begin{align*}
\int_{\Lambda}\left(A_{\sigma} f\right)(x) g(x) \sigma(d x) & =\int_{\Lambda} c(x) f(x) g(x) \sigma(d x) \\
& +\int_{\Lambda} a(f, g)(x) \sigma(d x)+\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} f(s) g(s) \tilde{\sigma}(d s) \tag{2.7}
\end{align*}
$$

4. (The second Green formula)

$$
\begin{align*}
\int_{\Lambda}\left(\left(A_{\sigma} f\right)(x)\right. & \left.g(x)-f(x)\left(A_{\sigma} g\right)(x)\right) \sigma(d x) \\
= & \int_{\partial \Lambda}\left(\frac{\partial}{\partial n^{a}} f(s) \cdot g(s)-f(s) \cdot \frac{\partial}{\partial n^{a}} g(s)\right) \tilde{\sigma}(d s) \tag{2.8}
\end{align*}
$$

## 3 The image of the second quantization operator in Poisson spaces

Let $(B, D(B))$ be a closable operator in a real separable Hilbert space $\mathcal{H}$ with a dense domain $D(B)$. Define the Fock space as follows

$$
\begin{equation*}
\operatorname{Exp} \mathcal{H}:=\bigoplus_{n=0}^{\infty} \operatorname{Exp}_{n} \mathcal{H}:=\bigoplus_{n=0}^{\infty} \mathcal{H}^{\hat{\otimes} n} \tag{3.1}
\end{equation*}
$$

where $\operatorname{Exp}_{0} \mathcal{H}:=\mathbb{R}$. Consider an operator $\left(B^{(n)},(D(B))^{\hat{\otimes} n}\right)$ in $\operatorname{Exp}_{n} \mathcal{H}, n \in \mathbb{N}$, such as

$$
\begin{equation*}
B^{(n)}:=B \otimes 1 \ldots \otimes 1+1 \otimes B \otimes 1 \ldots \otimes 1+\ldots+1 \otimes \ldots \otimes 1 \otimes B \tag{3.2}
\end{equation*}
$$

and set $B^{(0)}:=0$ in $\operatorname{Exp}_{0} \mathcal{H}$. Let

$$
\begin{equation*}
\operatorname{Exp}_{f i n}(D(B)):=\left\{\left(f_{0}, \ldots, f_{k}, 0, \ldots\right) \mid f_{j} \in(D(B))^{\hat{\otimes} j}, j=0, \ldots, k, k \in \mathbb{N}_{0}\right\} \tag{3.3}
\end{equation*}
$$

where $\mathbb{N}_{0}:=\mathbb{N} \cup\{0\}$. Define the second quantization $d \operatorname{Exp} B$ of the operator $B$ as an operator on the dense domain $\operatorname{Exp}_{\text {fin }}(D(B))$ in $\operatorname{Exp} \mathcal{H}$ such that

$$
\begin{equation*}
d \operatorname{Exp} B \upharpoonright \operatorname{Exp}_{n} \mathcal{H}=B^{(n)}, n \in \mathbb{N}_{0} \tag{3.4}
\end{equation*}
$$

Proposition 3.1. The operator $\left(d \operatorname{Exp} B, \operatorname{Exp}_{\text {fin }}(D(B))\right)$ is closable in $\operatorname{Exp} \mathcal{H}$.
Proof. It's easy to see that for any $n \in \mathbb{N}_{0} D\left(\left(B^{(n)}\right)^{*}\right)$ is a dense domain in $\operatorname{Exp}_{n} \mathcal{H}$.

In the following we retain the notation $d \operatorname{Exp} B$ for the closure of this operator.

Set $\mathcal{H}:=L^{2}(\sigma)$ and suppose that $\mathcal{D}:=C_{0}^{\infty}\left(\mathbb{R}^{d}\right) \subset D(B) \subset C^{\infty}\left(\mathbb{R}^{d}\right)$. There is a canonical isomorphism between $\operatorname{Exp} L^{2}(\sigma)$ and $L^{2}\left(\Gamma, \pi_{\sigma}\right)=: L^{2}\left(\pi_{\sigma}\right)$ (where $\Gamma:=\Gamma_{\mathbb{R}^{d}}, \pi_{\sigma}$ is a Poisson measure with intensity measure $\sigma$ ) such that

$$
\begin{equation*}
\operatorname{Exp}_{n} L^{2}(\sigma) \ni \varphi^{\otimes n} \mapsto Q_{n}\left(\varphi^{\otimes n}, \cdot\right) \in L^{2}\left(\pi_{\sigma}\right), \varphi \in \mathcal{D} \subset D(B), \tag{3.5}
\end{equation*}
$$

where $Q_{n}\left(\varphi^{\otimes n}, \cdot\right)$ is a Charlier polynomial on $\Gamma$ (see [1], 4] for details). Let $H_{B}^{P}$ be the image of the operator $d \operatorname{Exp} B$ under this isomorphism. The proof of the following proposition is analogous to that of Theorem 5.1 of [1].

