

Connections Between Symmetries and Conservation Laws

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Abstract. This paper presents recent work on connections between symmetries and conservation laws. After reviewing Noether's theorem and its limitations, we present the Direct Construction Method to show how to find directly the conservation laws for any given system of differential equations. This method yields the multipliers for conservation laws as well as an integral formula for corresponding conserved densities. The action of a symmetry (discrete or continuous) on a conservation law yields conservation laws. Conservation laws yield non-locally related systems that, in turn, can yield nonlocal symmetries and in addition be useful for the application of other mathematical methods. From its admitted symmetries or multipliers for conservation laws, one can determine whether or not a given system of differential equations can be linearized by an invertible transformation.

Key words: conservation laws; linearization; nonlocal symmetries; Noether's theorem

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1 Introduction

In this paper, we discuss several connections between symmetries and conservation laws. In the classical Noether's theorem [1], if a given system of differential equations (DEs) has a variational principle, then a continuous symmetry (point, contact or higher order) that leaves invariant the action functional to within a divergence yields a conservation law [2, 3, 4, 5]. A system of DEs (as written) has a variational principle if and only if its linearized system is self-adjoint [4]. There are several limitations to Noether's theorem: It is restricted to variational systems. Consequently, to be applicable to a given system as written, the given system must be of even order, have the same number of dependent variables as the number of equations in the system and have no dissipation. There is also the difficulty of finding symmetries admitted by the action functional. Moreover, the use of Noether's theorem to find conservation laws is coordinate-dependent.

The Direct Construction Method (DCM) gets around the limitations of Noether's theorem for finding conservation laws. A given system of DEs has a conservation law if and only if all Euler operators associated with its dependent variables annihilate the scalar product of multipliers with each DE where in the resulting expression the solutions of the system of DEs are replaced by arbitrary functions. The DCM [6, 7, 8, 9] yields an integral formula for the corresponding conservation law for any such set of multipliers. On the solution space of the given system of DEs, multipliers are symmetries provided its linearized system is self-adjoint; otherwise they are solutions of the adjoint of the linearized system. The use of adjoint linearization to find multipliers for conservation laws is extensively discussed in [10] and references therein. A comprehensive discussion of the computational aspects of the DCM and comparisons with three other approaches to calculate conservation laws appears in [11].

For any system of DEs, a symmetry (point or contact) of the system maps conservation laws to conservation laws. In particular, any symmetry of the system induces a symmetry that leaves invariant the determining equations for multipliers of its conservation laws. Moreover, there is a formula to determine the action of a symmetry (discrete or continuous) on the multipliers of a given conservation law and from this action one can determine in advance whether or not the resulting conservation law will be new [12].

For any system of DEs, useful conservation laws yield equivalent non-locally related systems of DEs [13]. Such non-locally related systems (and some of their subsystems) in turn can yield nonlocal conservation laws and nonlocal symmetries for the given system. This leads to the important problem of finding extended trees of equivalent non-locally related systems of DEs [14].

For any system of DEs, from its admitted symmetry group one can determine whether or not it can be mapped into a linear system and find the explicit mapping [5, 15, 16, 17]; from its admitted multipliers for conservation laws, one can also determine whether or not it can be mapped into a linear system—here the solution of the determining system for multipliers yields the adjoint of the resulting linear system [13, 18].

This paper is organized as follows. In section two we review Noether's theorem and discuss its limitations. In section three we present the Direct Construction Method to find conservation laws. In section four we consider the use of symmetries to find new conservation laws from known conservation laws. In section five we show how conservation laws lead to equivalent non-locally related systems for finding nonlocal symmetries. In section six we compare the roles of symmetries and multipliers for conservation laws in determining whether or not a given system can be linearized by an invertible transformation. Finally, several examples are presented in section seven to illustrate the work discussed in this paper.

2 Review of Noether's theorem

Consider a functional $J[U]$ in terms of n independent variables $x = (x_1, \dots, x_n)$ and m arbitrary functions $U = (U^1(x), \dots, U^m(x))$ of the independent variables x defined on a domain Ω ,

$$J[U] = \int_{\Omega} L[U] dx = \int_{\Omega} L(x, U, \partial U, \dots, \partial^k U) dx, \quad (1)$$

with x derivatives of U up to some order k . We let ∂U , $\partial^2 U$, etc. denote all derivatives of U^γ of a given order with respect to x_i . We denote partial derivatives $\partial/\partial x_i$ by subscripts i , i.e., $U_i^\gamma = \frac{\partial U^\gamma}{\partial x_i}$, $U_{ij}^\gamma = \frac{\partial^2 U^\gamma}{\partial x_i \partial x_j}$, etc. Corresponding total derivatives are denoted by $D_i = \frac{\partial}{\partial x_i} + U_i^\gamma \frac{\partial}{\partial U^\gamma} + \dots + U_{i i_1 \dots i_j}^\gamma \frac{\partial}{\partial U_{i_1 \dots i_j}^\gamma} + \dots$. We let $F[U]$ denote a function F depending on x , U and derivatives of U with respect to x . Throughout this paper we use the summation convention for any repeated indices. The *Euler operator* with respect to U^γ is denoted by

$$E_{U^\gamma} = \frac{\partial}{\partial U^\gamma} - D_i \frac{\partial}{\partial U_i^\gamma} + \dots + (-1)^j D_{i_1} \dots D_{i_j} \frac{\partial}{\partial U_{i_1 \dots i_j}^\gamma} + \dots$$

From the calculus of variations, it follows that if $U = u$ is an extremum of $J[U]$, then

$$(E_{U^\gamma} L[U])|_{U=u} = 0, \quad \gamma = 1, \dots, m. \quad (2)$$

The function $L[U]$ is called a *Lagrangian* and the integral $J[U]$ of (1) an *action* (or *variational integral*). The extremal system of differential equations (2) with independent variables x and dependent variables $u = (u^1(x), \dots, u^m(x))$ are called the *Euler–Lagrange equations* for an extremum $U = u(x)$ of $J[U]$.