Proposition 3.2. For all $F, G \in \mathcal{F P}(\mathcal{D}, \Gamma)$ the following formula holds

$$
\begin{align*}
\int_{\Gamma}\left(H_{B}^{P} F\right)(\gamma) \cdot G(\gamma) & \pi_{\sigma}(d \gamma) \\
= & \int_{\Gamma} \int_{\mathbb{R}^{d}} B \nabla^{P} F(\gamma, x) \cdot \nabla^{P} G(\gamma, x) \sigma(d x) \pi_{\sigma}(d \gamma) \tag{3.6}
\end{align*}
$$

Note that $\mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma) \subset D\left(H_{B}^{P}\right)$ and by standard approximation arguments one can extend (3.6) on $F, G \in \mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma)$ (see, e.g., [1]).

Set now $\mathcal{H}:=L^{2}(\Lambda, \sigma)$ and suppose that $\mathcal{D}(\Lambda) \subset D(B) \subset \mathcal{D}$, where $\mathcal{D}(\Lambda)=$ $C_{0}^{\infty}(\Lambda)$. We consider the canonical isomorphism between $\operatorname{Exp} L^{2}(\Lambda, \sigma)$ and $L^{2}\left(\Gamma_{\Lambda}, \pi_{\sigma}^{\Lambda}\right)$ such that

$$
\begin{equation*}
\operatorname{Exp}_{n} L^{2}(\Lambda, \sigma) \ni \varphi^{\otimes n} \mapsto Q_{n, \Lambda}\left(\varphi^{\otimes n}, \cdot\right) \in L^{2}\left(\pi_{\sigma}^{\Lambda}\right), \tag{3.7}
\end{equation*}
$$

where $Q_{\Lambda, n}\left(\varphi^{\otimes n}, \cdot\right)$ is a Charlier polynomial on $\Gamma_{\Lambda}$. Denote by $H_{B}^{P, \Lambda}$ the image of $d \operatorname{Exp} B$ under this isomorphism. Clearly (cf. (3.6)) for any $F, G \in$ $\mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$.

$$
\begin{align*}
& \int_{\Gamma_{\Lambda}}\left(H_{B}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
&=\int_{\Gamma_{\Lambda}} \int_{\Lambda} B \nabla^{P} F(\gamma, x) \cdot \nabla^{P} G(\gamma, x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \tag{3.8}
\end{align*}
$$

Suppose that the operator $B$ is associated with the differential expression $A_{\sigma}$ :

$$
B f=A_{\sigma} f, f \in D(B)
$$

In this situation we'll write $H_{A_{\sigma}}^{P}$ and $H_{A_{\sigma}}^{P, \Lambda}$ instead of $H_{B}^{P}$ and $H_{B}^{P, \Lambda}$, respectively.
Let $F \in \mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma)$, i.e.,

$$
\begin{gather*}
F(\gamma):=g_{F}\left(\left\langle\varphi_{1}, \gamma\right\rangle, \ldots,\left\langle\varphi_{N}, \gamma\right\rangle\right),  \tag{3.9}\\
\varphi_{j} \in \mathcal{D}, j=1, \ldots, N, g_{F} \in C_{p}^{\infty}\left(\mathbb{R}^{N}\right),
\end{gather*}
$$

where $C_{p}^{\infty}\left(\mathbb{R}^{N}\right)$ is the set of all $C^{\infty}$-functions $g$ on $\mathbb{R}^{N}$ such that $g$ and all its partial derivatives are polynomially bounded. Set

$$
\begin{gather*}
\hat{F}_{r}(\gamma):=\frac{\partial g_{F}}{\partial q_{r}}\left(\left\langle\varphi_{1}, \gamma\right\rangle, \ldots,\left\langle\varphi_{N}, \gamma\right\rangle\right), 1 \leq r \leq N,  \tag{3.10}\\
\hat{F}_{r, s}(\gamma):=\frac{\partial^{2} g_{F}}{\partial q_{r} \partial q_{s}}\left(\left\langle\varphi_{1}, \gamma\right\rangle, \ldots,\left\langle\varphi_{N}, \gamma\right\rangle\right), 1 \leq r, s \leq N . \tag{3.11}
\end{gather*}
$$

The direct calculation gives the next result:
Proposition 3.3. For any $F \in \mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma)$ the following formula holds

$$
\begin{aligned}
A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) & =\sum_{r, s=1}^{N} \hat{F}_{r, s}\left(\gamma+\varepsilon_{x}\right) a\left(\varphi_{r}, \varphi_{s}\right)(x) \\
& +\sum_{r=1}^{N} \hat{F}_{r}\left(\gamma+\varepsilon_{x}\right)\left(\left(A_{\sigma} \varphi_{r}\right)(x)+a\left(\varphi_{r}, \ln \rho\right)(x)-c(x) \varphi_{r}(x)\right) \\
& +F\left(\gamma+\varepsilon_{x}\right) c(x)
\end{aligned}
$$