Now consider the Euler–Lagrange equations (2) as a given system of differential equations

$$G_\gamma[u] = E_{u^\gamma} L[u] = 0, \quad \gamma = 1, \dots, m \quad (3)$$

whose solutions $u(x)$ are extrema of some action integral (1). A *local conservation law* of system (3) is a divergence expression

$$D_i \Phi^i[u] = 0 \quad (4)$$

that holds for all solutions $u(x)$ of system (3); $\Phi^i[u]$, $i = 1, \dots, n$ are called the *conserved densities*.

Consider a local symmetry generator in evolutionary form [3, 4, 5]

$$X = \eta^\sigma[U] \frac{\partial}{\partial U^\sigma}; \quad (5)$$

X defines a *point symmetry* generator if $\eta^\sigma[U]$ is of the linear form $\eta^\sigma[U] = \eta^\sigma(x, U) + \xi_i(x, U) U_i^\sigma$, $\sigma = 1, \dots, m$; X defines a *contact symmetry* generator if $m = 1$ and $\eta^1[U] = \eta[U]$ depends at most on first partials of U , i.e. $\eta[U] = \eta(x, U, \partial U)$; X defines a *higher order symmetry* generator if it is neither a point symmetry nor contact symmetry generator. Higher order symmetries were first mentioned by Emmy Noether in her celebrated paper [1]. The p th extension of a local symmetry generator (5) is given by

$$X^{(p)} = \eta^\sigma[U] \frac{\partial}{\partial U^\sigma} + (D_i \eta^\sigma[U]) \frac{\partial}{\partial U_i^\sigma} + \dots + (D_{i_1} \dots D_{i_p} \eta^\sigma[U]) \frac{\partial}{\partial U_{i_1 \dots i_p}^\sigma}.$$

A local symmetry generator (5) defines a *variational symmetry* of the action integral (1) if and only if

$$X^{(k)} L[U] = D_i f^i[U] \quad (6)$$

holds for some functions $f^i[U]$. If (5) defines a variational symmetry of (1), then one can derive the identity [3, 4, 5]

$$\eta^\sigma[U] G_\sigma[U] = \eta^\sigma[U] E_{U^\sigma} L[U] = D_i (f^i[U] - W^i[U, \eta[U]]) \quad (7)$$

where

$$W^i[U, \eta[U]] = \left(\eta^\sigma[U] E_{U_i^\sigma} + \sum_{j=2} (D_{i_1} \dots D_{i_{j-1}} \eta^\sigma[U]) E_{U_{i_1 \dots i_{j-1}}^\sigma} \right) L[U].$$

Consequently, a variational symmetry (5) of the action integral (1) yields the conservation law

$$D_i \Phi^i[u] = D_i (W^i[u, \eta[u]] - f^i[u]) = 0$$

of the given system of differential equations arising from the extrema of the action integral (1). In particular, from (7) we see that Noether's theorem yields a conservation law of the variational system (3) for any set of *multipliers* $\{\eta^\sigma[U]\}$ for which $\eta^\sigma[U] \frac{\partial}{\partial U^\sigma}$ is a variational symmetry of the action integral (1).

Since a variational symmetry leaves the Lagrangian $L[U]$ invariant to within a divergence, it follows that a variational symmetry of the action integral $J[U]$ leaves invariant its family of extrema $U = u$. Hence a variational symmetry (5) of the action integral (1) yields a symmetry $X = \eta^\sigma[u] \frac{\partial}{\partial u^\sigma}$ of the system of differential equations (3), i.e.,

$$X^{(2k)} G_\gamma[U] \Big|_{U=u} = 0, \quad \gamma = 1, \dots, m. \quad (8)$$

Conversely, a symmetry $X = \eta^\sigma[u] \frac{\partial}{\partial u^\sigma}$ of a system of DEs (3) derivable from an action integral (1) does not necessarily yield a variational symmetry $X = \eta^\sigma[U] \frac{\partial}{\partial U^\sigma}$ of the action integral (1).

2.1 Limitations of Noether's theorem

1. *The difficulty of finding variational symmetries.* To find variational symmetries for a given variational system (3), first one determines local symmetries $X = \eta^\sigma[u] \frac{\partial}{\partial u^\sigma}$ admitted by the Euler–Lagrange equations (3) through the extension of Lie's algorithm to include higher order symmetries. Then for each such admitted symmetry, one checks if $X = \eta^\sigma[U] \frac{\partial}{\partial U^\sigma}$ leaves invariant the Lagrangian $L[U]$ to within a divergence, i.e., equation (6) is satisfied.
2. *A given system of DEs is not variational as written.* A given system of differential equations, as written, is variational if and only if its linearized system (Fréchet derivative) is self-adjoint [4]. Consequently, it is necessary that a given system of DEs, as written, must be of even order, have the same number of equations in the system as its number of dependent variables and be non-dissipative to directly admit a variational principle.
3. *Artifices may make a given system of DEs variational.* Such artifices include:

- The use of multipliers. As an example, the PDE

$$u_{tt} + H'(u_x)u_{xx} + H(u_x) = 0,$$

as written, does not admit a variational principle since its linearized equation $\zeta_{tt} + H'(u_x)\zeta_{xx} + (H''(u_x) + H'(u_x))\zeta_x = 0$ is not self-adjoint. However, the equivalent PDE

$$e^x[u_{tt} + H'(u_x)u_{xx} + H(u_x)] = 0,$$

as written, is self-adjoint!

- The use of a contact transformation of the variables. As an example, the ODE

$$y'' + 2y' + y = 0, \tag{9}$$

as written, obviously does not admit a variational principle. But the point transformation $x \rightarrow X = x$, $y \rightarrow Y = ye^x$, maps ODE (9) into the variational ODE $Y'' = 0$. It is well-known that every second order ODE, written in solved form, can be mapped into $Y'' = 0$ by some contact transformation but there is no finite algorithm to find such a transformation.

- The use of a differential substitution. As an example, the Korteweg–de Vries (KdV) equation

$$u_{xxx} + uu_x + u_t = 0,$$

as written, obviously does not admit a variational principle since it is of odd order. But the well-known differential substitution $u = v_x$ yields the related transformed KdV equation $v_{xxxx} + v_x v_{xx} + v_{xt} = 0$ which is the Euler–Lagrange equation for an extremum $V = v$ of the action integral with Lagrangian $L[V] = \frac{1}{2}(V_{xx})^2 - \frac{1}{6}(V_x)^3 - \frac{1}{2}V_x V_t$.

4. *Noether's theorem is coordinate-dependent.* The use of Noether's theorem to obtain a conservation law is coordinate-dependent since the action of a contact transformation can transform a DE having a variational principle to one that does not have one. On the other hand it is well-known that conservation laws are coordinate-independent in the sense that a contact transformation maps a conservation law into a conservation law (see, for example [12].)
5. *Artifice of a Lagrangian.* One should be able to find the conservation laws of a given system of DEs directly without the need to find a related functional whether or not the given system is variational.

3 Direct construction method to find conservation laws

Through the Direct Construction Method, one can get around the limitations of Noether's theorem. Here one works directly with a given system of DEs whether or not its linearized system is self-adjoint. The starting point is the observation from equation (7) that conservation laws obtained through Noether's theorem arise from multipliers that must be local symmetries of the given system of DEs (3). Moreover (from equations (8)), the local symmetries of a system of DEs are the solutions of its linearized system. Now suppose we have a given system of M differential equations

$$G_\alpha[u] = G_\alpha(x, u, \partial u, \dots, \partial^K u) = 0, \quad \alpha = 1, \dots, M \quad (10)$$

with n independent variables $x = (x_1, \dots, x_n)$ and m dependent variables $u = (u^1, \dots, u^m)$. It is understood that the system of DEs (10) includes all its differential consequences. The local symmetry $X = \eta^\sigma[u] \frac{\partial}{\partial u^\sigma}$ is admitted by (10) if and only if

$$X^{(K)} G_\alpha[u] \Big|_{G_\gamma[u]=0, \gamma=1, \dots, M} = 0, \quad \alpha = 1, \dots, M. \quad (11)$$

Equations (11) are commonly called the linear *determining equations* for the local symmetries admitted by the given system (10). Let L be the linearizing operator (Fréchet derivative) for system (10) that arises from equations (11). Then the determining equations (11) can be expressed as a linear system

$$L\eta[u] = 0. \quad (12)$$

The given system (10), as written, has a variational principle if and only if L is a self-adjoint operator, i.e., $L^* = L$ where L^* is the *adjoint* of L .

A set of *multipliers* $\{\Lambda^\sigma[U]\}$ yields a conservation law of the given system (10) if and only if

$$\Lambda^\sigma[U] G_\sigma[U] = D_i \Phi^i[U]$$

for some $\{\Phi^i[U]\}$. Then the conservation law $D_i \Phi^i[u] = 0$ holds if $U = u$ solves the given system (10). The aim is to first find all such sets of multipliers $\{\Lambda^\sigma[U]\}$ for conservation laws of (10) and then obtain the corresponding set of conserved densities $\{\Phi^i[u]\}$ for each set of multipliers. The determining equations for multipliers result from the following theorem that can be proved by direct computation.

Theorem 1. *A set of multipliers $\{\Lambda^\sigma[U]\}$ yields a conservation law of the given system of differential equations (10) if and only if the equations*

$$E_{U^\gamma}(\Lambda^\sigma[U] G_\sigma[U]) = 0, \quad \gamma = 1, \dots, m \quad (13)$$

hold for arbitrary functions $U = (U^1(x), \dots, U^m(x))$.

The set of equations (13) are the linear determining equations for the multipliers for conservation laws admitted by the given system (10). One can show [4, 6, 7, 8, 9, 19] that if $\{\Lambda^\sigma[U]\}$ yields a conservation law for system (10), then for the adjoint operator L^* of the linearization operator L arising from the linear determining equations (12) for local symmetries of system (10), one has

$$L^* \Lambda[u] = 0.$$

If $L^* = L$, then $\Lambda^\sigma[u] \frac{\partial}{\partial u^\sigma}$ is a local symmetry admitted by system (10); if $L^* \neq L$, then $\Lambda^\sigma[u] \frac{\partial}{\partial u^\sigma}$ is not a local symmetry admitted by system (10). When $L^* = L$, the determining equations (13) include as a subset the set of determining equations for local symmetries admitted by system (10) and in this case any solution of the determining equations (13) must be a variational symmetry. Most importantly, given a set of multipliers $\{\Lambda^\sigma[U]\}$ yielding a conservation law for system (10), the integral formulas for the corresponding conserved densities $\{\Phi^i[u]\}$ are given in [6, 7, 8, 9, 19].

4 Use of symmetries to find new conservation laws from known conservation laws

Any symmetry (discrete or continuous) admitted by a given system of DEs (10) maps a conservation law of (10) into another conservation law of (10). Usually, the same conservation law of (10) is obtained. An admitted symmetry of system (10) induces a symmetry that leaves invariant the linear determining system (13) for its multipliers. Hence, it follows that if we determine the action of a symmetry on a set of multipliers $\{\Lambda^\sigma[U]\}$ for a known conservation law of (10) to obtain another set of multipliers $\{\hat{\Lambda}^\sigma[U]\}$, then a priori we see whether or not a new conservation is obtained for (10).

Suppose the invertible point transformation

$$x = x(\tilde{x}, \tilde{U}), \quad U = U(\tilde{x}, \tilde{U}), \quad (14)$$

with inverse

$$\tilde{x} = \tilde{x}(x, U), \quad \tilde{U} = \tilde{U}(x, U),$$

is admitted by a given system of DEs (10). Then

$$G_\alpha[U] = A_\alpha^\beta[\tilde{U}]G_\beta[\tilde{U}] \quad (15)$$

holds for some $\{A_\alpha^\beta[\tilde{U}]\}$. One can easily prove the following theorem [12, 20].