Remark 3.4. 1. By (3.9) for any function $F \in \mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma)$ and for any $\gamma \in \Gamma$ we have

$$
\begin{equation*}
F(\gamma+\varepsilon .) \in \mathcal{K}_{2} . \tag{3.12}
\end{equation*}
$$

2. Note that when we write, for example, $A_{\sigma} F(\gamma)$ we mean that the expression $A_{\sigma}$ acts on constant $F(\gamma)$ as on the constant function of $x \in \mathbb{R}^{d}$. But we can also consider the action of the expression $A_{\sigma}$ on function $F(\gamma)$ as function of some $x \in \gamma$ when others point of $\gamma$ are fixed. In this situation we will write $\left(A_{\sigma}\right)_{x} F(\gamma), x \in \gamma$. It's easy to see that for any $x \in \gamma$

$$
\begin{gather*}
\left(A_{\sigma}\right)_{x} F(\gamma)=\sum_{r, s=1}^{N} \hat{F}_{r, s}(\gamma) a\left(\varphi_{r}, \varphi_{s}\right)(x)  \tag{3.13}\\
+\sum_{r=1}^{N} \hat{F}_{r}(\gamma)\left(\left(A_{\sigma} \varphi_{r}\right)(x)+a\left(\varphi_{r}, \ln \rho\right)(x)-c(x) \varphi_{r}(x)\right)+F(\gamma) c(x)
\end{gather*}
$$

Proposition 3.5. For any $F \in \mathcal{F} C_{p}^{\infty}(\mathcal{D}, \Gamma)$ the following formula holds

$$
\begin{align*}
\left(H_{A_{\sigma}}^{P} F\right)(\gamma) & =\left\langle\left(A_{\sigma}\right) \cdot F(\gamma), \gamma\right\rangle-\langle F(\gamma-\varepsilon .) c(\cdot), \gamma\rangle \\
& -\int_{\mathbb{R}^{d}} F\left(\gamma+\varepsilon_{x}\right) c(x) \sigma(d x)+F(\gamma) \int_{\mathbb{R}^{d}} c(x) \sigma(d x) \tag{3.14}
\end{align*}
$$

Proof. By (3.6), Proposition 2.2 and the Mecke identity we have

$$
\begin{aligned}
& \int_{\Gamma}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}(d \gamma) \\
& =\int_{\Gamma} \int_{\mathbb{R}^{d}} A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) \cdot G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& -\int_{\Gamma} \int_{\mathbb{R}^{d}} A_{\sigma} F(\gamma) \cdot G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& -\int_{\Gamma} \int_{\mathbb{R}^{d}} A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) \cdot G(\gamma) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& +\int_{\Gamma} \int_{\mathbb{R}^{d}} A_{\sigma} F(\gamma) \cdot G(\gamma) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& =\int_{\Gamma} \int_{\mathbb{R}^{d}}\left(\sum_{r, s=1}^{N} \hat{F}_{r, s}\left(\gamma+\varepsilon_{x}\right) a\left(\varphi_{r}, \varphi_{s}\right)(x)\right. \\
& +\sum_{r=1}^{N} \hat{F}_{r}\left(\gamma+\varepsilon_{x}\right)\left(\left(A_{\sigma} \varphi_{r}\right)(x)+a\left(\varphi_{r}, \ln \rho\right)(x)-c(x) \varphi_{r}(x)\right) \\
& \left.+F\left(\gamma+\varepsilon_{x}\right) c(x)\right) G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& -\int_{\Gamma} \int_{\mathbb{R}^{d}} F(\gamma) c(x) G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& -\int_{\Gamma} \int_{\mathbb{R}^{d}} F\left(\gamma+\varepsilon_{x}\right) c(x) G(\gamma) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& +\int_{\Gamma} \int_{\mathbb{R}^{d}} F(\gamma) c(x) G(\gamma) \sigma(d x) \pi_{\sigma}(d \gamma) \\
& =\int_{\Gamma}\left(\sum_{r, s=1}^{N} \hat{F}_{r, s}(\gamma)\left\langle a\left(\varphi_{r}, \varphi_{s}\right), \gamma\right\rangle\right. \\
& +\sum_{r=1}^{N} \hat{F}_{r}(\gamma)\left\langle A_{\sigma} \varphi_{r}+a\left(\varphi_{r}, \ln \rho\right)-\varphi_{r} c, \gamma\right\rangle \\
& +F(\gamma)\langle c, \gamma\rangle) G(\gamma) \pi_{\sigma}(d \gamma) \\
& -\int_{\Gamma}\langle F(\gamma-\varepsilon .) c(\cdot), \gamma\rangle G(\gamma) \pi_{\sigma}(d \gamma)
\end{aligned}
$$