Theorem 2. *Under the point transformation (14), there exist functions $\Psi^i[\tilde{U}]$ such that*

$$J[\tilde{U}]D_i\Phi^i[U] = \tilde{D}_i\Psi^i[\tilde{U}] \quad (16)$$

where the Jacobian determinant

$$J[\tilde{U}] = \frac{D(x_1, \dots, x_n)}{D(\tilde{x}_1, \dots, \tilde{x}_n)} = \begin{vmatrix} \tilde{D}_1x_1 & \tilde{D}_1x_2 & \cdots & \tilde{D}_1x_n \\ \tilde{D}_2x_1 & \tilde{D}_2x_2 & \cdots & \tilde{D}_2x_n \\ \vdots & \vdots & \ddots & \vdots \\ \tilde{D}_nx_1 & \tilde{D}_nx_2 & \cdots & \tilde{D}_nx_n \end{vmatrix} \quad (17)$$

and

$$\Psi^{i_1}[\tilde{U}] = \pm \begin{vmatrix} \Phi^1[U] & \Phi^2[U] & \cdots & \Phi^n[U] \\ \tilde{D}_{i_2}x_1 & \tilde{D}_{i_2}x_2 & \cdots & \tilde{D}_{i_2}x_n \\ \vdots & \vdots & \ddots & \vdots \\ \tilde{D}_{i_n}x_1 & \tilde{D}_{i_n}x_2 & \cdots & \tilde{D}_{i_n}x_n \end{vmatrix}. \quad (18)$$

In (18), (i_1, i_2, \dots, i_n) denotes all cyclic permutations of the indices $1, 2, \dots, n$ and the \pm sign is chosen according to whether the permutation has even or odd order, i.e. $(i_1, i_2, \dots, i_n) = (1, 2, \dots, n), (2, 3, \dots, 1), \dots, (n, 1, \dots, n-1)$ are alternately even and odd. The coordinates of $\Phi^i[U]$ in (16) are expressed in terms of the point transformation (14).

The following corollary results immediately from Theorem 2.

Corollary 1. *If the point transformation (14) is admitted by the DE system (10), then a conservation law (4) of system (10) is mapped to a conservation law*

$$D_i\Psi^i[u] = 0$$

of system (10) with conserved densities $\Psi^i[u]$.

In [3, 4, 21] it is shown how to obtain a conservation law from a known conservation law through the action of the infinitesimal generator of an admitted continuous symmetry (in evolutionary (characteristic) form) on the given conservation law. Here it is necessary to use differential consequences of the given system (10) even in the case when the symmetry is equivalent to a point or contact symmetry. A priori, one is unable to determine whether or not a new conservation law is obtained.

Now suppose the set of multipliers $\{\Lambda^\sigma[U]\}$ yields a conservation law with densities $\Phi^i[u]$ of a given system of DEs (10), i.e.,

$$\Lambda^\sigma[U]G_\sigma[U] = D_i\Phi^i[U]. \quad (19)$$

One can prove the following theorem [12].

Theorem 3. *Suppose the point transformation (14) is admitted by (10) and $\{\Lambda^\sigma[U]\}$ is a set of multipliers for a conservation law of (10) with conserved densities $\Phi^i[u]$. Then*

$$\hat{\Lambda}^\beta[\tilde{U}]G_\beta[\tilde{U}] = \tilde{D}_i\Psi^i[\tilde{U}] \quad (20)$$

where

$$\hat{\Lambda}^\beta[\tilde{U}] = J[\tilde{U}]A_\alpha^\beta[\tilde{U}]\Lambda^\alpha[U], \quad \beta = 1, \dots, M, \quad (21)$$

with the coordinates of derivative terms in $\Lambda^\alpha[U]$ expressed in terms of the natural extension of transformation (14) to the derivatives of U [3, 4, 5, 19]. In (20), $\Psi^i[\tilde{U}]$ are given by the determinant (18) and, in (21), $A_\alpha^\beta[\tilde{U}]$ are given by (15) and the Jacobian determinant $J[\tilde{U}]$ is given by (17).

After replacing the coordinates \tilde{x}_i by x_i , \tilde{U}^σ by U^σ , \tilde{U}_i^σ by U_i^σ , etc. in (20), one obtains the following corollary.

Corollary 2. *If $\{\Lambda^\alpha[U]\}$ is a set of multipliers yielding a conservation law of system (10) admitting the point transformation (14), then $\{\hat{\Lambda}^\beta[U]\}$ yields a set of multipliers for a conservation law of (10) where $\hat{\Lambda}^\beta[U]$ is given by (21) after replacing \tilde{x}_i by x_i , \tilde{U}^σ by U^σ , \tilde{U}_i^σ by U_i^σ , etc. The set of multipliers $\{\hat{\Lambda}^\beta[U]\}$ yields a new conservation law of system (10) if and only if this set is nontrivial on all solutions $U = u(x)$ of system (10), i.e. $\hat{\Lambda}^\beta[u] \neq c\Lambda^\beta[u]$, $\beta = 1, \dots, M$, for some constant c .*

Now suppose the point transformation (14) is a one-parameter (ε) Lie group of point transformations

$$x = x(\tilde{x}, \tilde{U}; \varepsilon) = e^{\varepsilon\tilde{X}}\tilde{x}, \quad U = U(\tilde{x}, \tilde{U}; \varepsilon) = e^{\varepsilon\tilde{X}}\tilde{U} \quad (22)$$

in terms of its infinitesimal generator (and extensions) $\tilde{X} = \xi_j(\tilde{x}, \tilde{U})\frac{\partial}{\partial\tilde{x}_j} + \eta^\sigma(\tilde{x}, \tilde{U})\frac{\partial}{\partial\tilde{U}^\sigma}$. If (19) holds, then from (16) and the Lie group properties of (22), it follows that

$$J[U; \varepsilon]e^{\varepsilon X}(\Lambda^\sigma[U]G_\sigma[U]) = D_i\Psi^i[U; \varepsilon] \quad (23)$$

in terms of the (extended) infinitesimal generator $X = \xi_j(x, U)\frac{\partial}{\partial x_j} + \eta^\sigma(x, U)\frac{\partial}{\partial U^\sigma}$.

Now assume that the Lie group of point transformations (22) is admitted by a given system of DEs (10). Then for some $\{a_{\rho\alpha}^\beta[U]\}$, one has

$$e^{\varepsilon X}G_\alpha[U] = \sum \varepsilon^\rho a_{\rho\alpha}^\beta G_\beta[U].$$

Then, after expanding both sides of equation (23) in terms of power series in ε , one obtains an expression of the form

$$\sum \varepsilon^p \hat{\Lambda}^\sigma[U; p] G_\sigma[U] = \sum \varepsilon^p D_i \left(\frac{1}{p!} \frac{d^p}{d\varepsilon^p} \Psi^i[U; \varepsilon] \right) \Big|_{\varepsilon=0}. \quad (24)$$

Corresponding to the sequence of sets of multipliers $\{\hat{\Lambda}^\sigma[U; p]\}$, $p = 1, 2, \dots$, arising in expression (24), one obtains a sequence of conservation laws

$$D_i \left(\frac{d^p}{d\varepsilon^p} \Psi^i[u; \varepsilon] \right) \Big|_{\varepsilon=0} = 0, \quad p = 1, 2, \dots$$

for system (10) from its known conservation law $D_i \Phi^i[u] = 0$.

The results presented in this section extend to contact transformations in the case of a scalar DE [12].

5 Connections between conservation laws and nonlocal symmetries

Through useful conservation laws, one can obtain a tree of non-locally related systems and subsystems for a given system of DEs. Such non-locally related systems lead to the possibility of using Lie's algorithm to seek nonlocal symmetries for the given system. For details of the work discussed in this section see [5, 13, 14] and references therein.

In this section, let $x = (x_1, x_2) = (t, x)$; $u = (u^1, \dots, u^m)$. Consider a given system of PDEs

$$G_\alpha[u] = G_\alpha(t, x, u, \partial u, \dots, \partial^k u) = 0, \quad \alpha = 1, \dots, M. \quad (25)$$

Suppose $G_1[u]$ is written as a conservation law:

$$G_1[u] = D_t T[u] + D_x X[u] = 0. \quad (26)$$

Then through the conservation law (26), one can introduce an auxiliary potential dependent variable v . It is easy to see that system (25) is *non-locally* equivalent to the potential system

$$\begin{aligned} v_t + X[u] &= 0, \\ v_x - T[u] &= 0, \\ G_\alpha[u] &= 0, \quad \alpha = 2, \dots, M. \end{aligned} \quad (27)$$

Suppose a point symmetry (not in evolutionary form)

$$X = \xi(x, t, u, v) \frac{\partial}{\partial x} + \tau(x, t, u, v) \frac{\partial}{\partial t} + \eta^j(x, t, u, v) \frac{\partial}{\partial u^j} + \phi(x, t, u, v) \frac{\partial}{\partial v}$$

is admitted by the equivalent potential system (27). Then X yields a *nonlocal* symmetry of the given system (25) if and only if $(\xi_v)^2 + (\tau_v)^2 + \sum_{j=1}^m (\eta_v^j)^2 \neq 0$.

Now suppose none of the equations of a given system of PDEs (25) is written as a conservation law. Suppose the set of multipliers $\{\Lambda^\sigma[U]\}$ yields a conservation law of system (25), i.e.,

$$\Lambda^\sigma[U] G_\sigma[U] = D_t T[U] + D_x X[U]. \quad (28)$$

Suppose for some β , $\Lambda^\beta[U] \neq 0$ and suppose every solution of the system $\Lambda^\beta[u] = 0$, $G_\sigma[u] = 0$, $\sigma = 1, \dots, M$ with $\sigma \neq \beta$, is also a solution of the given system of PDEs (25). Then the conservation law (28) is said to be a *useful* conservation law of system (25). In the case of a useful conservation law (without loss of generality $\beta = 1$), the given system of PDEs (25) is non-locally equivalent to a resulting potential system of the form (27) since this equivalence follows when one replaces $\Lambda^1[u] G_1[u] = 0$ by the potential equations (27a,b).

6 Connections between symmetries and conservation laws for linearization

In [5, 16], the following two theorems are proven for determining whether or not a given nonlinear system of PDEs (10) can be invertibly mapped to some linear system of PDEs and to find such a mapping when it exists in terms of admitted symmetries of the nonlinear system. Here we present the theorems when the number of dependent variables $m \geq 2$.

Theorem 4 (Necessary conditions for the existence of an invertible mapping). *If there exists an invertible transformation that maps a given nonlinear system of PDEs (10) to some linear system of PDEs, then*

1. *the mapping must be a point transformation of the form*

$$\begin{aligned} z_j &= \phi_j(x, u), & j &= 1, 2, \dots, n, \\ w^\gamma &= \psi^\gamma(x, u), & \gamma &= 1, 2, \dots, m; \end{aligned} \quad (29)$$

2. *the given system of PDEs (10) must admit an infinite-parameter Lie group of point transformations having an infinitesimal generator*

$$X = \xi_i[u] \frac{\partial}{\partial x_i} + \eta^\nu[u] \frac{\partial}{\partial u^\nu} \quad (30)$$

with each $\xi_i[u]$, $\eta^\nu[u]$ characterized by

$$\xi_i[u] = \sum_{\sigma=1}^m \alpha_i^\sigma(x, u) F^\sigma(x, u), \quad \eta^\nu[u] = \sum_{\sigma=1}^m \beta_\nu^\sigma(x, u) F^\sigma(x, u), \quad (31)$$

where α_i^σ , β_ν^σ are specific functions of x and u , and $F = (F^1, F^2, \dots, F^m)$ is an arbitrary solution of some linear system of PDEs

$$LF = 0 \quad (32)$$

with L representing a linear differential operator in terms of some independent variables $X = (X_1(x, u), X_2(x, u), \dots, X_n(x, u))$.

Theorem 5 (Sufficient conditions for the existence of an invertible mapping). *Suppose a given nonlinear system of PDEs (10) admits an infinitesimal generator (30) whose coefficients are of the form (31) with F being an arbitrary solution of some linear system (32) with specific independent variables $X = (X_1(x, u), X_2(x, u), \dots, X_n(x, u))$. If the linear homogeneous system of m first order PDEs for the scalar function Φ given by*

$$\alpha_i^\sigma(x, u) \frac{\partial \Phi}{\partial x_i} + \beta_\nu^\sigma(x, u) \frac{\partial \Phi}{\partial u^\nu} = 0, \quad \sigma = 1, 2, \dots, m, \quad (33)$$

has n functionally independent solutions $X_1(x, u), X_2(x, u), \dots, X_n(x, u)$, and the linear system of m^2 first order PDEs

$$\alpha_i^\sigma(x, u) \frac{\partial \psi^\gamma}{\partial x_i} + \beta_\nu^\sigma(x, u) \frac{\partial \psi^\gamma}{\partial u^\nu} = \delta^{\gamma\sigma}, \quad \gamma, \sigma = 1, 2, \dots, m, \quad (34)$$

where $\delta^{\gamma\sigma}$ is the Kronecker symbol, has a solution $\psi = (\psi^1(x, u), \psi^2(x, u), \dots, \psi^m(x, u))$, then the invertible mapping given by

$$\begin{aligned} z_j &= \phi_j(x, u) = X_j(x, u), & j &= 1, 2, \dots, n, \\ w^\gamma &= \psi^\gamma(x, u), & \gamma &= 1, 2, \dots, m, \end{aligned}$$

transforms the nonlinear system of PDEs (10) to the linear system of PDEs given by

$$Lw = g(z)$$

for some function $g(z)$.