$$
\begin{aligned}
& -\int_{\Gamma}\left(\int_{\mathbb{R}^{d}} F\left(\gamma+\varepsilon_{x}\right) c(x) \sigma(d x)\right) G(\gamma) \pi_{\sigma}(d \gamma) \\
& +\int_{\Gamma}\left(F(\gamma) \int_{\mathbb{R}^{d}} c(x) \sigma(d x)\right) G(\gamma) \pi_{\sigma}(d \gamma)
\end{aligned}
$$

So, by (3.13) the proof is fulfilled.
Proposition 3.6. For any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$ the following formula holds

$$
\begin{align*}
\left(H_{A_{\sigma}}^{P, \Lambda} F\right)(\gamma) & =\left\langle\left(A_{\sigma}\right) \cdot F(\gamma), \gamma\right\rangle-\langle F(\gamma-\varepsilon \cdot) c(\cdot), \gamma\rangle \\
& -\int_{\Lambda} F\left(\gamma+\varepsilon_{x}\right) c(x) \sigma(d x)+F(\gamma) \int_{\Lambda} c(x) \sigma(d x)  \tag{3.15}\\
& -\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s)
\end{align*}
$$

Proof. The proof is analogous to that of (3.14).

## 4 Green Formulas

We start with the so-called second Green formula:
Proposition 4.1. For any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$ the following formula holds

$$
\begin{gather*}
\int_{\Gamma_{\Lambda}}\left(\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot G(\gamma)-F(\gamma) \cdot\left(H_{A_{\sigma}}^{P} G\right)(\gamma)\right) \pi_{\sigma}^{\Lambda}(d \gamma)  \tag{4.1}\\
=\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda}\left(\frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) \cdot G\left(\gamma+\varepsilon_{s}\right)-F\left(\gamma+\varepsilon_{s}\right) \cdot \frac{\partial}{\partial n^{a}} G\left(\gamma+\varepsilon_{s}\right)\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma) \\
-\int_{\Gamma_{\Lambda}} \int_{\mathbb{R}^{d} \backslash \Lambda}\left(F\left(\gamma+\varepsilon_{x}\right) G(\gamma)-G\left(\gamma+\varepsilon_{x}\right) F(\gamma)\right) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{gather*}
$$

Proof. By (3.14), (3.12) and the Mecke identity we have

$$
\begin{aligned}
\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot & G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\Gamma_{\Lambda}}\left\langle\left(A_{\sigma}\right) . F(\gamma), \gamma\right\rangle \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\langle F(\gamma-\varepsilon .) c(\cdot), \gamma\rangle \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\left(\int_{\mathbb{R}^{d}} F\left(\gamma+\varepsilon_{x}\right) c(x) \sigma(d x)\right) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}}\left(F(\gamma) \int_{\mathbb{R}^{d}} c(x) \sigma(d x)\right) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\Gamma_{\Lambda}} \int_{\Lambda} A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) \cdot G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{aligned}
$$

$$
\begin{aligned}
& -\int_{\Gamma_{\Lambda}} \int_{\Lambda} F(\gamma) c(x) G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}} \int_{\mathbb{R}^{d}} F\left(\gamma+\varepsilon_{x}\right) G(\gamma) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}} \int_{\mathbb{R}^{d}} F(\gamma) G(\gamma) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{aligned}
$$

and from the corresponding formula for $\int_{\Gamma_{\Lambda}} F(\gamma) \cdot\left(H_{A_{\sigma}}^{P} G\right)(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)$ one has

$$
\begin{gathered}
\int_{\Gamma_{\Lambda}}\left(\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot G(\gamma)-F(\gamma) \cdot\left(H_{A_{\sigma}}^{P} G\right)(\gamma)\right) \pi_{\sigma}^{\Lambda}(d \gamma) \\
=\int_{\Gamma_{\Lambda}} \int_{\Lambda}\left(A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) \cdot G\left(\gamma+\varepsilon_{x}\right)-F\left(\gamma+\varepsilon_{x}\right) \cdot A_{\sigma} G\left(\gamma+\varepsilon_{x}\right)\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
-\int_{\Gamma_{\Lambda}} \int_{\Lambda}\left(F(\gamma) G\left(\gamma+\varepsilon_{x}\right)-G(\gamma) F\left(\gamma+\varepsilon_{x}\right)\right) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
-\int_{\Gamma_{\Lambda}} \int_{\mathbb{R}^{d}}\left(F\left(\gamma+\varepsilon_{x}\right) G(\gamma)-G\left(\gamma+\varepsilon_{x}\right) F(\gamma)\right) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{gathered}
$$