The theorems for linearization through admitted infinite-parameter groups have their counterparts for admitted multipliers for conservation laws [13]. This follows first from the observation that for any linear operator L and its adjoint operator L^* , one has the well-known relationship

$$VLU - UL^*V = D_i P^i[U, V] \quad (35)$$

for some $P^i[U, V]$ that depend bilinearly on any functions $U(x)$, $V(x)$ and their derivatives to some finite order. Hence if a given system of PDEs (10) is linear, i.e., it is of the form

$$Lu = 0, \quad (36)$$

then *any* set of multipliers $\{\Lambda^\sigma(x)\}$ satisfying

$$L^*\Lambda(x) = 0 \quad (37)$$

yields a conservation law for system (35).

The above observation combined with the result that a point transformation (29) maps any conservation law of a given system of PDEs into a conservation law of another system of PDEs in terms of the variables of the point transformation [12], shows that if a given nonlinear system of PDEs (10) is linearizable by a point transformation (29), then it is necessary that (10) admits a set of multipliers $\{\Lambda^\sigma[U]\}$ of the form

$$\Lambda^\sigma[U] = A_\rho^\sigma[U]F^\rho(X); \quad (38)$$

A_ρ^σ are specific functions of x , U and derivatives of U to some finite order, and $X = (X_1(x, U), X_2(x, U), \dots, X_n(x, U))$ are specific functions of x and U with the property that if $U = u$ is any solution of the given nonlinear system (10), then for some linear operator L^* in terms of independent variables $z = X = (X_1(x, u), X_2(x, u), \dots, X_n(x, u))$, one has

$$L^*F = 0. \quad (39)$$

If a given nonlinear system (10) can be mapped by a point transformation (29) into a linear system $Lw = 0$, then the operator L^* in (39) is the adjoint of L .

The results presented in this section extend to contact transformations in the case of a scalar DE [5, 13, 15, 16, 17, 18].

In summary, if a given nonlinear system of PDEs can be mapped invertibly into a linear system, then its admitted symmetries yield the linear system and the mapping to the linear system whereas its admitted multipliers yield the adjoint linear system. Further details on linearization through conservation laws will appear in [18].

7 Examples

Now we illustrate all of the above connections between symmetries and conservation laws through several examples involving the nonlinear telegraph (NLT) equation

$$u_{tt} - (F(u)u_x)_x - (G(u))_x = 0. \quad (40)$$

For *any* $(F(u), G(u))$ pair, from PDE (40) written as a conservation law, we obtain potential systems

$$\begin{aligned} G_1[u, v] &= v_t - F(u)u_x - G(u) = 0, \\ G_2[u, v] &= v_x - u_t = 0; \end{aligned} \tag{41}$$

$$\begin{aligned} H_1[u, v, w] &= G_1[u, v] = 0, \\ H_2[u, v, w] &= w_t - v = 0, \\ H_3[u, v, w] &= w_x - u = 0. \end{aligned} \tag{42}$$

The point symmetry classification of the NLT scalar equation (40) appears in [18]. The point symmetry classification of the potential system (41) yields nonlocal symmetries for the NLT equation (40) for a very wide class of $(F(u), G(u))$ pairs [23]. For *specific* $(F(u), G(u))$ pairs, the conservation law classification results for potential system (41) yield additional useful conservation laws and an extended tree of non-locally related systems [14, 24]. For the potential system (41), it is interesting to compare its conservation law and point symmetry classifications [25].

7.1 Nonlocal symmetries

The point symmetry

$$X = \xi(x, t, u, v) \frac{\partial}{\partial x} + \tau(x, t, u, v) \frac{\partial}{\partial t} + \eta(x, t, u, v) \frac{\partial}{\partial u} + \phi(x, t, u, v) \frac{\partial}{\partial v} \tag{43}$$

is admitted by the potential system (41) if and only if the system of determining equations

$$\begin{aligned} \xi_v - \tau_u &= 0, \\ \eta_u - \phi_v + \xi_x - \tau_t &= 0, \\ G(u)[\eta_v + \tau_x] + \eta_t - \phi_x &= 0, \\ \xi_u - F(u)\tau_v &= 0, \\ \phi_u - G(u)\tau_u - F(u)\eta_v &= 0, \\ G(u)\xi_v + \xi_t - F(u)\tau_x &= 0, \\ F(u)[\phi_v - \tau_t + \xi_x - \eta_u - 2G(u)\tau_v] - F'(u)\eta &= 0, \\ G(u)[\phi_v - \tau_t - G(u)\tau_v] - F(u)\eta_x - G'(u)\eta + \phi_t &= 0, \end{aligned}$$

holds for *arbitrary* values of x, t, u, v . One can show [23] that the potential system (41) yields a nonlocal symmetry of the NLT equation (40) if and only if

$$\begin{aligned} (c_3u + c_4)F'(u) - 2(c_1 - c_2 - G(u))F(u) &= 0, \\ (c_3u + c_4)G'(u) + G^2(u) - (c_1 - 2c_2 + c_3)G(u) - c_5 &= 0, \end{aligned} \tag{44}$$

for arbitrary constants $c_i, i = 1, \dots, 5$. Moreover, the potential system (41) allows (40) to be linearizable (by a nonlocal transformation) if and only if in (44), one has $c_1 = 0, c_5 = c_2(c_3 - c_2)$. For any $(F(u), G(u))$ pair satisfying ODE system (44), the potential system (41) admits a point symmetry of the form (43) with

$$\begin{aligned} \xi &= c_1x + \int F(u)du, \\ \tau &= c_2t + v, \\ \eta &= c_3u + c_4, \\ \phi &= c_5t + (c_1 - c_2 + c_3)v, \end{aligned}$$

which in turn yields a nonlocal symmetry admitted by the given NLT equation (40) since τ has an essential dependence on v .