So, (3.12) and (2.8) imply the statement of proposition.
Let us consider for any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)(F$ has the form (3.9), $G$ has the analogous form $\left.G(\gamma):=g_{G}\left(\left\langle\psi_{1}, \gamma\right\rangle, \ldots,\left\langle\psi_{M}, \gamma\right\rangle\right)\right)$ the bilinear form

$$
\begin{equation*}
a^{\Gamma}(F, G)(\gamma):=\sum_{j=1}^{N} \sum_{k=1}^{M} \hat{F}_{j}(\gamma) \hat{G}_{k}(\gamma)\left\langle a\left(\varphi_{j}, \psi_{k}\right), \gamma\right\rangle \tag{4.2}
\end{equation*}
$$

and the corresponding (pre-)Dirichlet form:

$$
\begin{equation*}
\mathcal{E}_{a, \pi_{\sigma}^{\Lambda}}^{\Gamma}(F, G):=\int_{\Gamma_{\Lambda}} a^{\Gamma}(F, G)(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \tag{4.3}
\end{equation*}
$$

In the following we always suppose, that

$$
\begin{equation*}
c \equiv 0 \text { on } \mathbb{R}^{d} \backslash \Lambda \tag{4.4}
\end{equation*}
$$

Using this condition we can prove the so-called first Green formula.
Proposition 4.2. For any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$

$$
\begin{align*}
\mathcal{E}_{a, \pi_{\sigma}^{\Lambda}}^{\Gamma}(F, G) & =\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\left(H_{c}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)  \tag{4.5}\\
& -\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) G\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{align*}
$$

where we understand $c$ as the operator of the multiplication on the function $c(\cdot)$ in $L^{2}(\Lambda, \sigma)$.

Proof. We have for any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$

$$
\begin{aligned}
\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) & \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\Gamma_{\Lambda}} \int_{\Lambda} A_{\sigma} F\left(\gamma+\varepsilon_{x}\right) \cdot G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}} \int_{\Lambda} F(\gamma) c(x) G\left(\gamma+\varepsilon_{x}\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}} \int_{\Lambda} F\left(\gamma+\varepsilon_{x}\right) G(\gamma) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}} \int_{\Lambda} F(\gamma) G(\gamma) c(x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\Gamma_{\Lambda}} \int_{\Lambda} a\left(F\left(\gamma+\varepsilon_{x}\right), G\left(\gamma+\varepsilon_{x}\right)\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}} \int_{\Lambda} c(x) \nabla^{P} F(\gamma, x) \nabla^{P} G(\gamma, x) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) G\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{aligned}
$$

Next, by Mecke identity one has

$$
\int_{\Gamma_{\Lambda}} \int_{\Lambda} a\left(F\left(\gamma+\varepsilon_{x}\right), G\left(\gamma+\varepsilon_{x}\right)\right) \sigma(d x) \pi_{\sigma}^{\Lambda}(d \gamma)=\int_{\Gamma_{\Lambda}} a^{\Gamma}(F, G)(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

so, by (3.8) we have

$$
\begin{align*}
\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) & \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\Gamma_{\Lambda}} a^{\Gamma}(F, G)(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& +\int_{\Gamma_{\Lambda}}\left(H_{c}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)  \tag{4.6}\\
& +\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) G\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{align*}
$$

## 5 The symmetric extensions of the minimal operator

Let us consider the minimal operator:

$$
\begin{equation*}
H_{\min }:=\left(H_{A_{\sigma}}^{P}, \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}(\Lambda), \Gamma_{\Lambda}\right)\right) . \tag{5.1}
\end{equation*}
$$

By the first Green formula the operator $H_{\min }$ is symmetric in $L^{2}\left(\Gamma_{\Lambda}\right)$. We define the maximal operator by the standard way:

$$
\begin{equation*}
H_{\max }:=\left(H_{\min }\right)^{*} \tag{5.2}
\end{equation*}
$$

We can formulate the following proposition

## Proposition 5.1.

$$
\mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right) \subset D\left(H_{\max }\right)
$$

and for any $G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$

$$
H_{\max } G=H_{A_{\sigma}}^{P, \Lambda} G
$$

Proof. For any $F \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}(\Lambda), \Gamma_{\Lambda}\right), G \in F C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$ we have

$$
\begin{gathered}
\int_{\Gamma_{\Lambda}}\left(\left(H_{\min } F\right)(\gamma) \cdot G(\gamma)-F(\gamma) \cdot\left(H_{A_{\sigma}}^{P} G\right)(\gamma)\right) \pi_{\sigma}^{\Lambda}(d \gamma) \\
=\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda}\left(\frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) \cdot G\left(\gamma+\varepsilon_{s}\right)-F\left(\gamma+\varepsilon_{s}\right) \cdot \frac{\partial}{\partial n^{a}} G\left(\gamma+\varepsilon_{s}\right)\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma) \\
=-\int_{\Gamma_{\Lambda}} F(\gamma) \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} G\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{gathered}
$$

Hence $G \in D\left(H_{\max }\right)$ and by (4.4)

$$
H_{\max } G=\left(H_{A_{\sigma}}^{P} G\right)(\gamma)-\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} G\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s)=H_{A_{\sigma}}^{P, \Lambda} G(\gamma)
$$

Therefore we have the following inclusion

$$
\begin{equation*}
H_{\min } \subset H_{A_{\sigma}}^{P, \Lambda} \subset H_{\max } \tag{5.3}
\end{equation*}
$$

Suppose that $\mathcal{A}$ be an algebra, such that $\mathcal{D}(\Lambda) \subset \mathcal{A} \subset \mathcal{D}$.
Remark 5.2. If the operator $\left(A_{\sigma}, \mathcal{A}\right)$ is symmetric in $L^{2}(\Lambda)$, then $\left(H_{A_{\sigma}}^{P, \Lambda}, \mathcal{F} \mathcal{P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right)$ is the symmetric extension of $H_{\min }$ in $L^{2}\left(\Gamma_{\Lambda}\right)$.

To describe symmetric extensions of $H_{\min }$, which are defined by the differential expression $H_{A_{\sigma}}^{P}$, we start with the following simple proposition.
Proposition 5.3. Suppose that

$$
\mathcal{F} C_{p}^{\infty}\left(\mathcal{D}(\Lambda), \Gamma_{\Lambda}\right) \subset \mathcal{F} \subset \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)
$$

and $\left(H_{A_{\sigma}}^{P}, \mathcal{F}\right)$ be a symmetric extensions of $H_{\min }$. Then for any $F \in \mathcal{F}$

$$
\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s)=0 .\left(\bmod \pi_{\sigma}^{\Lambda}\right)
$$

Proof. It's easy to see from (5.3) and the fact, that the symmetric extension of $H_{\text {min }}$ is the symmetric restriction of $H_{\text {max }}$.

Now we can formulate our main result.
Theorem 5.4. The operator $\left(H_{A_{\sigma}}^{P}, \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right)$ is symmetric in $L^{2}\left(\Gamma_{\Lambda}\right)$ if and only if the operator $\left(A_{\sigma}, \mathcal{A}\right)$ is symmetric in $L^{2}(\Lambda)$ and for any $\varphi \in \mathcal{A}$

$$
\begin{equation*}
\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \tilde{\sigma}(d s)=0 \tag{5.4}
\end{equation*}
$$

Moreover, in this case

$$
\left(H_{A_{\sigma}}^{P}, \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right)=\left(H_{A_{\sigma}}^{P, \Lambda}, \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right) .
$$

Proof. Suppose that the operator $\left(A_{\sigma}, \mathcal{A}\right)$ is symmetric in $L^{2}(\Lambda)$ and for any $\varphi \in \mathcal{A}$ the condition (5.4) is valid. For $\varphi, \psi \in \mathcal{A}$ consider $F(\gamma)=Q_{n, \Lambda}\left(\varphi^{\otimes n}, \gamma\right)$, $G(\gamma)=Q_{m, \Lambda}\left(\psi^{\otimes m}, \gamma\right)$. By (4.1) and condition (4.4) one has $\left(H_{A_{\sigma}}^{P} F, G\right)_{L^{2}\left(\pi_{\sigma}^{\Lambda}\right)}-\left(F, H_{A_{\sigma}}^{P} G\right)_{L^{2}\left(\pi_{\sigma}^{\Lambda}\right)}$

$$
=\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda}\left[n \frac{\partial}{\partial n^{a}} \varphi(s) Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right)\right.
$$

$$
\times\left(Q_{m, \Lambda}\left(\psi^{\otimes m}, \gamma\right)+m \psi(s) Q_{m-1, \Lambda}\left(\psi^{\otimes(m-1)}, \gamma\right)\right)
$$

$$
-\left(Q_{n, \Lambda}\left(\varphi^{\otimes n}, \gamma\right)+n \varphi(s) Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right)\right)
$$

$$
\left.\times m \frac{\partial}{\partial n^{a}} \psi(s) Q_{m-1, \Lambda}\left(\psi^{\otimes(m-1)}, \gamma\right)\right] \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

$$
=n \int_{\Gamma_{\Lambda}} Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right) Q_{m, \Lambda}\left(\psi^{\otimes m}, \gamma\right) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

$$
\times \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \tilde{\sigma}(d s)
$$

$$
+n m \int_{\Gamma_{\Lambda}} Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right) Q_{m-1, \Lambda}\left(\psi^{\otimes(m-1)}, \gamma\right) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

$$
\times \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \cdot \psi(s) \tilde{\sigma}(d s)
$$