7.2 Conservation laws

The multipliers $[\Lambda_1, \Lambda_2] = [\alpha(x, t, U, V), \beta(x, t, U, V)]$ yield a conservation law of the NLT potential system (41) if and only if for the Euler operators E_U, E_V one has

$$\begin{aligned} E_U(\alpha G_1[U, V] + \beta G_2[U, V]) &\equiv 0, \\ E_V(\alpha G_1[U, V] + \beta G_2[U, V]) &\equiv 0, \end{aligned} \quad (45)$$

for *arbitrary* differentiable functions $U(x, t), V(x, t)$. Equations (45) reduce to the following determining equations for the multipliers:

$$\begin{aligned} \beta_V - \alpha_U &= 0, \\ \beta_U - F(U)\alpha_V &= 0, \\ \beta_x - \alpha_t - G(U)\alpha_V &= 0, \\ F(U)\alpha_x - \beta_t - [G(U)\alpha]_U &= 0. \end{aligned} \quad (46)$$

For any solution of the determining equations (46), one can show that the corresponding conserved densities are given by

$$\begin{aligned} X &= - \int_a^u \alpha(x, t, s, b) ds - \int_b^v \beta(x, t, u, s) ds - G(a) \int^x \alpha(s, t, a, b) ds, \\ T &= \int_a^u \beta(x, t, s, b) ds + \int_b^v \alpha(x, t, u, s) ds, \end{aligned}$$

in terms of arbitrary constants a and b .

In solving the determining equations (46), it turns out that three cases arise in terms of classifying functions $d(U)$ and $h(U)$ given by

$$\begin{aligned} d(U) &= G'^2 F''' - 3G' G'' F'' + (3G''^2 - G' G''') F', \\ h(U) &= G'^2 G^{(4)} - 4G' G'' G''' + 3G''^3. \end{aligned}$$

In particular, one must separately consider the cases $d(U) = h(U) = 0$; $d(U) \neq 0, h(U) \equiv 0$; $d(U) \neq 0, h(U) \neq 0$. For further details, see [24].

Example 1. As a first example, consider the NLT potential system

$$\begin{aligned} v_t + (1 - 2e^{2u})u_x - e^u &= 0, \\ v_x - u_t &= 0. \end{aligned} \quad (47)$$

After solving the corresponding determining equations (44), one can show that potential system (47) admits the set of multipliers

$$\begin{aligned} \alpha &= e^{-\frac{1}{2}(U+t/\sqrt{2})} \sin\left(\frac{1}{2}\left(V + \frac{x+2e^U}{\sqrt{2}}\right)\right), \\ \beta &= -e^{-\frac{1}{2}(U+t/\sqrt{2})} \left[\sqrt{2}e^U \sin\left(\frac{1}{2}\left(V + \frac{x+2e^U}{\sqrt{2}}\right)\right) + \cos\left(\frac{1}{2}\left(V + \frac{x+2e^U}{\sqrt{2}}\right)\right) \right] \end{aligned} \quad (48)$$

with corresponding densities

$$\begin{aligned} T &= -2e^{-\frac{1}{2}(u+t/\sqrt{2})} \cos\left(\frac{1}{2}\left(v + \frac{x+2e^u}{\sqrt{2}}\right)\right), \\ X &= 2e^{-\frac{1}{2}(u+t/\sqrt{2})} \left[\sqrt{2}e^u \cos\left(\frac{1}{2}\left(v + \frac{x+2e^u}{\sqrt{2}}\right)\right) - \sin\left(\frac{1}{2}\left(v + \frac{x+2e^u}{\sqrt{2}}\right)\right) \right]. \end{aligned} \quad (49)$$

Example 2. As a second example, consider the NLT potential system

$$\begin{aligned} v_t - (\operatorname{sech}^2 u)u_x + \tanh u &= 0, \\ v_x - u_t &= 0. \end{aligned} \quad (50)$$

The potential system (50) admits the set of multipliers

$$\alpha = e^x(2x + t^2 - V^2 - 2 \log(\cosh U)), \quad \beta = 2e^x(V \tanh U - t) \quad (51)$$

with corresponding densities

$$\begin{aligned} T &= e^x \left[2tu - \frac{1}{3}v^3 + v(t^2 + 2x - 2 \log(\cosh u)) \right], \\ X &= e^x [(v^2 - t^2 - 2x + 2(1 + \log(\cosh u))) \tanh u - 2(tv + u)]. \end{aligned} \quad (52)$$

7.3 Symmetry action on conservation laws to yield new conservation laws

Example 3. The NLT potential system (47) obviously admits the discrete reflection symmetry

$$(t, x, u, v) = (-\tilde{t}, \tilde{x}, \tilde{u}, -\tilde{v}) \quad (53)$$

and the one-parameter (ε) family of translations

$$(t, x, u, v) = (\tilde{t}, \tilde{x}, \tilde{u}, \tilde{v} + \varepsilon). \quad (54)$$

Applying these symmetries to (48), (49), one is able to find three more conservation laws for system (47).

The reflection symmetry (53) applied to (48) yields a second set of multipliers

$$[\alpha_2, \beta_2] = [-\alpha(x, -t, U, -V), \beta(x, -t, U, -V)] \quad (55)$$

with corresponding densities

$$[T_2, X_2] = [T(x, -t, u, -v), -X(x, -t, u, -v)]. \quad (56)$$

The translation symmetry (54) applied to (48), (49) yields a third set of multipliers (from the $O(\varepsilon)$ term in the corresponding expression (24))

$$[\alpha_3, \beta_3] = [\alpha(x, t, U, V + \pi), \beta(x, t, U, V + \pi)] \quad (57)$$

and conserved densities

$$[T_3, X_3] = [T(x, t, u, v + \pi), X(x, t, u, v + \pi)]. \quad (58)$$

Finally, the reflection symmetry (53) applied to (57), (58) yields a fourth set of multipliers $[\alpha_4, \beta_4] = [-\alpha(x, -t, U, -V - \pi), \beta(x, -t, U, -V - \pi)]$ with corresponding densities $[T_4, X_4] = [T(x, -t, u, -v - \pi), -X(x, -t, u, -v - \pi)]$.

Example 4. The NLT potential system (50) obviously admits the one-parameter (ε) family of translations

$$(t, x, u, v) = (\tilde{t} + \varepsilon, \tilde{x}, \tilde{u}, \tilde{v}) \quad (59)$$

as well as the point symmetry associated with the infinitesimal generator

$$X = v \frac{\partial}{\partial t} + \tanh u \frac{\partial}{\partial x} + \frac{\partial}{\partial u} + t \frac{\partial}{\partial v}. \quad (60)$$

Applying these two symmetries to (51), (52), one is able to generate three more conservation laws as follows.

From the $O(\varepsilon)$ and $O(\varepsilon^2)$ terms that result from the corresponding expression (24) when (59) is applied to (51), (52), one respectively obtains two new sets of multipliers $[\alpha_2, \beta_2] = [te^x, -e^x]$, $[\alpha_3, \beta_3] = [e^x, 0]$ and conserved densities $[T_2, X_2] = [e^x(tv + u), -e^x(v + t \tanh u)]$, $[T_3, X_3] = [e^x v, -e^x \tanh u]$. The $O(\varepsilon)$ term that arises from expression (24) when symmetry (60) is applied to $[\alpha_2, \beta_2]$, $[T_2, X_2]$ yields a fourth set of multipliers $[\alpha_4, \beta_4] = [e^x V, -e^x \tanh U]$ and corresponding conserved densities $[T_4, X_4] = [e^x(\frac{1}{2}v^2 + \log(\cosh u)), -e^x v \tanh u]$.

7.4 Linearization

Now we consider two examples of NLT potential systems that can be linearized and see how this can be done from admitted symmetries and observed from admitted multipliers for conservation laws.

Example 5. The quasilinear NLT potential system

$$\begin{aligned} v_t - F(u)u_x &= 0, \\ v_x - u_t &= 0, \end{aligned} \tag{61}$$

admits the point symmetry $X = A(u, v)\frac{\partial}{\partial x} + B(u, v)\frac{\partial}{\partial t}$ with

$$\begin{aligned} A_u - F(u)B_v &= 0, \\ A_v - B_u &= 0. \end{aligned} \tag{62}$$

In the symmetry Theorems 1 and 2, let $x_1 = t$, $x_2 = x$, $u^1 = u$, $u^2 = v$. Then the conditions in Theorem 1 yield $F^1 = B(u, v)$, $F^2 = A(u, v)$, $\alpha_1^1 = \alpha_2^2 = 1$, $\alpha_2^1 = \alpha_1^2 = 0$, $\beta_v^\sigma = 0$ with the linear system given by (62) for $(X_1, X_2) = (u, v)$. Here system (33) becomes $\Phi_t = \Phi_x = 0$ with functionally independent solutions $X_1 = u$, $X_2 = v$; system (34) becomes $\frac{\partial \psi^1}{\partial t} = \frac{\partial \psi^2}{\partial x} = 1$, $\frac{\partial \psi^1}{\partial x} = \frac{\partial \psi^2}{\partial t} = 0$ and has as a solution $\psi^1 = t$, $\psi^2 = x$. This yields the well-known hodograph transformation $z_1 = u$, $z_2 = v$, $w^1 = t$, $w^2 = x$ that maps (61) to the linear system $x_u - F(u)t_v = 0$, $x_v - t_u = 0$.

On the other hand, system (61) admits multipliers $[\Lambda^1, \Lambda^2] = [a(U, V), b(U, V)]$ satisfying the linear system of PDEs

$$\begin{aligned} b_V + a_U &= 0, \\ b_U + F(U)a_V &= 0. \end{aligned} \tag{63}$$

Observe that system (63) is the adjoint of system (62).

Example 6. The NLT potential system

$$\begin{aligned} v_t - u^{-2}u_x - u^{-1} &= 0, \\ v_x - u_t &= 0, \end{aligned} \tag{64}$$

admits the point symmetry $X = -u^{-1}A(\hat{u}, v)\frac{\partial}{\partial x} + B(\hat{u}, v)\frac{\partial}{\partial t} + A(\hat{u}, v)\frac{\partial}{\partial u}$, $\hat{u} = x + \log u$, with

$$\begin{aligned} A_v + B_{\hat{u}} &= 0, \\ A_{\hat{u}} + B_v - A &= 0. \end{aligned} \tag{65}$$

Here the conditions in Theorem 1 yield $F^1 = B(\hat{u}, v)$, $F^2 = A(\hat{u}, v)$, $\alpha_1^1 = 1$, $\alpha_2^2 = -u^{-1}$, $\alpha_1^2 = \alpha_2^1 = 0$, $\beta_1^2 = 1$, $\beta_1^1 = \beta_2^2 = \beta_2^1 = 0$ with the linear system given by (65) for $(X_1, X_2) = (\hat{u} = x +$

$\log u, v$). System (33) becomes $\Phi_t = 0$, $\Phi_u - u^{-1}\Phi_x = 0$ with functionally independent solutions $X_1 = \hat{u}$, $X_2 = v$; system (34) becomes $\frac{\partial\psi^1}{\partial t} = 1$, $\frac{\partial\psi^1}{\partial u} - u^{-1}\frac{\partial\psi^1}{\partial x} = 0$, $\frac{\partial\psi^2}{\partial t} = 0$, $\frac{\partial\psi^2}{\partial u} - u^{-1}\frac{\partial\psi^2}{\partial x} = 1$ and has as a solution $\psi^1 = t$, $\psi^2 = u$. This yields the transformation $z_1 = \hat{u} = x + \log u$, $z_2 = v$, $w^1 = t$, $w^2 = u$ that maps (64) to the linear system $\frac{\partial w^2}{\partial z_2} + \frac{\partial w^1}{\partial z_1} = 0$, $\frac{\partial w^2}{\partial z_1} + \frac{\partial w^1}{\partial z_2} - w^2 = 0$.

On the other hand, system (64) admits multipliers for conservation laws given by $[\Lambda^1, \Lambda^2] = [Ua(\hat{U}, V), b(\hat{U}, V)]$, $\hat{U} = x + \log U$, with

$$\begin{aligned} a_V + b_{\hat{U}} &= 0, \\ a_{\hat{U}} + b_V + a &= 0. \end{aligned} \tag{66}$$

Again observe that system (66) is the adjoint of system (65).

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