$$
-m \int_{\Gamma_{\Lambda}} Q_{n, \Lambda}\left(\varphi^{\otimes n}, \gamma\right) Q_{m-1, \Lambda}\left(\psi^{\otimes(m-1)}, \gamma\right) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

$$
\times \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \psi(s) \tilde{\sigma}(d s)
$$

$$
-n m \int_{\Gamma_{\Lambda}} Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right) Q_{m-1, \Lambda}\left(\psi^{\otimes(m-1)}, \gamma\right) \pi_{\sigma}^{\Lambda}(d \gamma)
$$

$$
\times \int_{\partial \Lambda} \varphi(s) \cdot \frac{\partial}{\partial n^{a}} \psi(s) \tilde{\sigma}(d s)
$$

$$
=0
$$

by (2.8) and condition (5.4).
Next, we have by $(\overline{3.14}),(3.15)$ and (4.4)

$$
\begin{aligned}
\left(H_{A_{\sigma}}^{P} F\right)(\gamma)-\left(H_{A_{\sigma}}^{P, \Lambda} F\right)(\gamma) & =\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} F\left(\gamma+\varepsilon_{s}\right) \tilde{\sigma}(d s) \\
& =n Q_{n-1, \Lambda}\left(\varphi^{\otimes(n-1)}, \gamma\right) \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \tilde{\sigma}(d s)=0
\end{aligned}
$$

by condition (5.4).
Conversely, let the operator $\left(H_{A_{\sigma}}^{P}, \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right)$ be symmetric in $L^{2}\left(\Gamma_{\Lambda}\right)$. For any $\varphi, \psi \in \mathcal{A}$ consider the functions $F(\gamma)=Q_{1, \Lambda}(\varphi, \gamma)=\langle\varphi, \gamma\rangle-\langle\varphi\rangle_{\sigma, \Lambda}$ and $G(\gamma)=Q_{1, \Lambda}(\psi, \gamma)=\langle\psi, \gamma\rangle-\langle\psi\rangle_{\sigma, \Lambda}$. Then $F, G \in \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)$. Clearly,

$$
\int_{\Gamma_{\Lambda}} Q_{1, \Lambda}(\varphi, \gamma) \pi_{\sigma}^{\Lambda}(d \gamma)=\int_{\Gamma_{\Lambda}} Q_{1, \Lambda}(\psi, \gamma) \pi_{\sigma}^{\Lambda}(d \gamma)=0 .
$$

By (4.1) and (4.4) we have

$$
\begin{aligned}
0 & =\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda}\left(\frac{\partial}{\partial n^{a}} \varphi(s) \cdot\left(Q_{1, \Lambda}(\psi, \gamma)+\psi(s)\right)\right. \\
& \left.-\left(Q_{1, \Lambda}(\varphi, \gamma)+\varphi(s)\right) \cdot \frac{\partial}{\partial n^{a}} \psi(s)\right) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& =\int_{\partial \Lambda}\left(\frac{\partial}{\partial n^{a}} \varphi(s) \cdot \psi(s)-\varphi(s) \cdot \frac{\partial}{\partial n^{a}} \psi(s)\right) \tilde{\sigma}(d s) \\
& =\int_{\Lambda}\left(\left(A_{\sigma} \varphi\right)(x) \cdot \psi(x)-\varphi(x) \cdot\left(A_{\sigma} \psi\right)(x)\right) \sigma(d x) .
\end{aligned}
$$

Now set for any $\varphi \in \mathcal{A} F(\gamma)=Q_{1, \Lambda}(\varphi, \gamma), G(\gamma)=Q_{0, \Lambda} \equiv 1$. Then by (4.1) and (4.4) one has

$$
0=\int_{\Gamma_{\Lambda}} \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \tilde{\sigma}(d s) \pi_{\sigma}^{\Lambda}(d \gamma)=\int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi(s) \tilde{\sigma}(d s) .
$$

Example 5.5. The simple example of such algebra $\mathcal{A}$ is the algebra $\mathcal{D}_{\mathcal{N}}$ of functions, which satisfy Neumann-type boundary condition:

$$
\begin{equation*}
\frac{\partial}{\partial n^{a}} \varphi(s)=0, \quad s \in \partial \Lambda, \quad \varphi \in \mathcal{D}_{\mathcal{N}} . \tag{5.5}
\end{equation*}
$$

It's easy to see that for any $F \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}_{\mathcal{N}}, \Gamma_{\Lambda}\right), G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right)$

$$
\begin{aligned}
\mathcal{E}_{a, \pi_{\sigma}^{\Lambda}}^{\Gamma}(F, G) & =\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\left(H_{c}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) .
\end{aligned}
$$

Corollary 5.6. Suppose that $\left(A_{\sigma}, \mathcal{A}\right)$ is the essentially self-adjoint operator in $L^{2}(\Lambda)$. Then $\left(H_{A_{\sigma}}^{P}, \mathcal{F P}\left(\mathcal{A}, \Gamma_{\Lambda}\right)\right)$ is the essentially self-adjoint operator in $L^{2}\left(\Gamma_{\Lambda}\right)$ if and only if for any $\varphi \in \mathcal{A}$ the condition (5.4) is fulfilled.

Proof. The result follows immediately from Corollary 5.4 and the fact that the second quantization of an essentially self-adjoint operator is the essentially selfadjoint operator in the corresponding Hilbert space.

Now we find the explicit formula for the action of the Friedrichs extension of $H_{\min }$ on the smooth cylinder functions. Denote by $\mathcal{D}_{D}$ the set of all functions from $\mathcal{D}$ satisfying the Dirichlet boundary condition on $\partial \Lambda$ and let $A_{\sigma, D}$ be the Friedrichs extension of $\left(A_{\sigma}, \mathcal{D}(\Lambda)\right)$. The following proposition gives a formula for the action of $H_{A_{\sigma, D}}^{P}$ on the smooth cylinder functions.

Theorem 5.7. Suppose that $\mathcal{D}_{D}$ is an essential domain for $A_{\sigma, D}$. Then the closure of the operator $\left(H_{A_{\sigma, D}}^{P}, \mathcal{F} \mathcal{P}\left(\mathcal{D}_{D}, \Gamma_{\Lambda}\right)\right)$ defined by the differential expression

$$
\begin{align*}
\left(H_{A_{\sigma, D}}^{P} F\right)(\gamma) & :=\left(H_{A_{\sigma}}^{P} F\right)(\gamma) \\
& -\sum_{j=1}^{N} \frac{\partial g_{F}}{\partial q_{j}}\left(\left\langle\varphi_{1}, \gamma\right\rangle, \ldots,\left\langle\varphi_{N}, \gamma\right\rangle\right) \cdot \int_{\partial \Lambda} \frac{\partial}{\partial n^{a}} \varphi_{j}(s) \tilde{\sigma}(d s), \tag{5.6}
\end{align*}
$$

coincides with the Friedrichs extension of $H_{\min }$ in $L^{2}\left(\Gamma_{\Lambda}\right)$. Moreover, for any $F, G \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}_{D}, \Gamma_{\Lambda}\right)$

$$
\begin{aligned}
\mathcal{E}_{a, \pi_{\sigma}^{\Lambda}}^{\Gamma}(F, G) & =\int_{\Gamma_{\Lambda}}\left(H_{A_{\sigma, D}}^{P} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\left(H_{c}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{aligned}
$$

Proof. First, we recall that, for $F \in \mathcal{F} \mathcal{P}\left(\mathcal{D}(\Lambda), \Gamma_{\Lambda}\right)$

$$
H_{A_{\sigma}}^{P} F=H_{A_{\sigma}}^{P, \Lambda} F
$$

and we can consider $H_{A_{\sigma}}^{P, \Lambda}$ as the image of the second quantization of the symmetric (in $L^{2}(\Lambda)$ ) operator $\left(A_{\sigma}, \mathcal{D}(\Lambda)\right)$. Therefore, the Friedrichs extension $H_{A_{\sigma, D}}^{P}$ of the minimal operator $H_{\min }$ is the image of the second quantization of $A_{\sigma, D}$. In particular, $\mathcal{F} \mathcal{P}\left(\mathcal{D}_{D}, \Gamma_{\Lambda}\right)$ is an essential domain of $H_{A_{\sigma, D}}^{P}$. So, the assertion directly follows from Proposition 5.1, the operator inclusion $H_{\text {min }} \subset H_{A_{\sigma, D}}^{P} \subset H_{\text {max }}$ and the fact that for $F \in \mathcal{F} \mathcal{P}\left(\mathcal{D}_{D}, \Gamma_{\Lambda}\right)$

$$
H_{A_{\sigma, D}}^{P} F=H_{A_{\sigma}}^{P, \Lambda} F=H_{\max } F
$$

The last statement is a direct consequence of (4.5).

Remark 5.8. Note that for any $F \in \mathcal{F} C_{p}^{\infty}\left(\mathcal{D}, \Gamma_{\Lambda}\right), G \in F C_{p}^{\infty}\left(\mathcal{D}_{D}, \Gamma_{\Lambda}\right)$

$$
\begin{aligned}
\mathcal{E}_{a, \pi_{\sigma}^{\Lambda}}^{\Gamma}(F, G) & =\int_{\Gamma_{\Lambda}}\left(H_{\max } F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma) \\
& -\int_{\Gamma_{\Lambda}}\left(H_{c}^{P, \Lambda} F\right)(\gamma) \cdot G(\gamma) \pi_{\sigma}^{\Lambda}(d \gamma)
\end{aligned}
$$

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